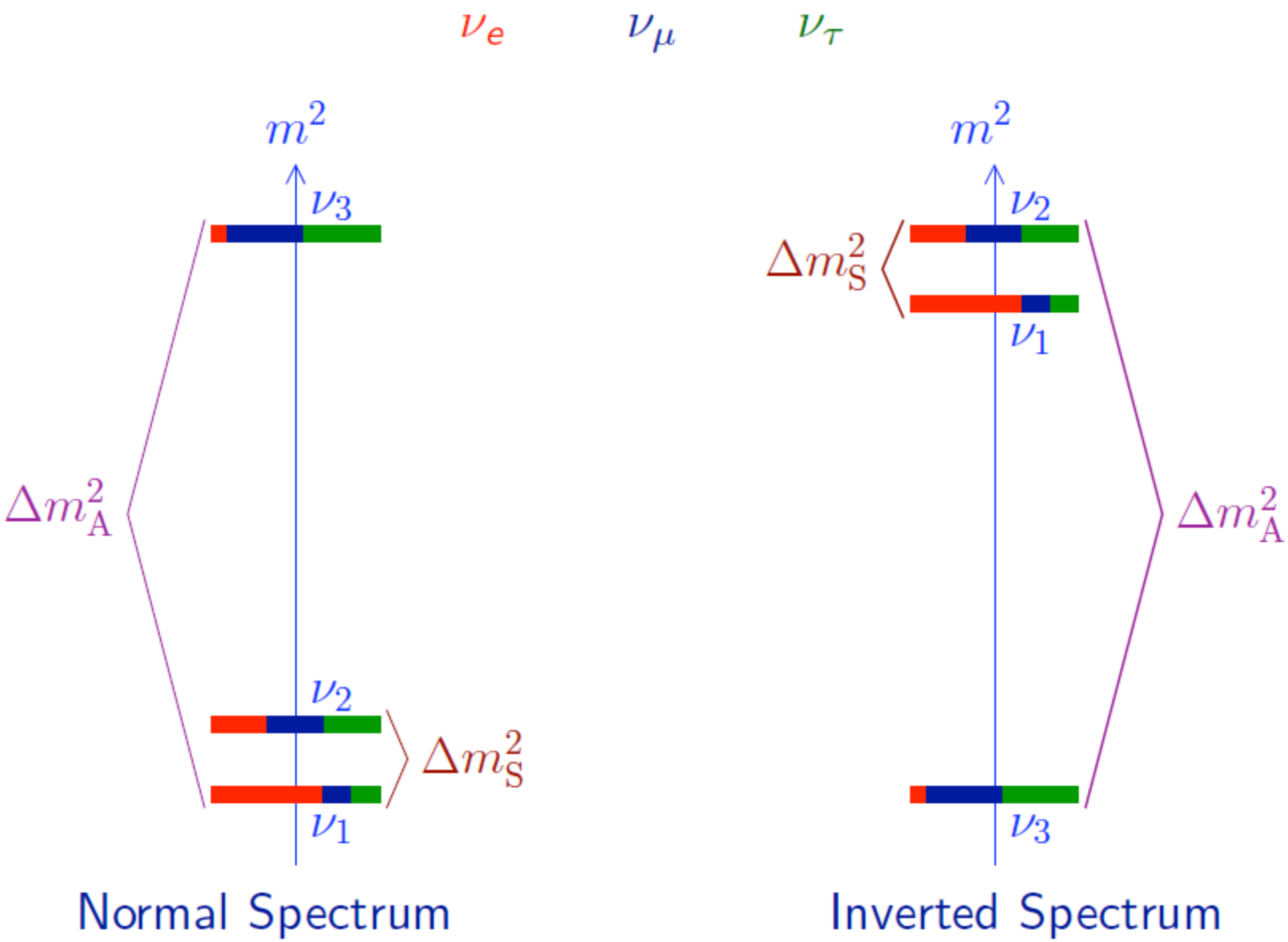
Lecture 1: Direct measurement of neutrino masses

PhD Cycle XXXIV

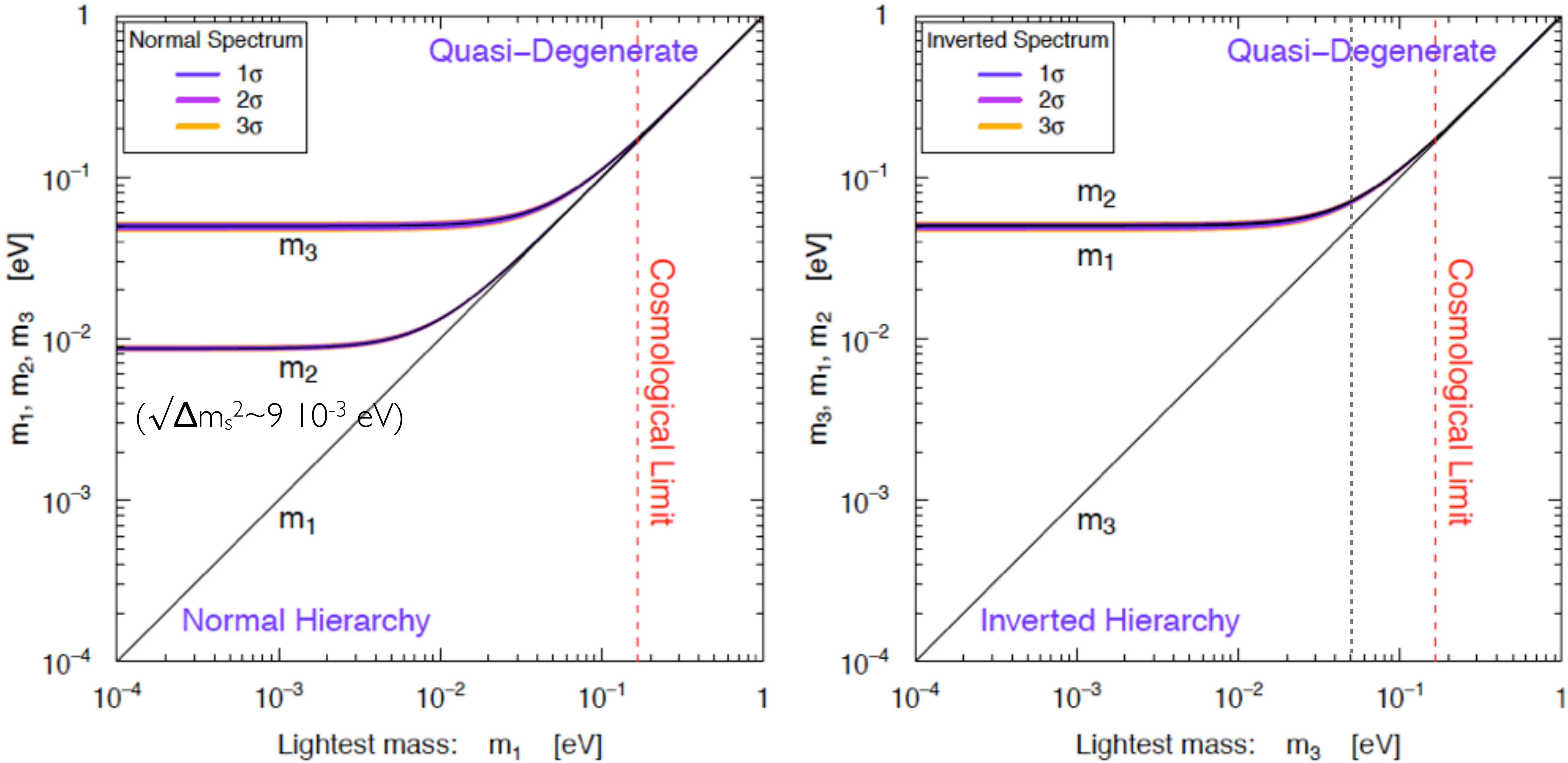
Three-Neutrino Mixing Paradigm



absolute scale is not determined by neutrino oscillation data

Parameter	best-fit
Δm^2_{21} $[10^{-5}$ eV $^2]$	7.37
$\Delta m^2_{31(23)}$ $[10^{-3}$ eV $^2]$	2.56 (2.54)

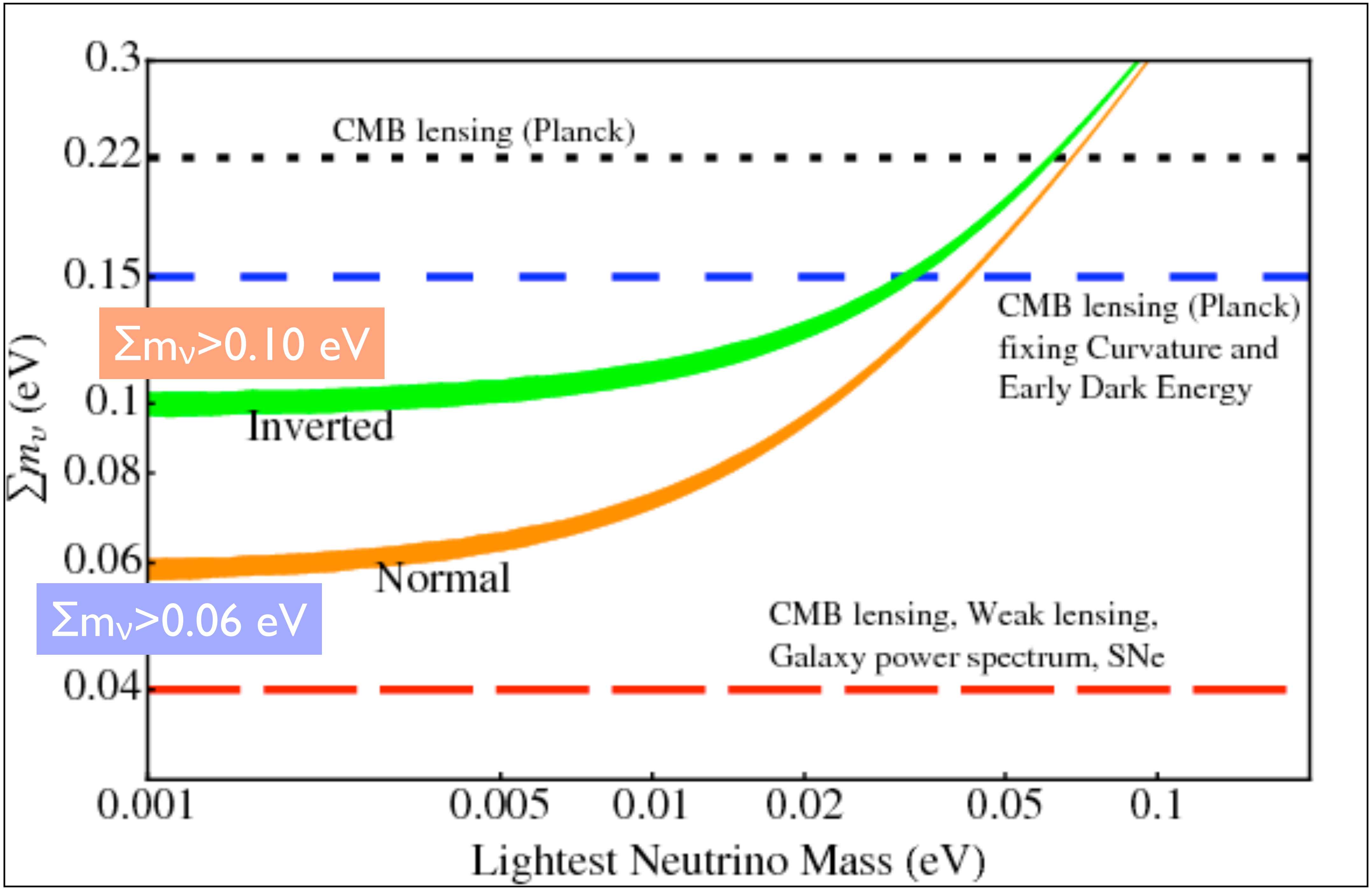
Absolute Values of Neutrino Masses



$$m_2^2 = m_1^2 + \Delta m_{21}^2 = m_1^2 + \Delta m_S^2$$
$$m_3^2 = m_1^2 + \Delta m_{31}^2 = m_1^2 + \Delta m_A^2$$

$$m_1^2 = m_3^2 - \Delta m_{31}^2 = m_3^2 + \Delta m_A^2$$
$$m_2^2 = m_1^2 + \Delta m_{21}^2 \simeq m_3^2 + \Delta m_A^2$$

Quasi-Degenerate for $m_1 \simeq m_2 \simeq m_3 \simeq m_\nu \gg \sqrt{\Delta m_A^2} \simeq 5 \times 10^{-2}\,\text{eV}$



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Observables sensitive to m_ν

The absolute mass scale can be measured through:
(numbers on the right are current upper limits)

- tritium beta decay

$$m_\beta \equiv \left[\sum |U_{ei}|^2 m_i^2\right]^{1/2}$$

(2.05 – 2.3 eV @ 95%CL)

- neutrinoless double beta decay

$$m_{\beta\beta} \equiv \left|\sum U_{ei}^2 m_i\right|$$

(0.06 – 0.16 eV @ 90%CL)

- cosmological observations

$$\sum m_\nu \equiv \sum_i m_i$$

(0.2 – 0.7 eV @ 95%CL)

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TODAY

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NOT HERE

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LAST LECTURE

$$\sum m_\nu \equiv \sum_i m_i$$ (0.2 – 0.7 eV @ 95%CL)

Direct mass measurements

Beta decay: direct ν_e mass

The most sensitive known method to measure the electron neutrino mass is by observing the electron spectrum in nuclear β-decay

$$\mathcal{N}(A,Z) \to \mathcal{N}(A,Z+1) + e^- + \bar{\nu}_e \,, \tag{14.1}$$

where A and Z are, respectively, the mass and atomic numbers of the parent nucleus.

As we have seen in sections 6.1 and 6.2.1, the electron neutrino, in general, does not have a definite mass, but is a mixture of massive neutrinos. However, following the tradition, in this section we treat the electron neutrino as a mass eigenstate. We will discuss the effects of neutrino mixing in nuclear β-decay in section 14.1.1.

The differential decay rate in allowed[74] β-decays is given by

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}E_e} = \frac{G_\mathrm{F}^2\, m_e^5}{2\,\pi^3} \cos^2\theta_\mathrm{C}\, |\mathcal{M}|^2\, F(Z,E_e)\, E_e\, p_e\, E_\nu\, p_\nu \,, \tag{14.2}$$

where θ_C is the Cabibbo angle, $\mathcal{M}$ is the nuclear matrix element, E_e (E_ν) and p_e (p_ν) are the electron (neutrino) energy and momentum, and $F(Z,E_e)$ is the Fermi

the surrounding electrons). The two factors $E_i\, p_i$ in eqn (14.2), with $i = e, \nu$, come from the phase-space factor of the final state: $\mathrm{d}^3p_i = p_i^2\,\mathrm{d}p_i\,\mathrm{d}\cos\theta_i\,\mathrm{d}\phi_i = p_i\,E_i\,\mathrm{d}E_i\,\mathrm{d}\cos\theta_i\,\mathrm{d}\phi_i$, where θ_i and ϕ_i are the polar angular coordinates of $\vec{p}_i$.

Fundamentals of Neutrino Physics and Astrophysics

by: Carlo Giunti and Chung W. Kim, Oxford UP 2007

massless. On the other hand, if the electron neutrino has a mass m_{ν_e}, the maximal kinetic energy of the electron is

$$T_{\max} = Q_\beta - m_{\nu_e} \quad (14.5)$$

Since the neutrino momentum is given by

$$p_\nu = \sqrt{E_\nu^2 - m_{\nu_e}^2} = \sqrt{(Q_\beta - T)^2 - m_{\nu_e}^2}\,, \quad (14.6)$$

the differential decay rate in eqn (14.2) can be written, for $T \leq T_{\max}$, as[75]

$$\frac{d\Gamma}{dT} = \frac{G_F^2\, m_e^5}{2\,\pi^3} \cos^2\theta_C\, |\mathcal{M}|^2\, F(Z, E_e)\, E_e\, p_e\, (Q_\beta - T) \sqrt{(Q_\beta - T)^2 - m_{\nu_e}^2}\,, \quad (14.8)$$

➡ What we measure:

Kurie plot

$$K(T) = \sqrt{\frac{d\Gamma/dT}{\frac{(\cos\vartheta_C G_F)^2}{2\pi^3}|\mathcal{M}|^2 F(E)\, pE}} = \left[(Q - T)\sqrt{(Q-T)^2 - m_{\nu_e}^2}\right]^{1/2}$$

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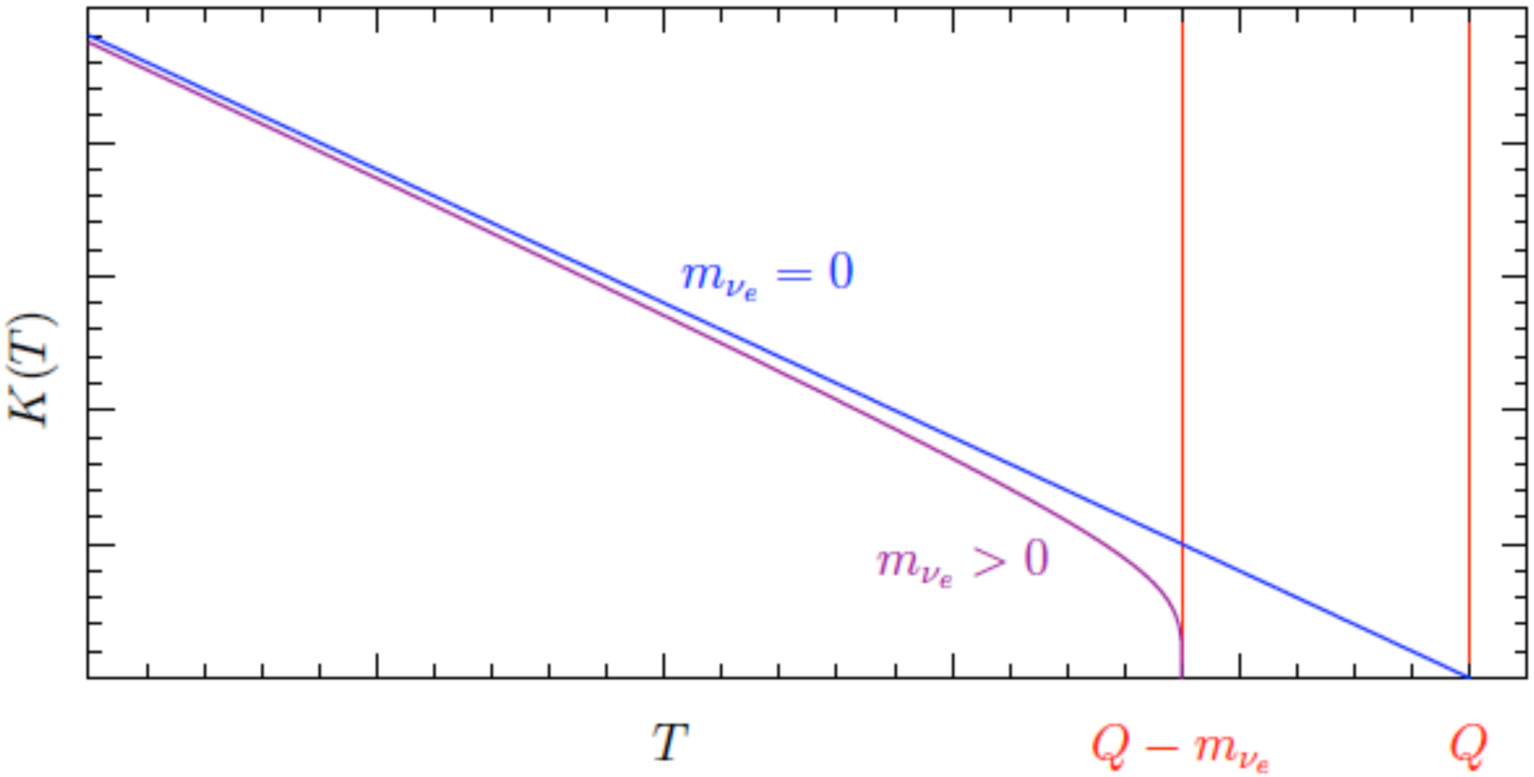


$m_{\nu_e} < 2.2\,\mathrm{eV}$ (95% C.L.)

Mainz & Troitsk

[Weinheimer, hep-ex/0210050]

future: KATRIN

[www.katrin.kit.edu]

sensitivity: $m_{\nu_e} \simeq 0.2\,\mathrm{eV}$

C. Giunti – Neutrino Mass: Overview of $\beta\beta_{0\nu}$, Cosmology and Direct Measurements – 14 May 2012 – 4/17

The mass is extracted from a fit to the end-point

Pros and cons

end-point of the electron spectrum is that very few events occur near the end-point. We can estimate the relative number of events occurring in an interval of energy ΔT below the end-point as follows. Below the end-point we have

$$T \simeq Q_\beta \quad \Longrightarrow \quad \begin{cases} E_e \simeq Q_\beta + m_e \\ p_e = \sqrt{E_e^2 - m_e^2} \simeq \sqrt{Q_\beta \, (Q_\beta + 2\, m_e)} \, . \end{cases} \tag{14.9}$$

Ignoring the neutrino mass and the Fermi function, we have

$$\left. \frac{\mathrm{d}\Gamma}{\mathrm{d}T} \right|_{\substack{T \simeq Q_\beta \\ m_{\nu_e} = 0}} \propto (Q_\beta + m_e) \, \sqrt{Q_\beta \, (Q_\beta + 2\, m_e)} \, (Q_\beta - T)^2 \, , \tag{14.10}$$

and

$$\int_{Q_\beta - \Delta T}^{Q_\beta} \frac{\mathrm{d}\Gamma}{\mathrm{d}T} \, \mathrm{d}T \propto (Q_\beta + m_e) \, \sqrt{Q_\beta \, (Q_\beta + 2\, m_e)} \, (\Delta T)^3 \, . \tag{14.11}$$

The total number of events is proportional to

$$\int_0^{Q_\beta} \frac{\mathrm{d}\Gamma}{\mathrm{d}T} \, \mathrm{d}T \propto \int_0^{Q_\beta} (T + m_e) \, \sqrt{T \, (T + 2 m_e)} \, (Q_\beta - T)^2 \, \mathrm{d}T \, , \tag{14.12}$$

where we have neglected again the Fermi function and the neutrino mass. Since we are interested in an order-of-magnitude estimate, we consider $Q_\beta \gg m_e$, which leads to the approximation

$$\int_0^{Q_\beta} \frac{\mathrm{d}\Gamma}{\mathrm{d}T} \, \mathrm{d}T \propto Q_\beta^5 \, . \tag{14.13}$$

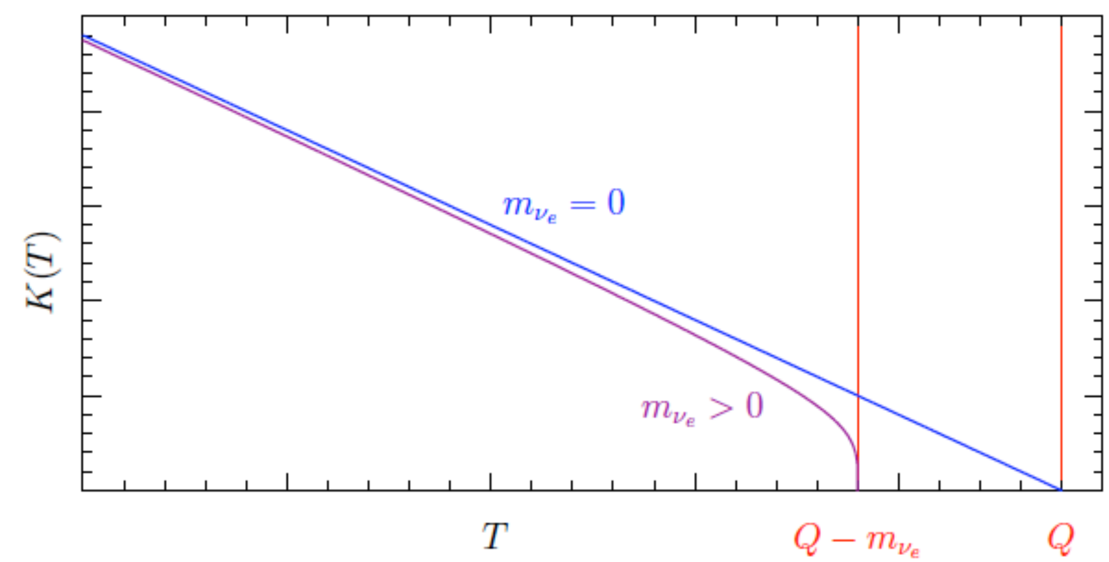
Thus, the relative number of events occurring in an interval of energy ΔT below the end-point is given by

$$\frac{n(\Delta T)}{n} \propto \left(\frac{\Delta T}{Q_\beta} \right)^3 \, . \tag{14.14}$$

One can obtain the same result considering $Q_\beta \ll m_e$. From eqn (14.14) it is clear that in order to maximi…

+ : This method relies **purely** on 3-body kinematics, without any assumption on the nature of the **ν** (e.g. Dirac/Majorana)

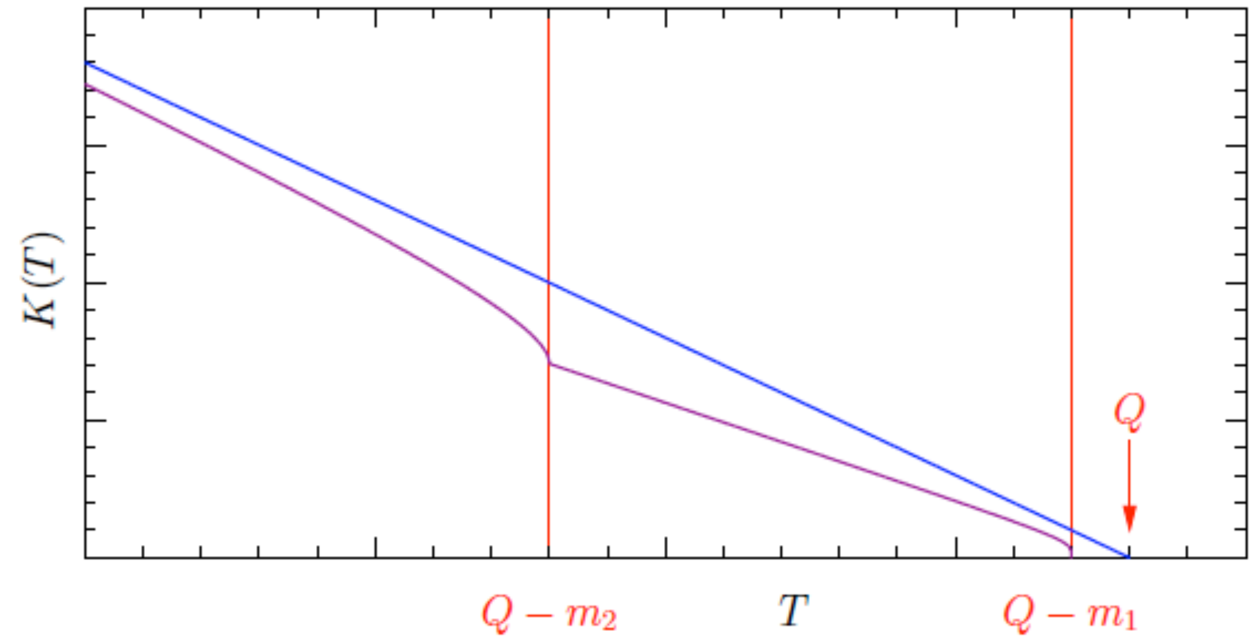
- : statistics, especially at the end point



The smaller the Q_β, the better for the relative stat error just below the end-point

Added complication: mix

Neutrino Mixing $\Longrightarrow K(T) = \left[(Q-T)\sum_k |U_{ek}|^2 \sqrt{(Q-T)^2 - m_k^2}\right]^{1/2}$



analysis of data is different from the no-mixing case: $2N-1$ parameters

$$\left(\sum_k |U_{ek}|^2 = 1\right)$$

if experiment is not sensitive to masses ($m_k \ll Q - T$)

effective mass: $$m_\beta^2 = \sum_k |U_{ek}|^2 m_k^2$$

$$K^2 = (Q-T)^2 \sum_k |U_{ek}|^2 \sqrt{1 - \frac{m_k^2}{(Q-T)^2}} \simeq (Q-T)^2 \sum_k |U_{ek}|^2 \left[1 - \frac{1}{2}\frac{m_k^2}{(Q-T)^2}\right]$$

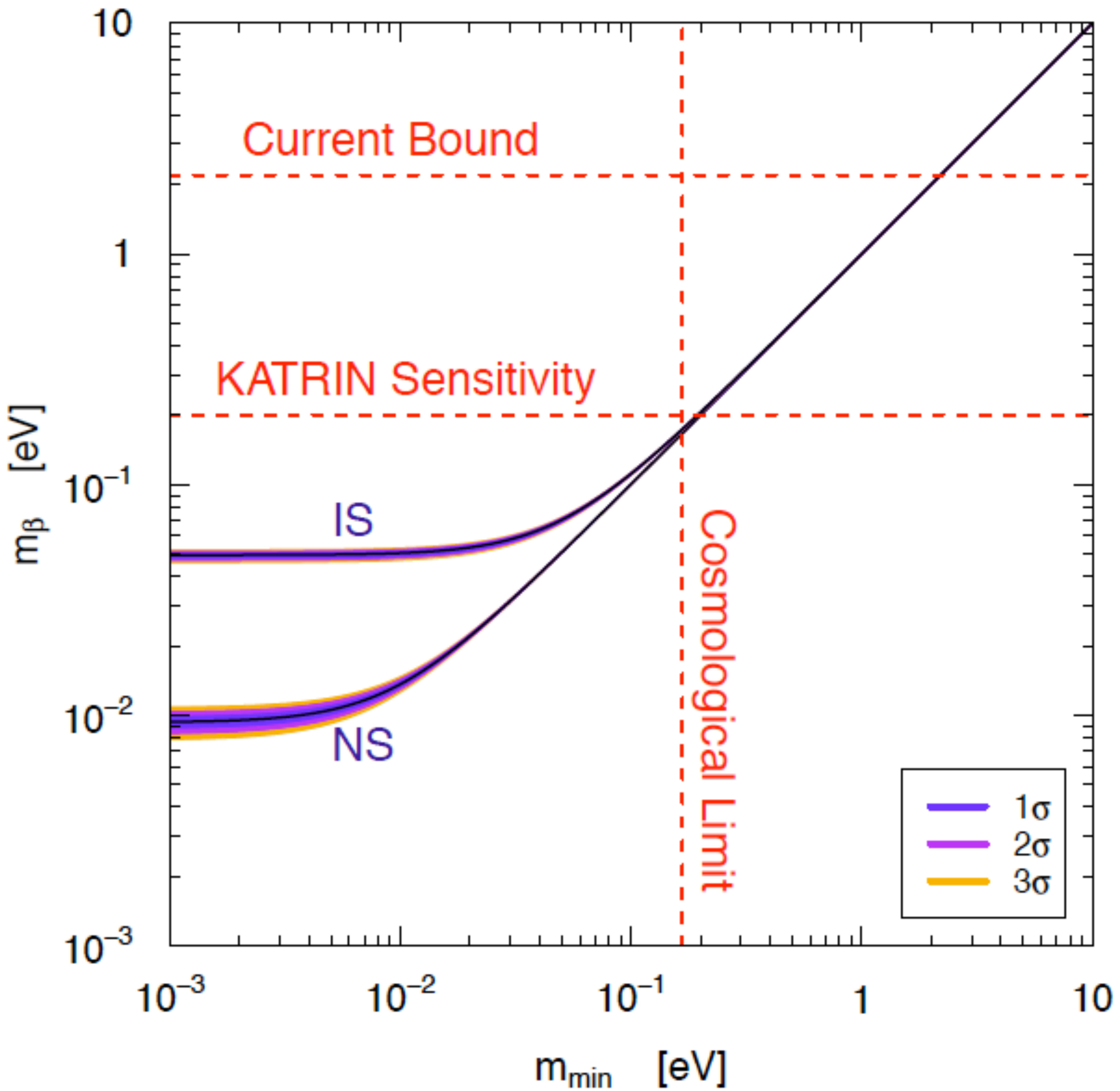
$$= (Q-T)^2 \left[1 - \frac{1}{2}\frac{m_\beta^2}{(Q-T)^2}\right] \simeq (Q-T)\sqrt{(Q-T)^2 - m_\beta^2}$$

Kurie plot

$$K(T) = \sqrt{\frac{d\Gamma/dT}{\frac{(\cos\vartheta_C G_F)^2}{2\pi^3}|\mathcal{M}|^2 F(E)\, pE}} = \left[(Q-T)\sqrt{(Q-T)^2 - m_{\nu_e}^2}\right]^{1/2}$$

Predictions of 3ν-Mixing Paradigm

$$m_\beta^2 = |U_{e1}|^2\, m_1^2 + |U_{e2}|^2\, m_2^2 + |U_{e3}|^2\, m_3^2$$



- Quasi-Degenerate:

$$m_\beta^2 \simeq m_\nu^2 \sum_k |U_{ek}|^2 = m_\nu^2$$

- Inverted Hierarchy:

$$m_\beta^2 \simeq (1 - s_{13}^2)\Delta m_A^2 \simeq \Delta m_A^2$$

- Normal Hierarchy:

$$m_\beta^2 \simeq s_{12}^2 c_{13}^2 \Delta m_S^2 + s_{13}^2 \Delta m_A^2$$
$$\simeq 2 \times 10^{-5} + 6 \times 10^{-5}\, \text{eV}^2$$

- If $m_\beta \lesssim 4 \times 10^{-2}\,\text{eV}$

$\Downarrow$

Normal Spectrum

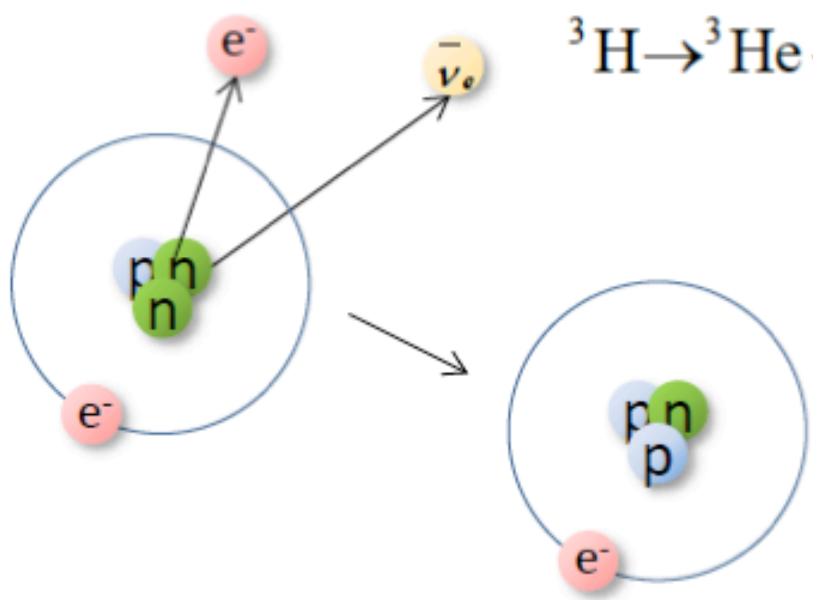
Remember: $\sqrt{\Delta m_A^2} \simeq 5 \times 10^{-2}\,\text{eV}$

Beta decay of ^{3}H

$$Q = M_{^3H} - M_{^3He} - m_e = 18.58\,\text{keV}$$

- $\tau_{1/2} \cong 12.3$ years (4*10^8 atoms for 1 Bq)

reasonably short lifetime

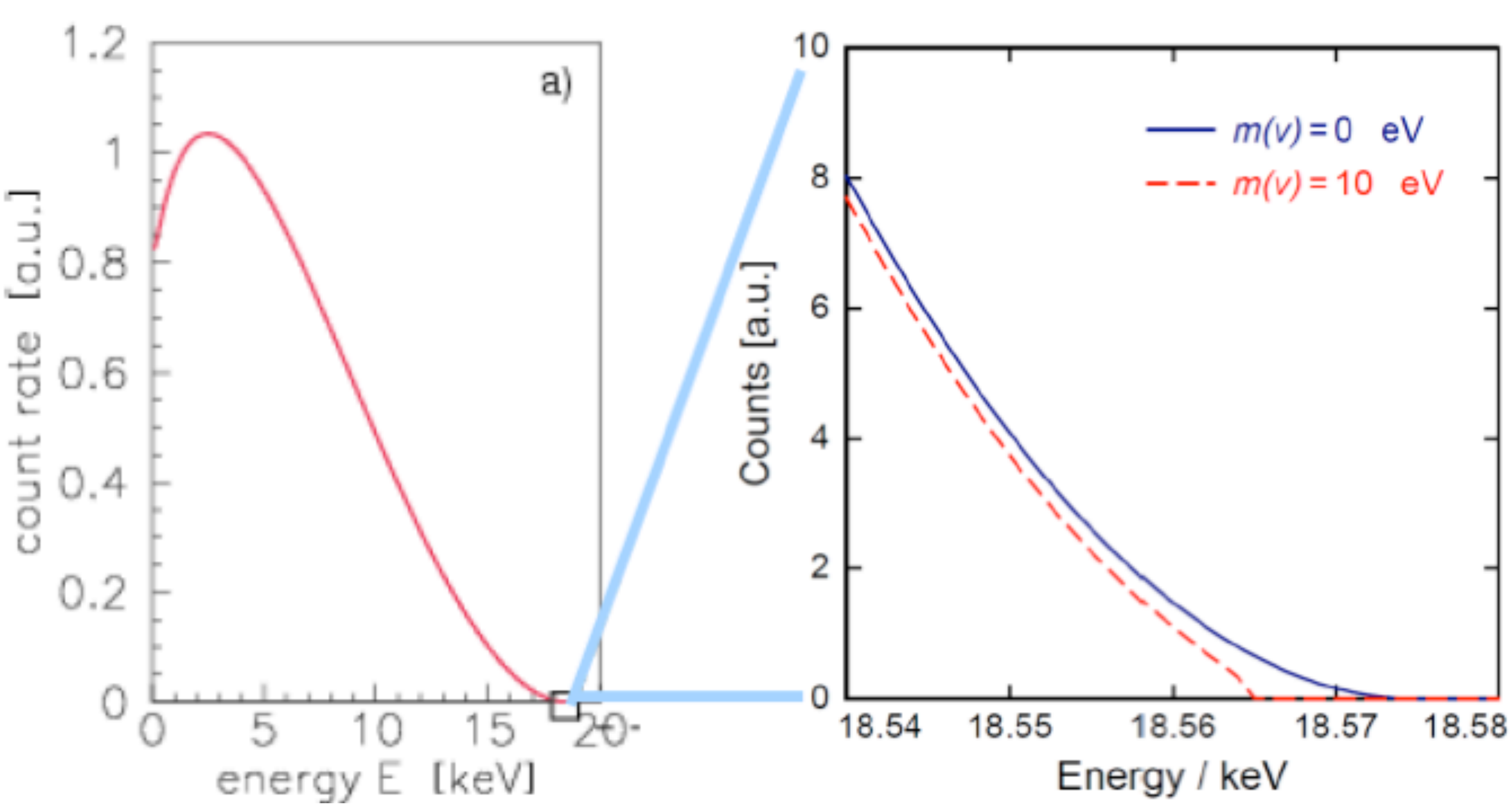


$$^3H \rightarrow {}^3He + e^- + \bar{\nu}_e$$

^{3}H: chosen because:

✓ low Q ⇒ enhanced $\frac{n(\Delta T)}{n} \propto \left(\frac{\Delta T}{Q_\beta}\right)^3$

✓ simple atomic structure (small uncertainties on $|\mathcal{M}|^2 F(Z, E_e)$)

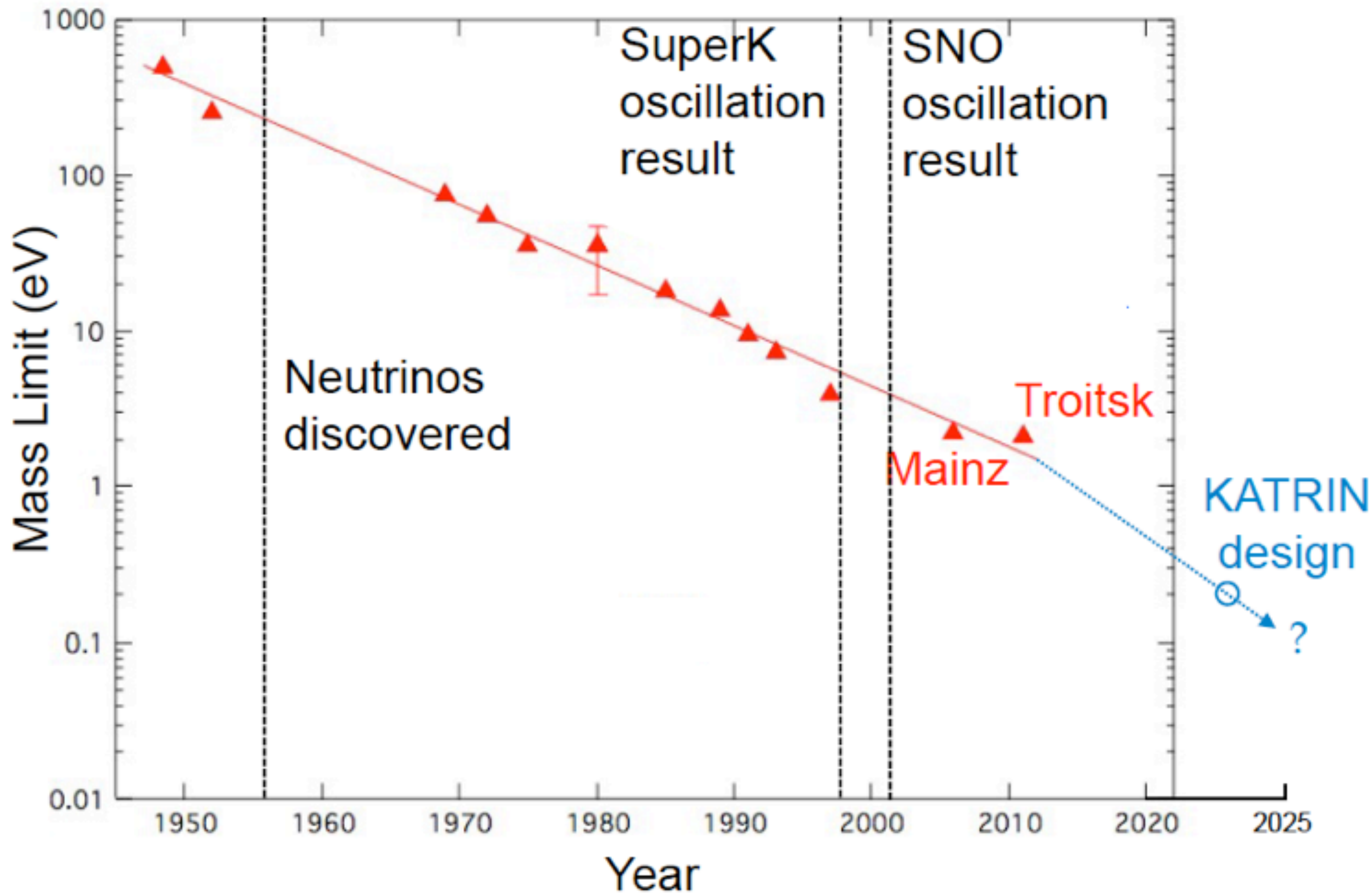


Only a small fraction of events in the last eV below the endpoint: 2 *10^{-13}

Triutium is present as bi-atomic molecules

TRITIUM

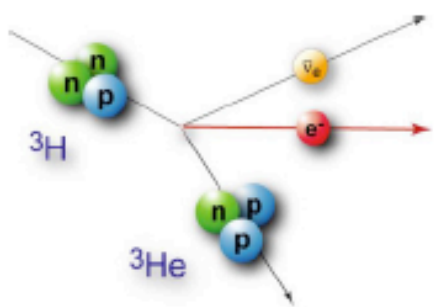
$m_{\nu,\mathrm{eff}}^2$: A Brief History in Tritium

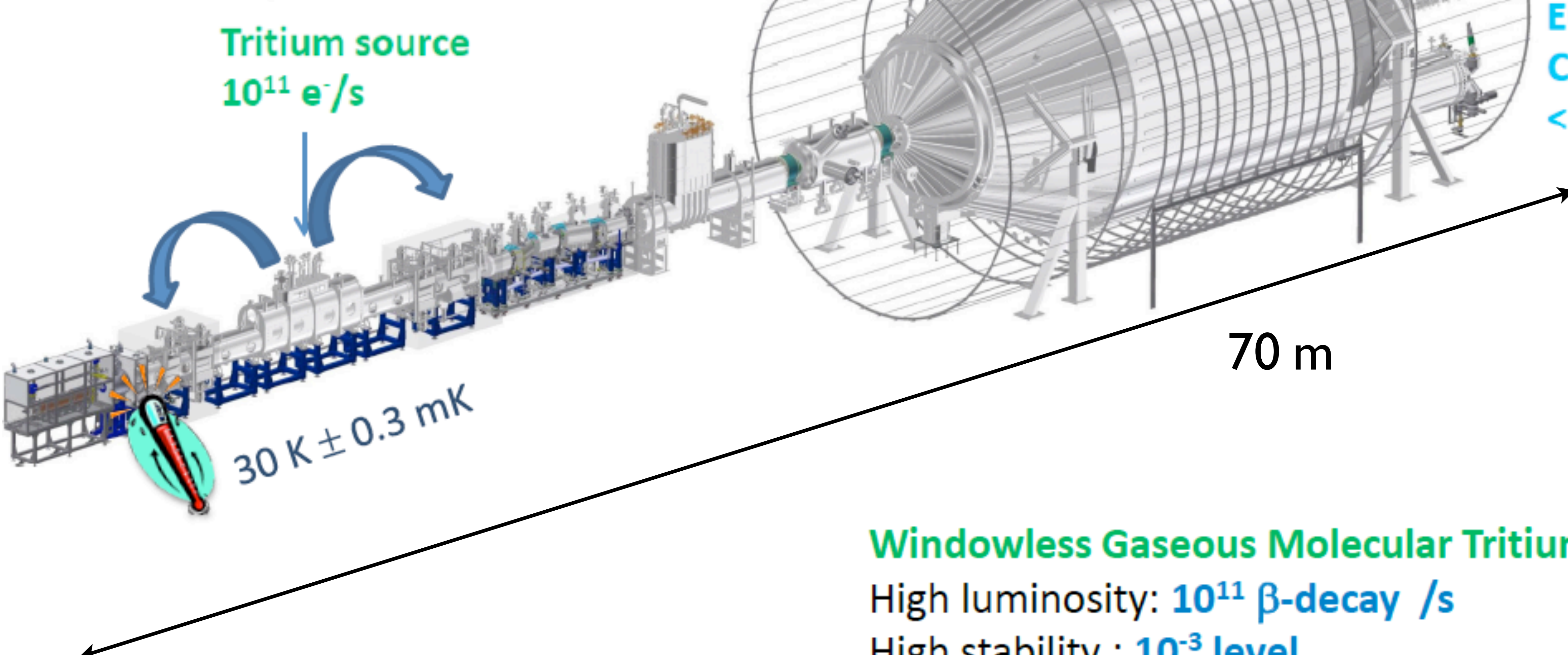


Adapted from J. Wilkerson, Neutrino 2012

The KATRIN experiment

The source should be transparent to
the emitted electrons: gas or thin layers





Windowless Gaseous Molecular Tritium Source

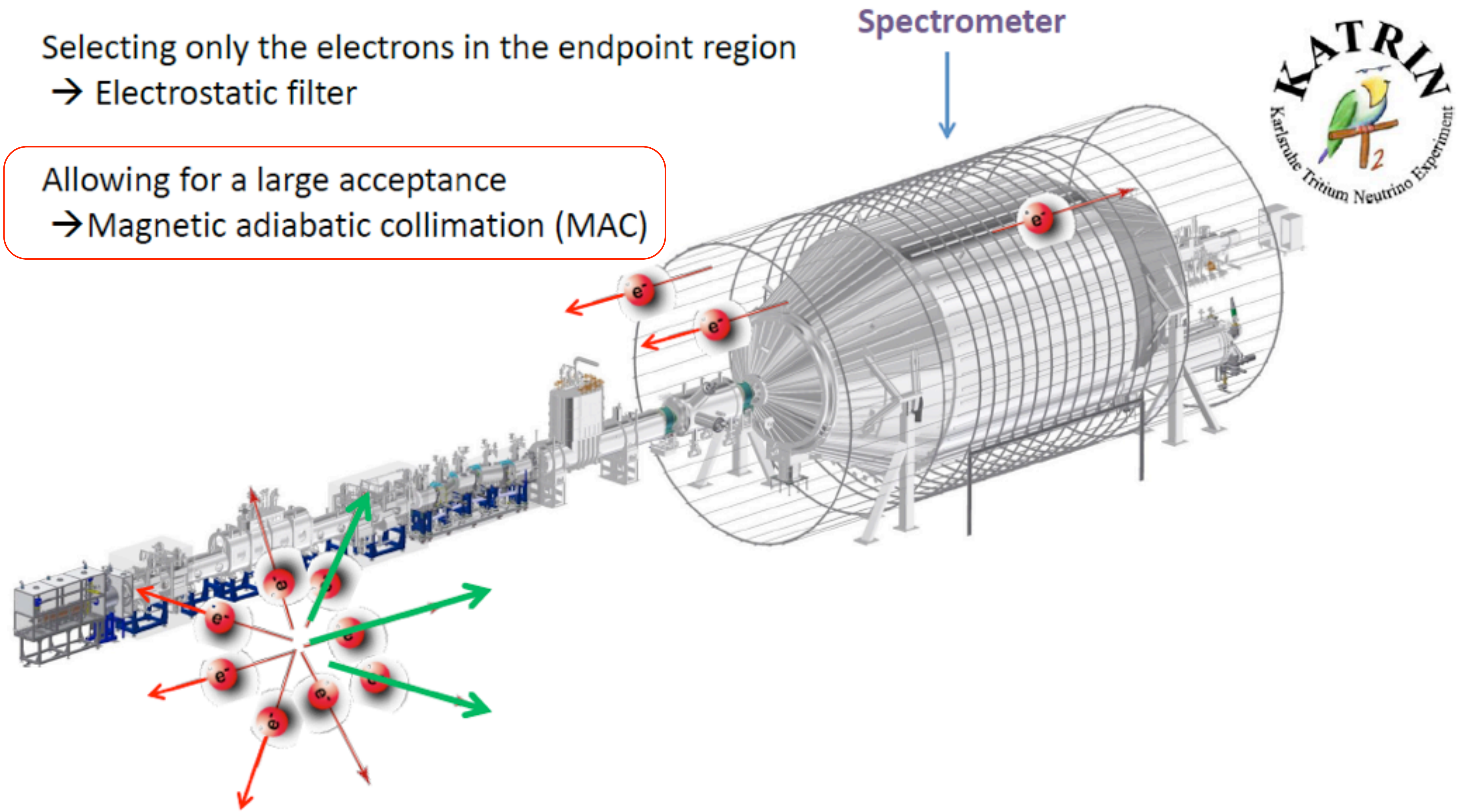
High luminosity: **10^{11} β-decay /s**

High stability : **10^{-3} level**

The KATRIN experiment

Selecting only the electrons in the endpoint region
→ Electrostatic filter

Allowing for a large acceptance
→Magnetic adiabatic collimation (MAC)



MAC-E filter principle allows for
< 1 eV energy cut off

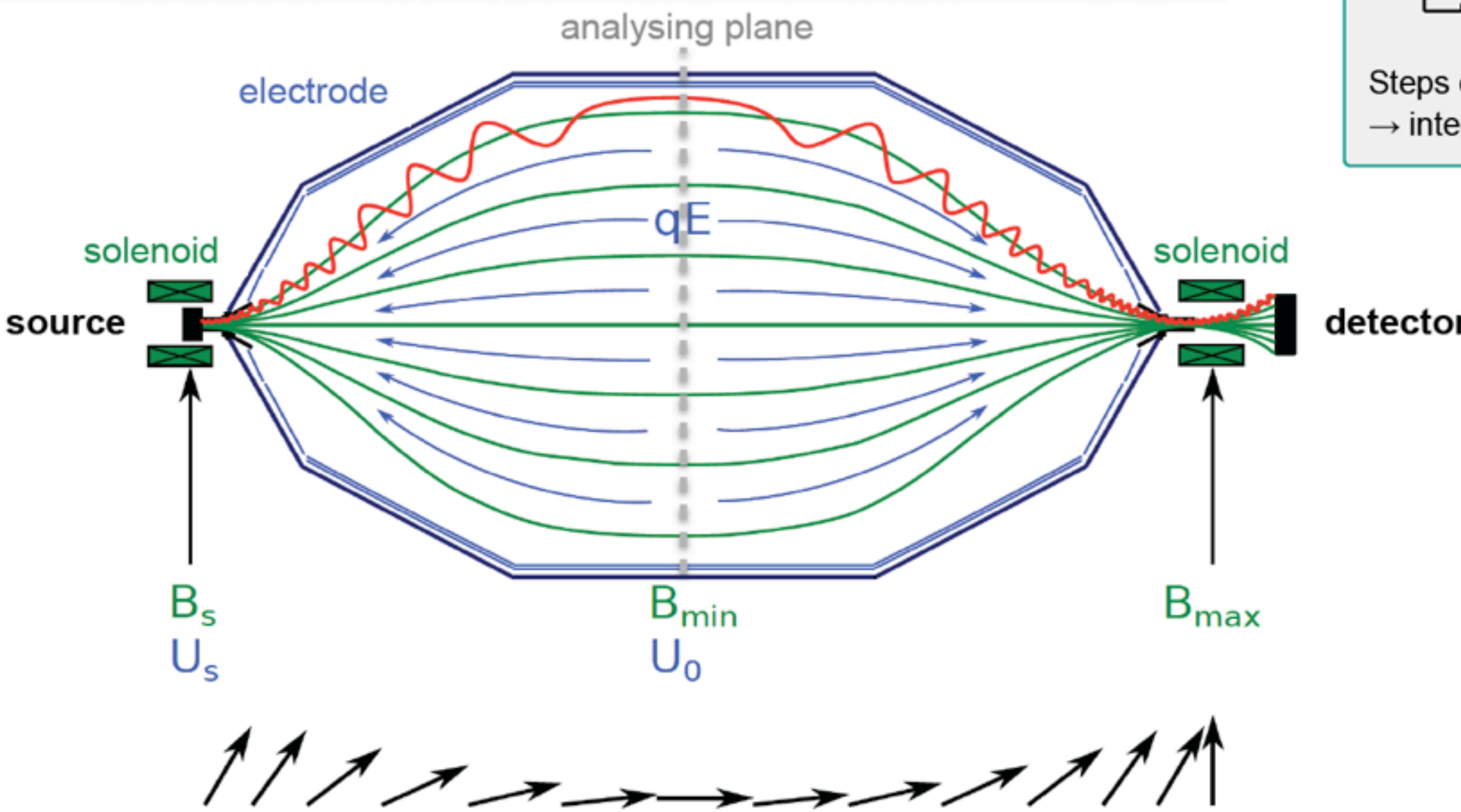
High-resolution β spectrometer

Karlsruhe Institute of Technology

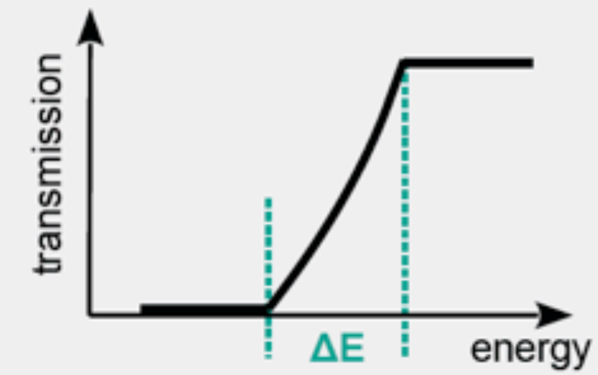
Magnetic Adiabatic Collimation & Electrostatic Filter

- ➜ integrating electrostatic filter ($E_{kin} > eU_0$)
- ➜ “clean” (analytic) response function
- ➜ $\Delta E < 1$ eV at 18.6 keV

$$\frac{\Delta E}{E} = \frac{B_{\mathrm{min}}}{B_{\mathrm{max}}}$$

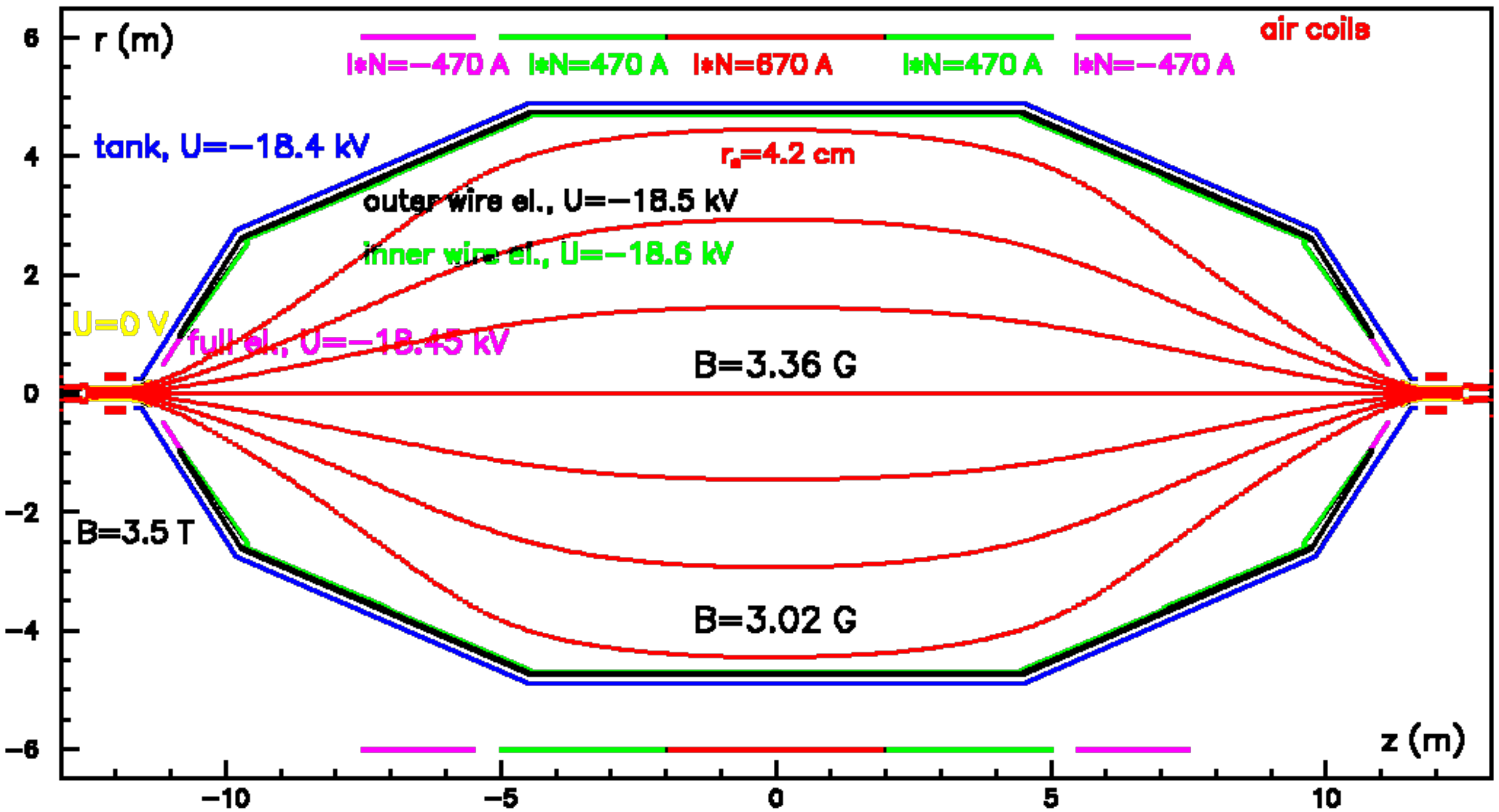


Sharp high-pass filter:

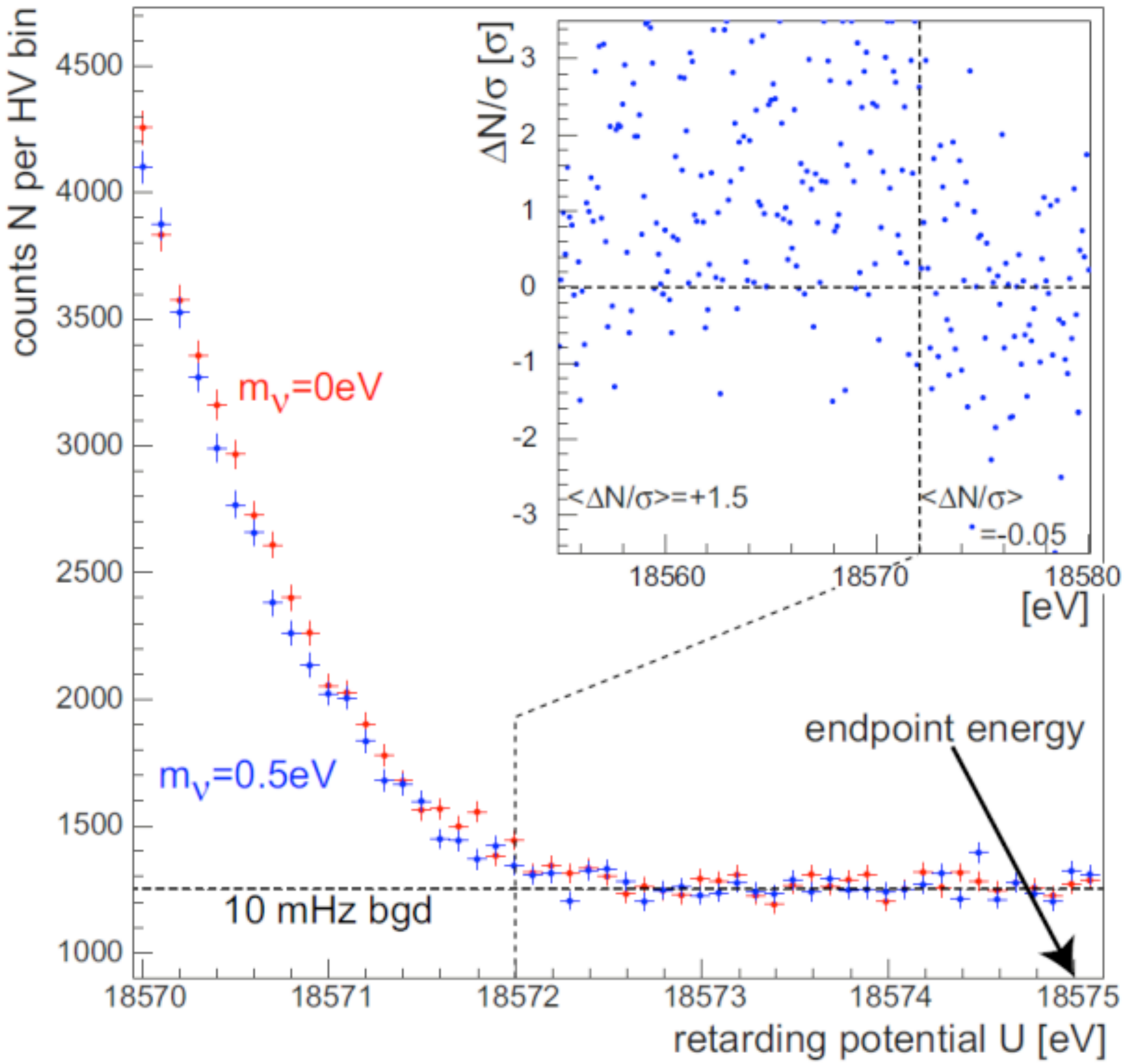


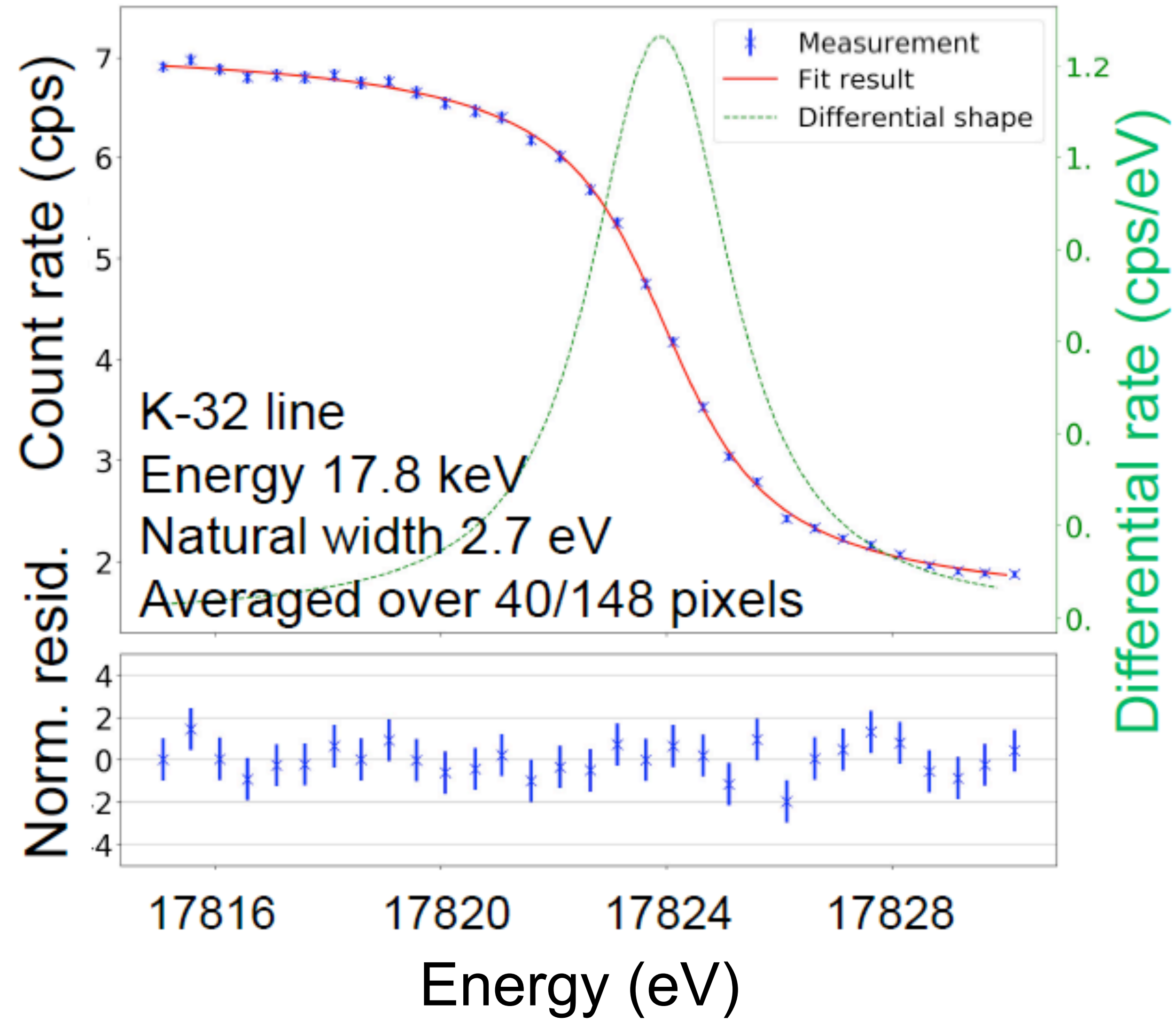
Steps of filter potential
→ integrated β spectrum

[Beamson et al. 1980; Kruit & Read 1983; **Lobashev 1985; Picard et al. 1992**]



- The larger the spectrometer, the smaller (more "adiabatic") the gradient of the e- momentum
- ~constant along B lines
- spectrometer acts as an **integrating** high-energy pass filter by virtue of E field (threshold effect)

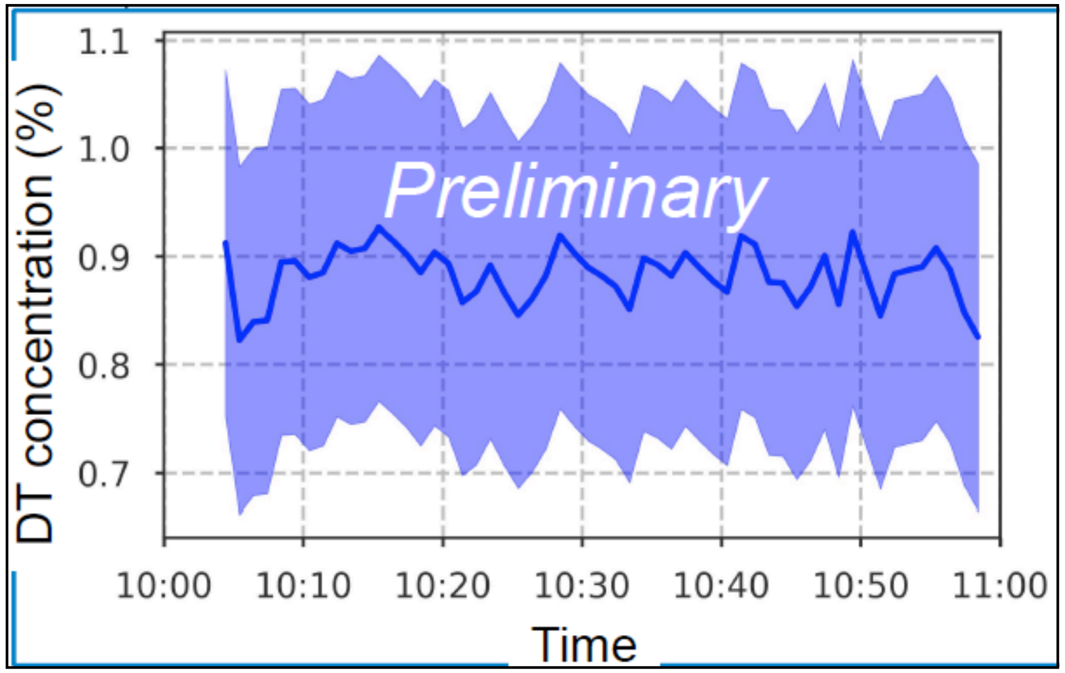


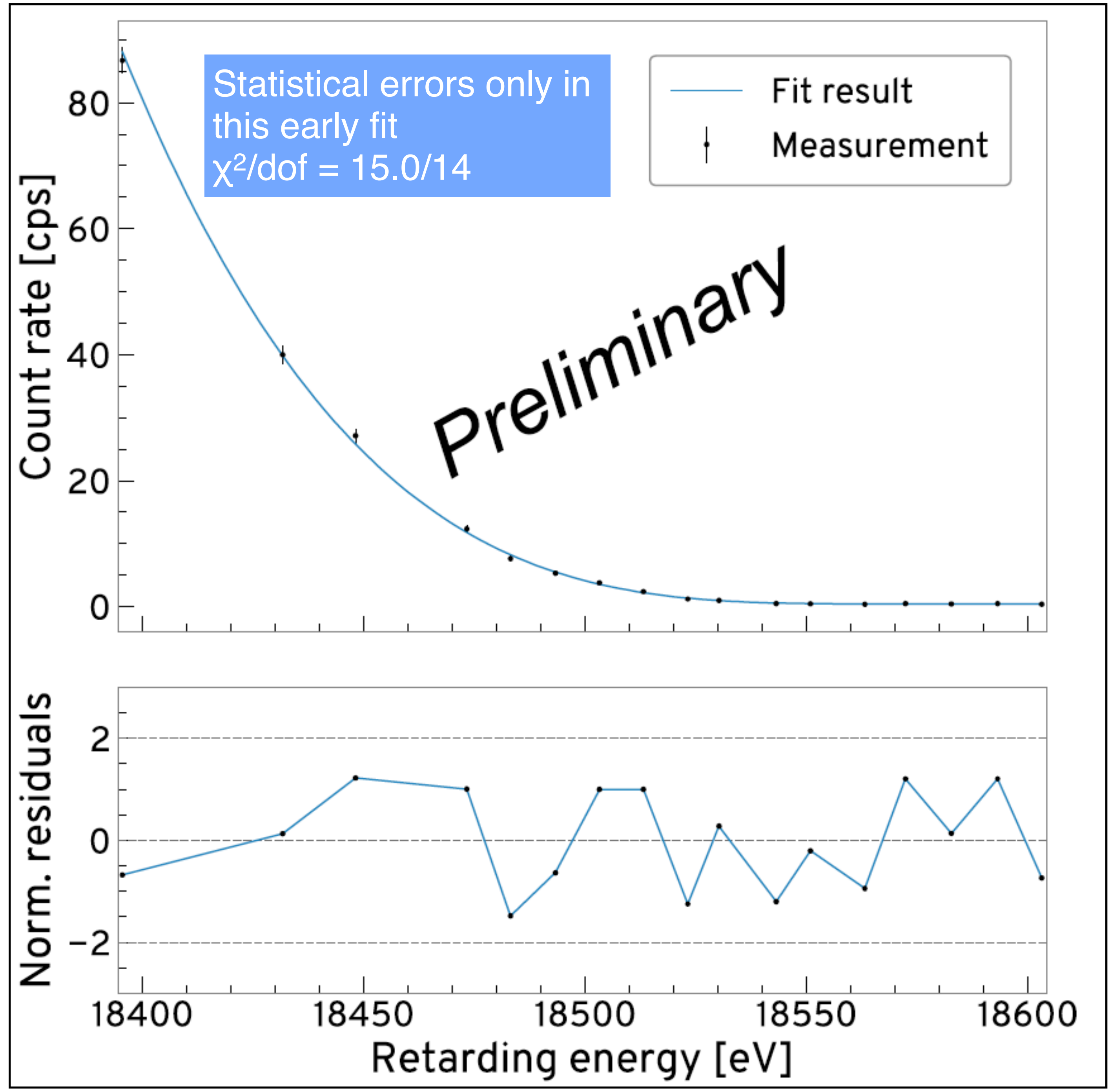


"Calibration": example with a real July 2017 pre-run with ^{83}Kr de-excitation lines

- **Integrated** e- spectrum
- Will measure N(e-) vs electric potential applied
- Measurement up to 30 eV below end-point (Q-value)

First tritium injection: 18th May, 7:48 am UTC





- Commissioning in past months
  - checks of temp and pressure, Tritium concentration stability
  - whether number of electrons over threshold as expected
  - absolute residual:
    - observed - expected (fit)

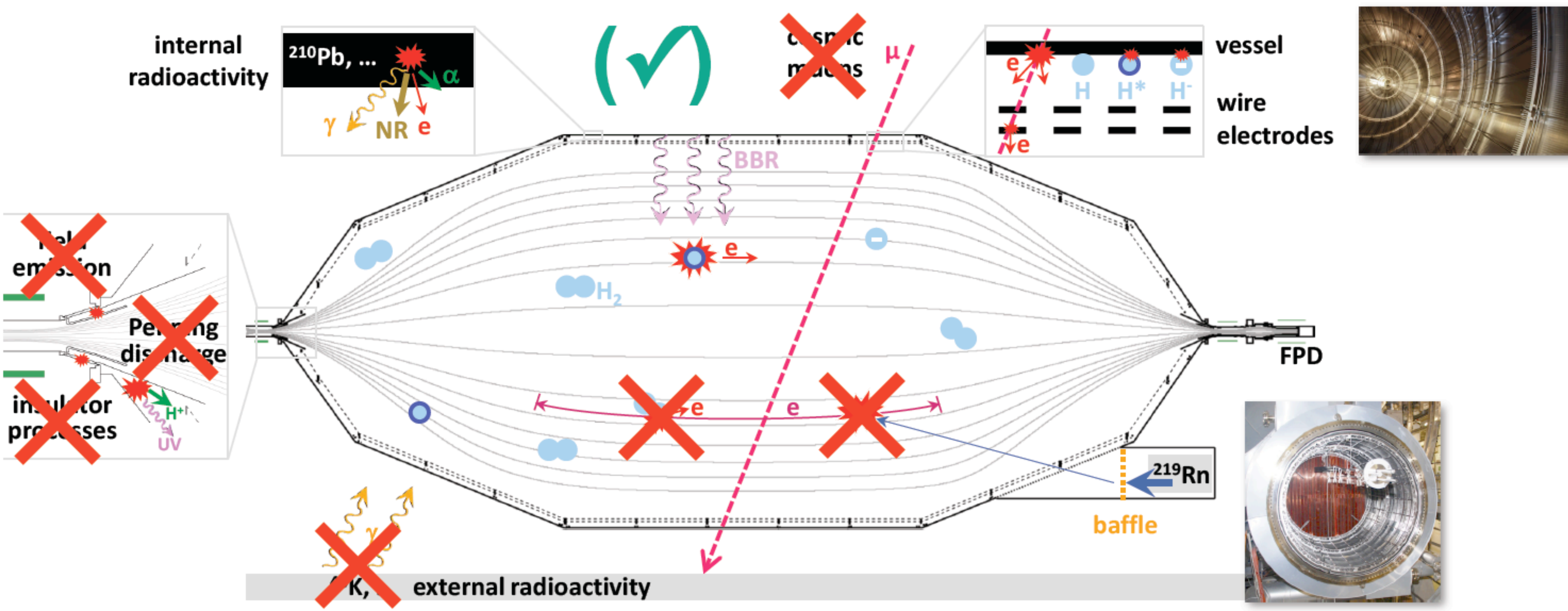
expect $m_{\nu,eff}$ data in early 2019

Main sources of background:

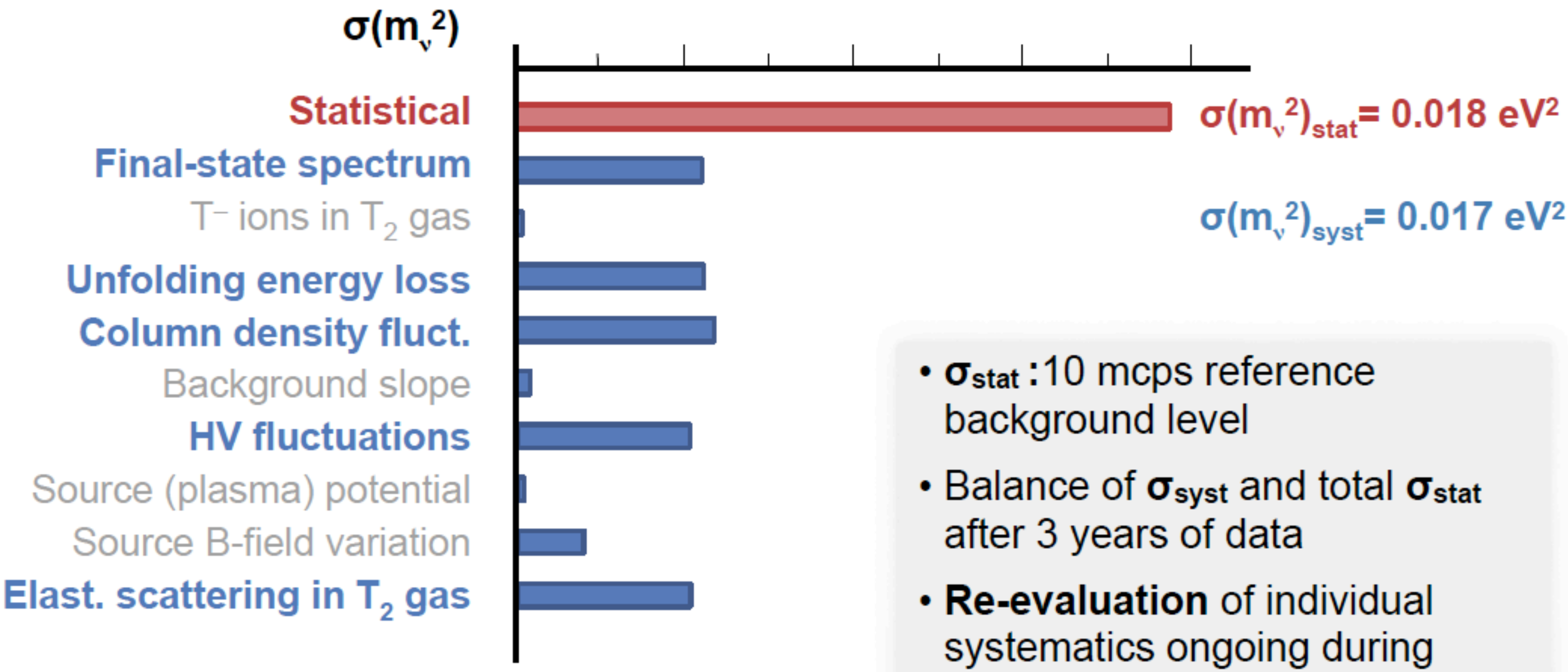
- electrons produced by ionization in the residual gas: veto system in place
- electrons from interactions of cosmic rays: veto system in place
- γ-rays from natural radioactivity emanating from material surrounding the detector and from the detector itself: this is still open (^{210}Pb on spectrometer walls)

Backgrounds vary over a large range of energies, β-electrons concentrated.

Rely on accurate E determination to separate and reject bkgs



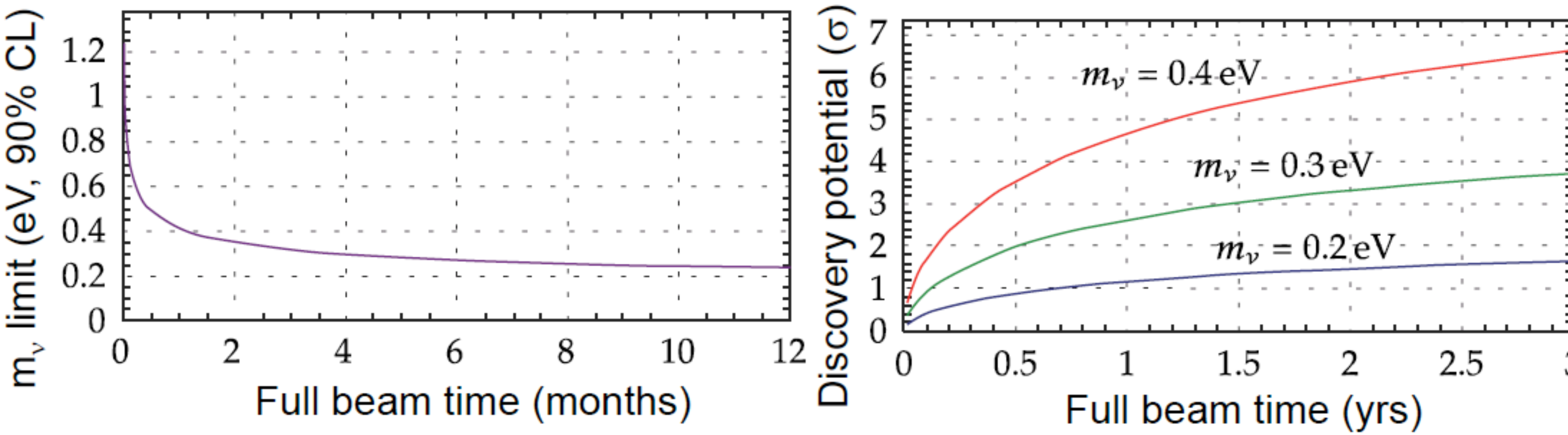
KATRIN's uncertainty budget (design sensitivity, ~2004):



- **σ_{stat}** :10 mcps reference background level
- Balance of **σ_{syst}** and total **σ_{stat}** after 3 years of data
- **Re-evaluation** of individual systematics ongoing during system characterisation

Old projections, but still believed to be mostly accurate

- Full sensitivity ($\sigma_{syst} = \sigma_{stat}$) after 3 beam years (~5 calendar years)



G. Drexlin et al., Adv. High Energy Phys. ***2013*** *(2013) 293986*

Excluding masses down to 0.24 eV possible after 3 years of full beam

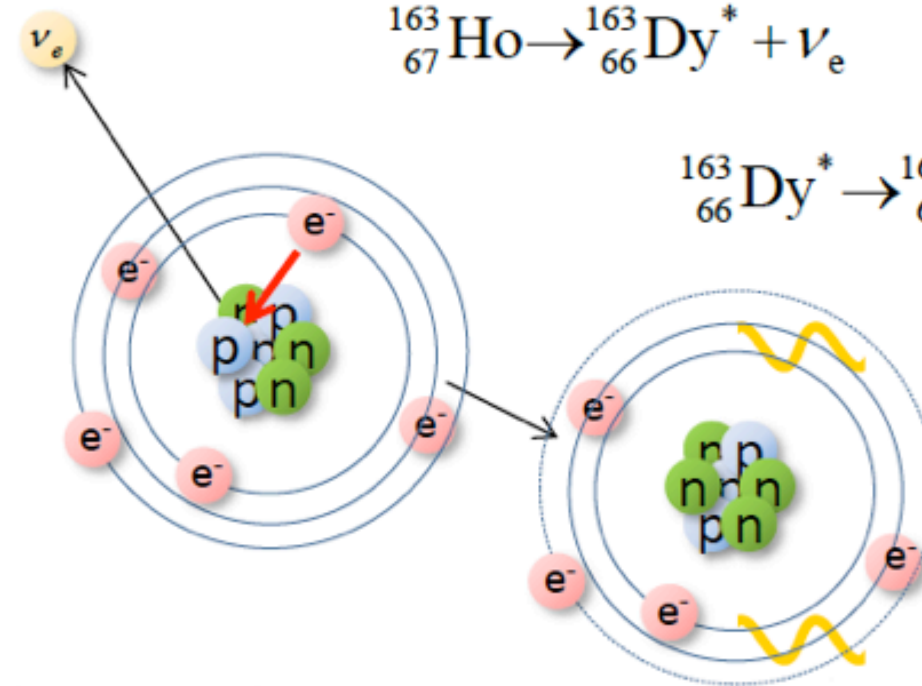
Electron capture in ^{163}Ho

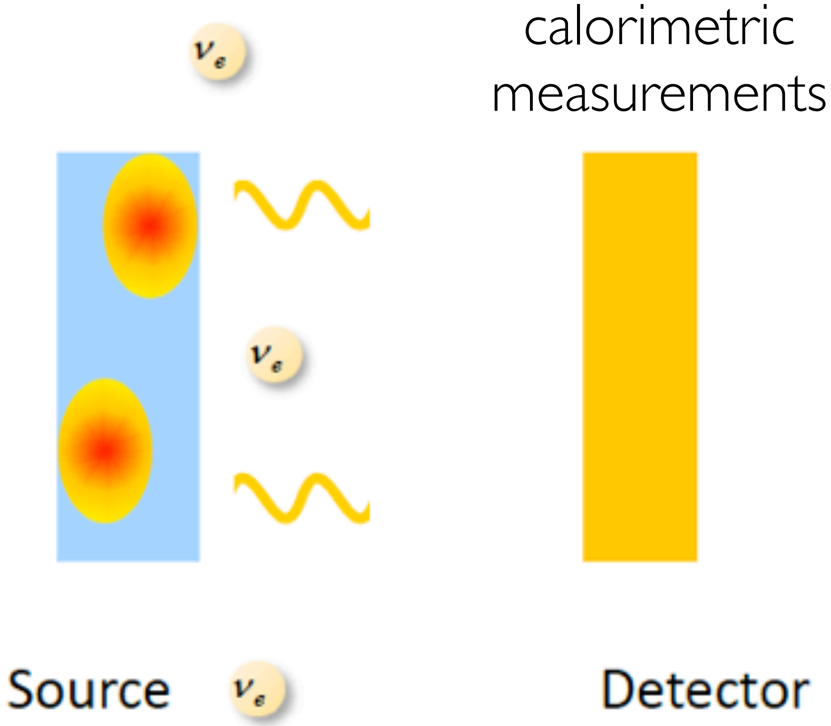
Atomic de-excitation:

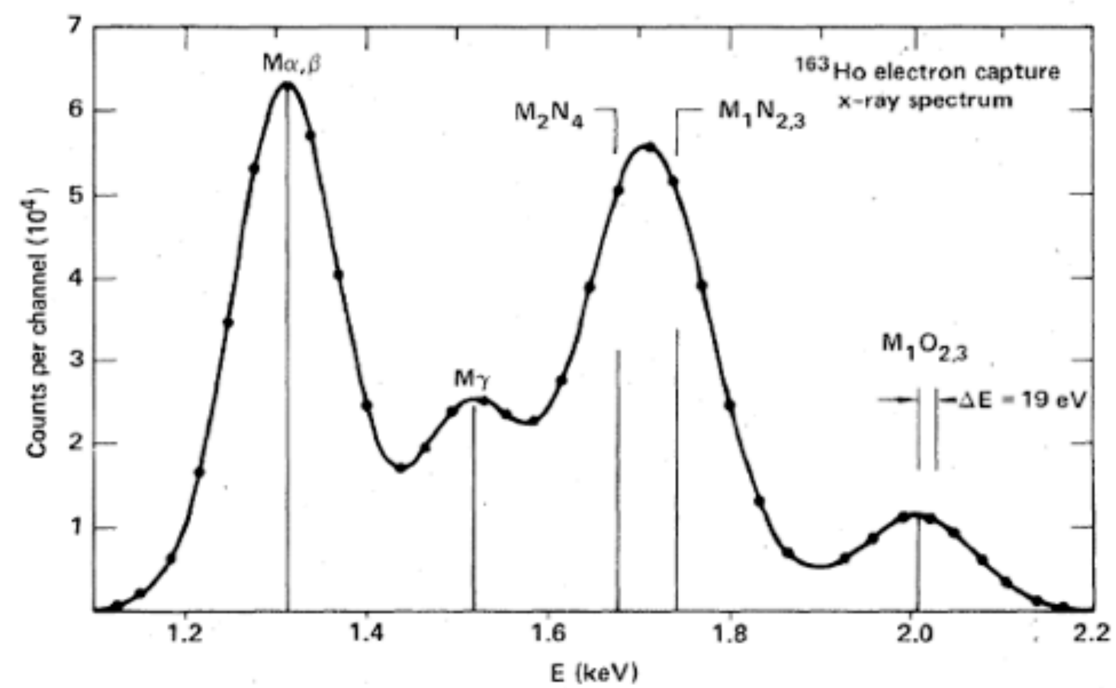
- **X-ray emission**
- Auger electrons
- Coster-Kronig transitions

$$^{163}_{67}\mathrm{Ho} \rightarrow ^{163}_{66}\mathrm{Dy}^* + \nu_e$$

$$^{163}_{66}\mathrm{Dy}^* \rightarrow ^{163}_{66}\mathrm{Dy} + E_C$$







Volume 118B, number 4, 5, 6 PHYSICS LETTERS 9 December 1982

CALORIMETRIC MEASUREMENTS OF ^{163}HOLMIUM DECAY AS TOOLS TO DETERMINE THE ELECTRON NEUTRINO MASS

A. DE RÚJULA and M. LUSIGNOLI [1]
CERN, Geneva, Switzerland

Atomic de-excitation:

- **X-ray emission**
- Auger electrons
- Coster-Kronig transitions

$$^{163}_{67}\mathrm{Ho} \rightarrow ^{163}_{66}\mathrm{Dy}^* + \nu_e$$

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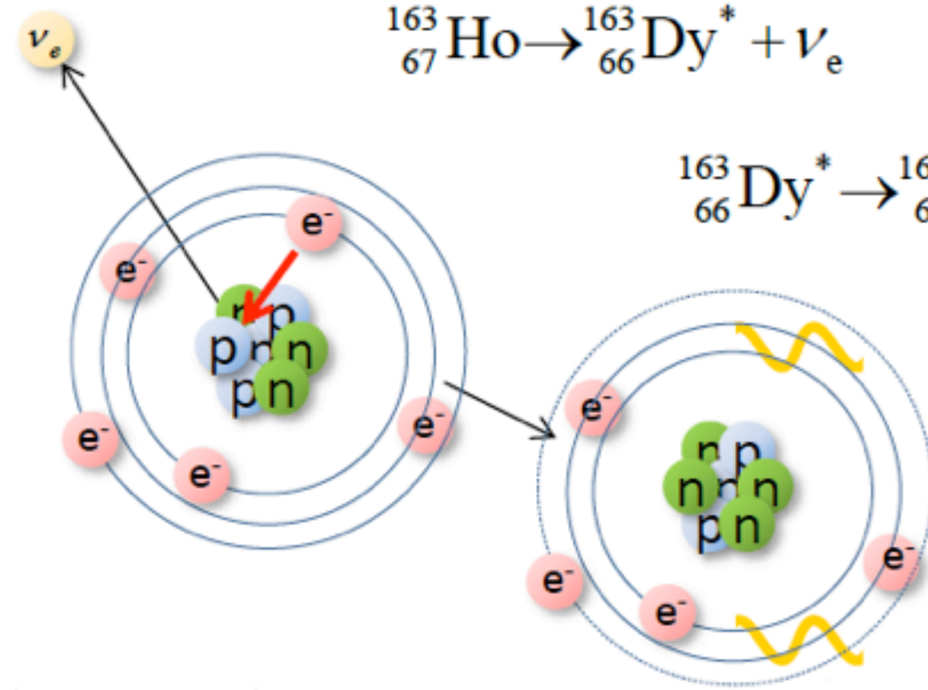
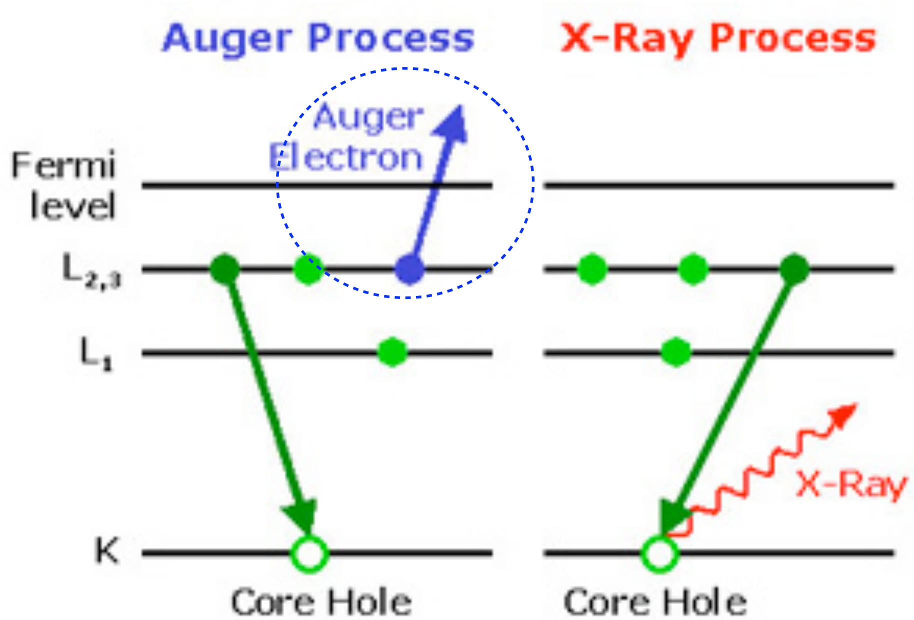
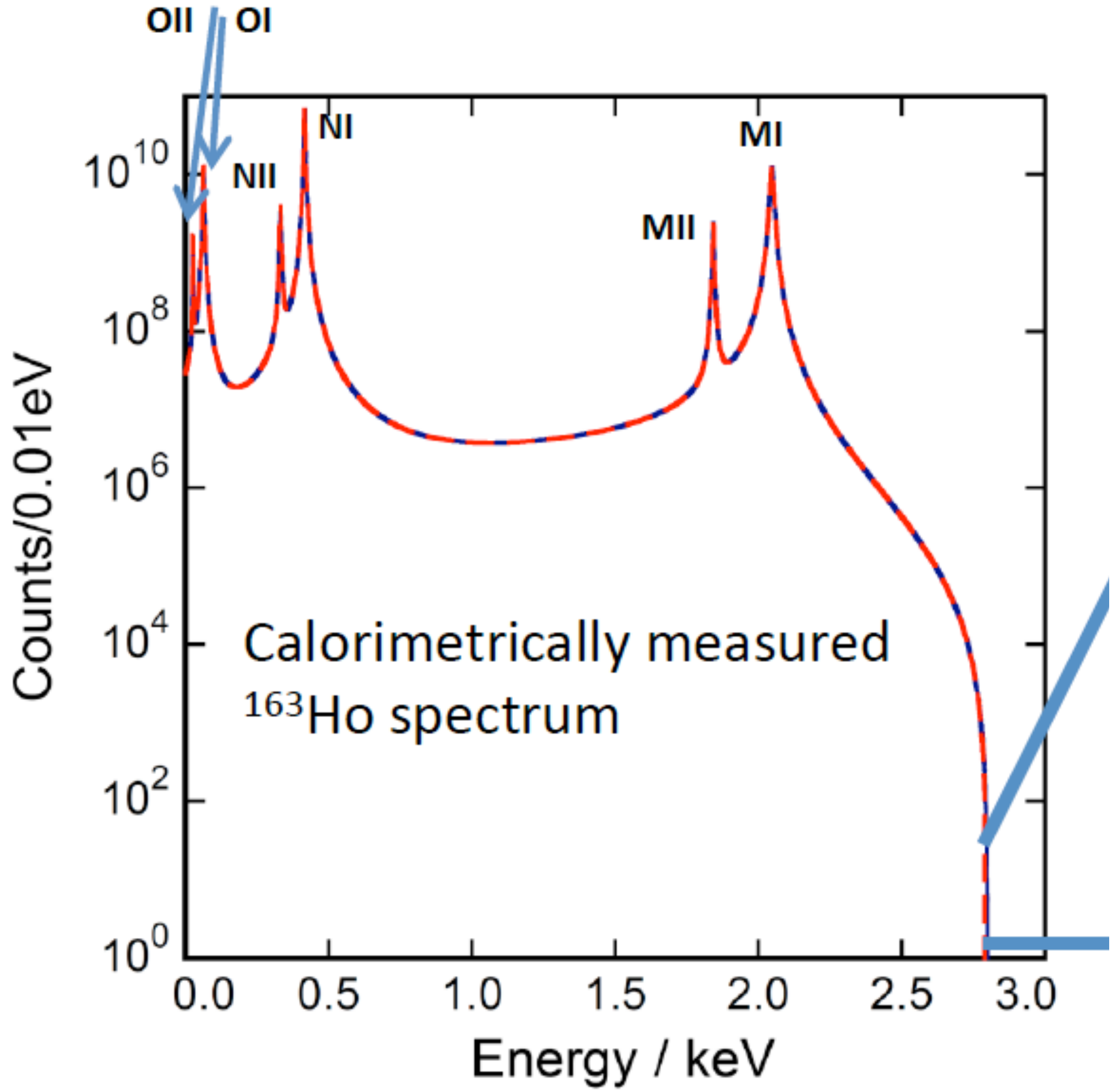


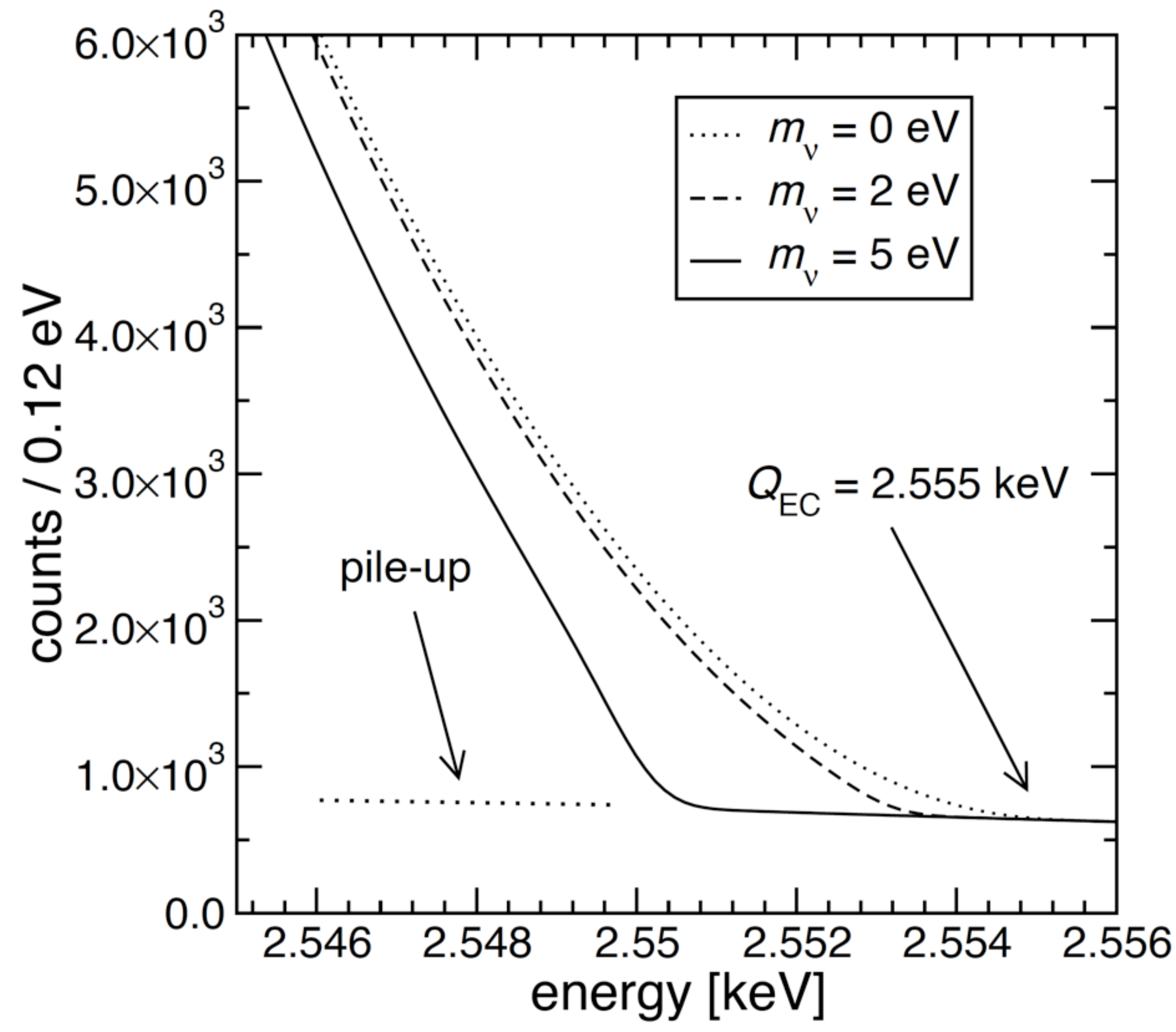


- $\tau_{1/2} \cong 4570$ years (2*10^{11} atoms for 1 Bq)
- Q_{EC} = $(2.833 \pm 0.030^{stat} \pm 0.015^{syst})$ keV

S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

- Q_{EC} function of m(ν) ⇒ measurement of E(e) sensitive at end-point
- One sees Auger electron and measures E(e) inclusively, with source = detector
  - no risk of loss or of mis-modelling energy at source
- E freed by de-exciting ^{163}Dy has ~lowest known "Q value": 2.8 keV
- *problem*: lifetime! need smart format of detector to maximize statistics



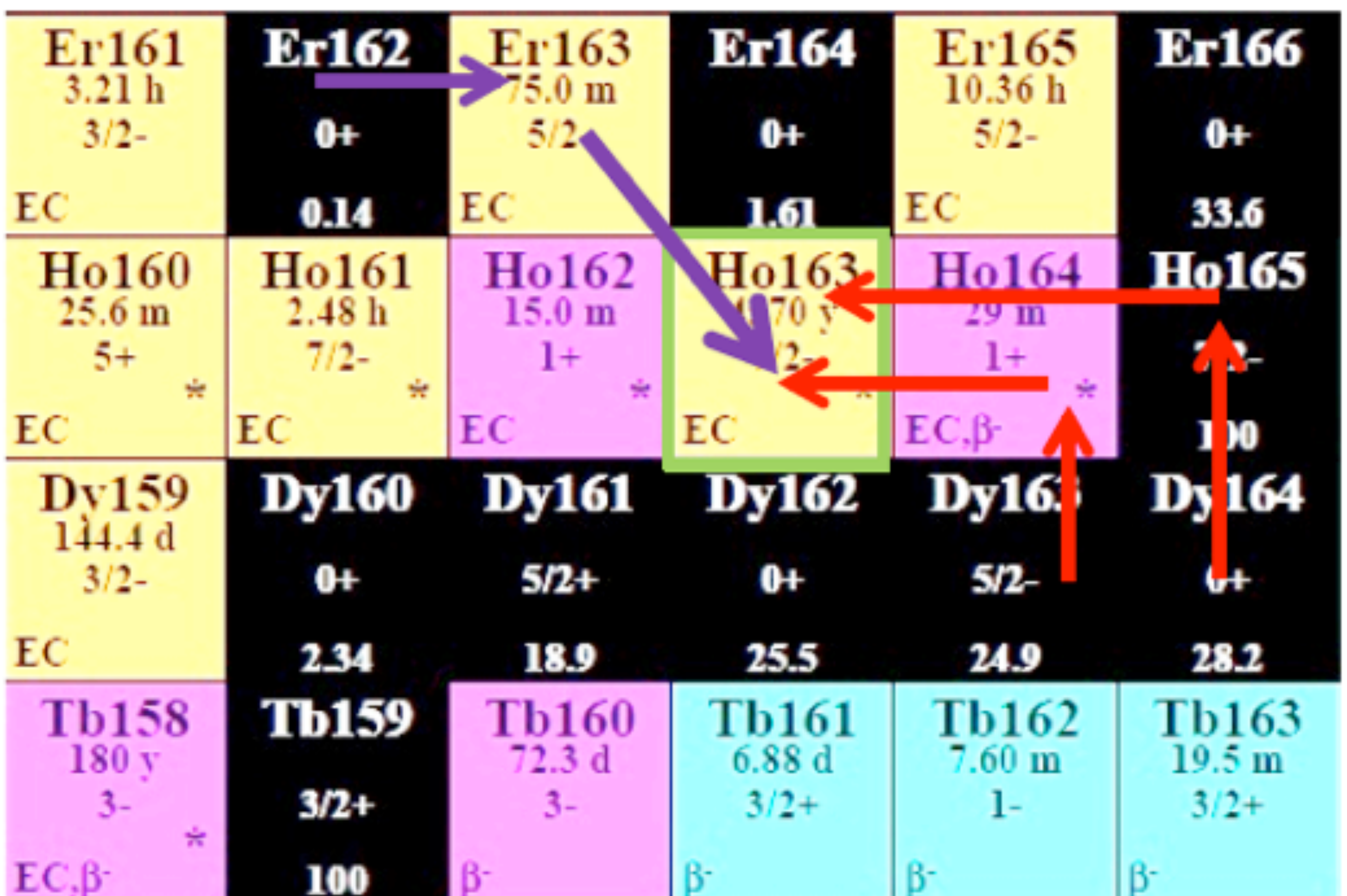


$$\frac{d\lambda_{EC}}{dE_c} = \frac{G_\beta^2}{4\pi^2}(Q - E_c)\sqrt{(Q - E_c)^2 - m_\nu^2} \times$$
$$\sum_i n_i C_i \beta_i^2 B_i \frac{\Gamma_i}{2\pi} \frac{1}{(E_c - E_i)^2 + \Gamma_i^2/4},$$

Equivalent of beta-decay but with B-W peaks corresponding to energy levels

Main issues

- Measuring the energy with micro-calorimeters with high resolution
- Estimate doubles=overlaps and their bias on Q end-point
- Keep background under control

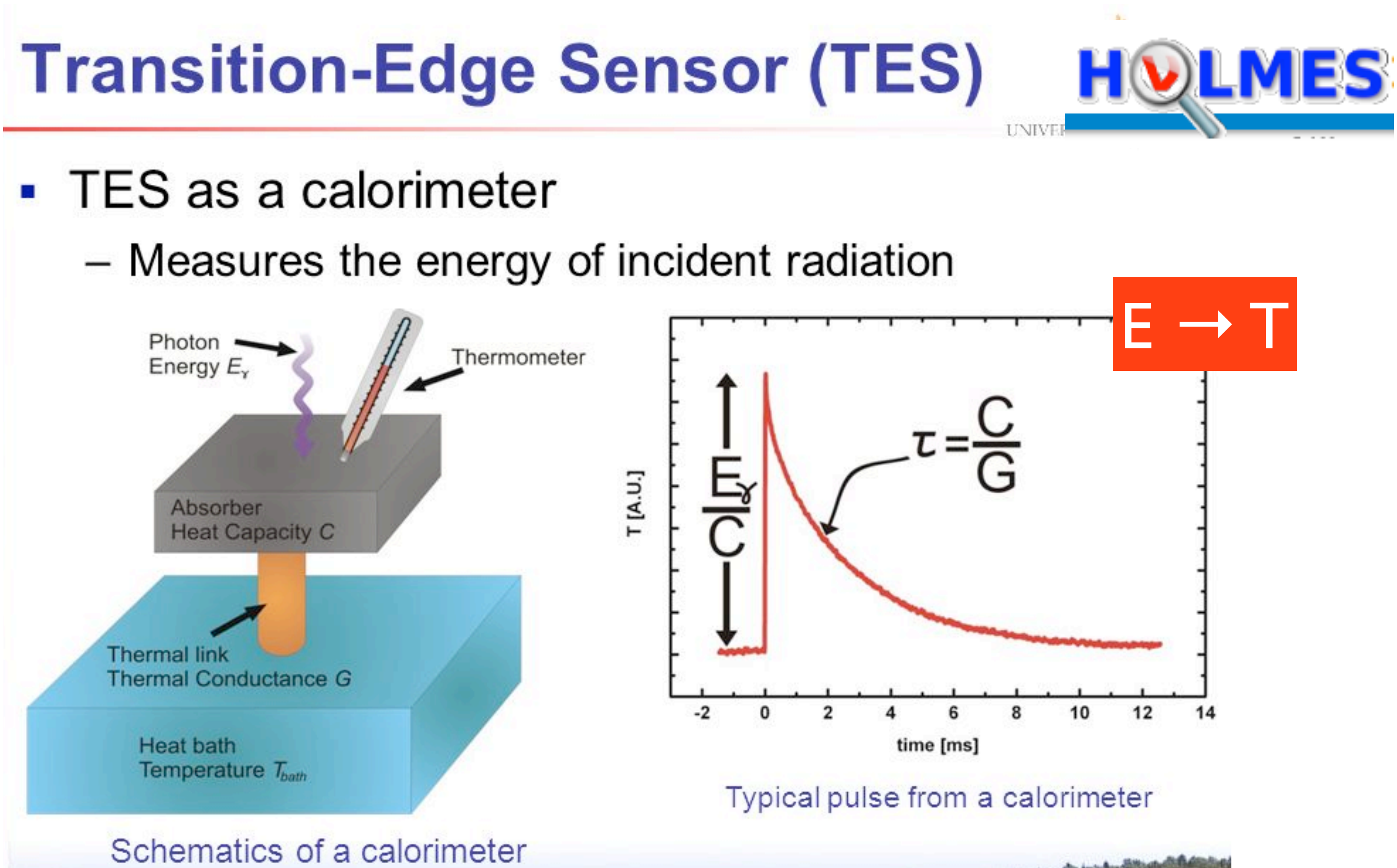


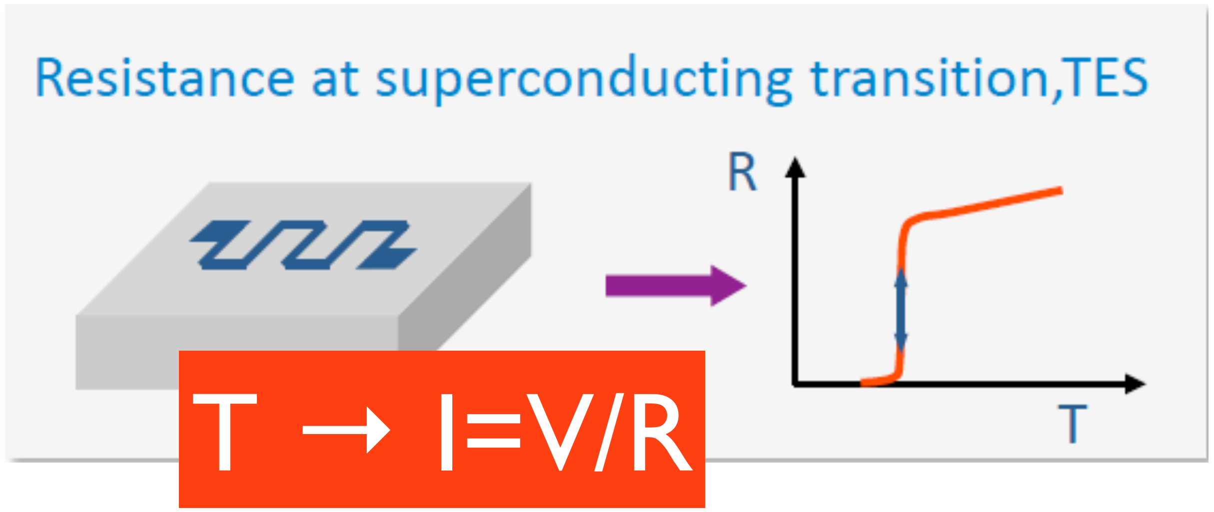
$$^{162}\mathrm{Er}(\mathrm{n},\gamma)^{163}\mathrm{Er} \rightarrow {}^{163}\mathrm{Ho} + \nu_e$$

- ^{163}Ho source mostly from neutron irradiation of ^{163}Er
- decays quickly (τ ~75 min) and large x-sec: effective process
- but mind radioactive impurities from other elements emitting below 5 keV!

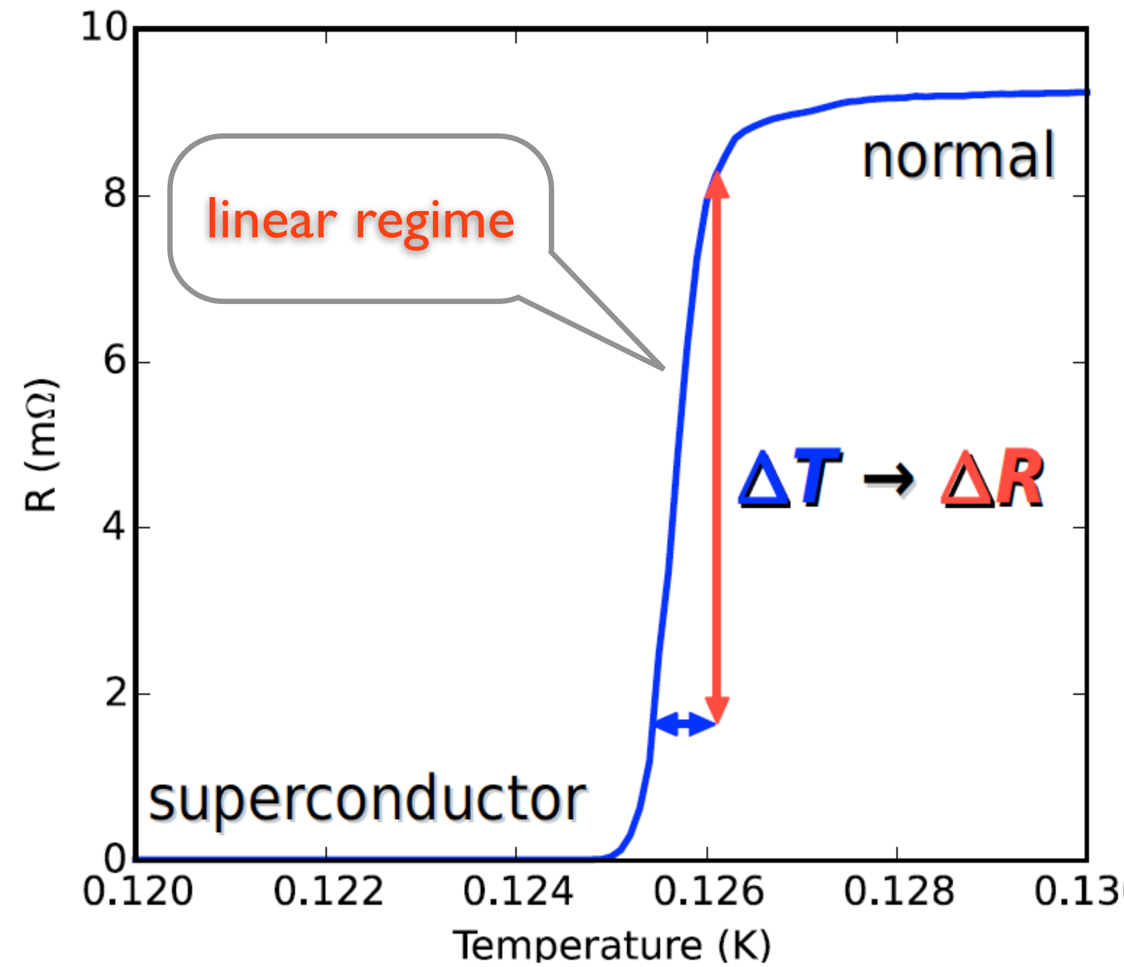
Microcalorimetry

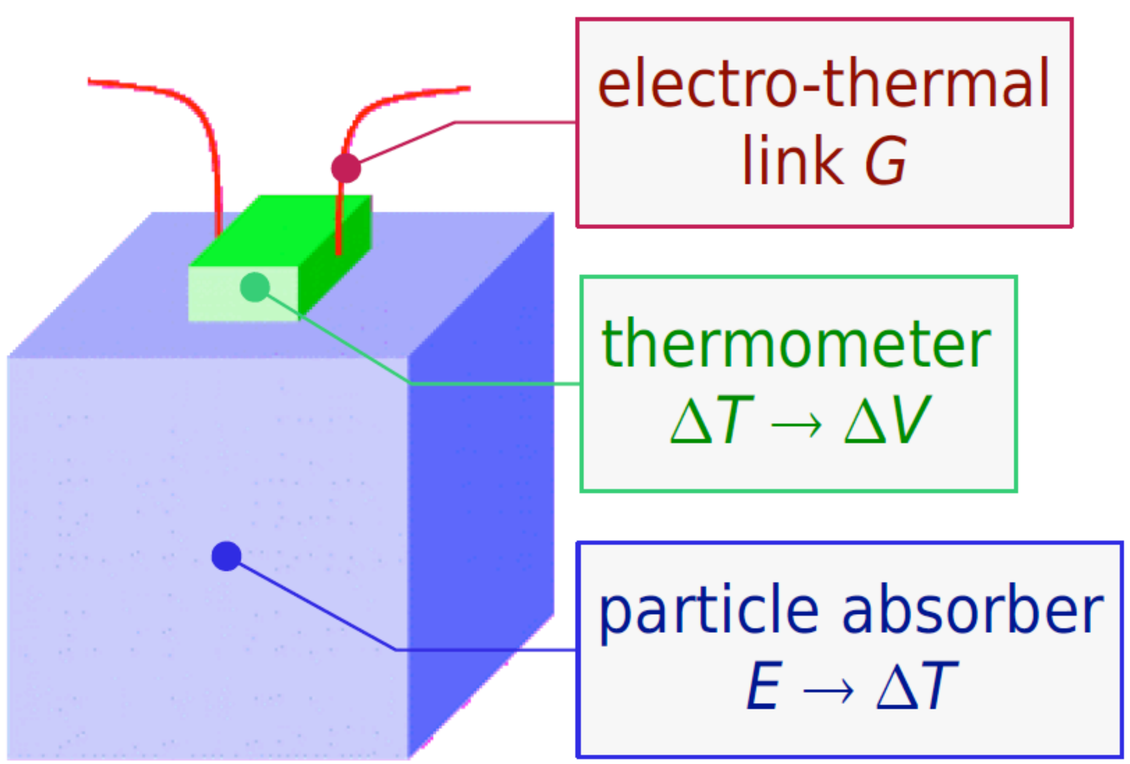
- Small, segmented devices needed for high E resolution (and to avoid two events overlapping in time in same reading element)





- Exploit super-conducting property to rapidly (but about linearly) change R with T

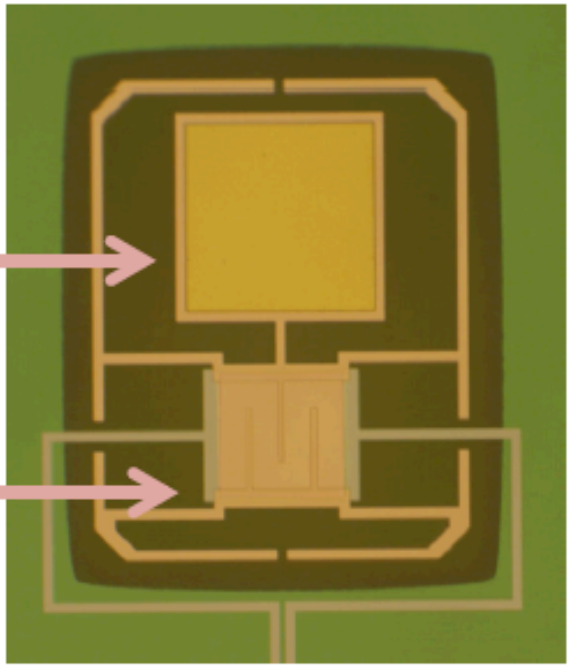




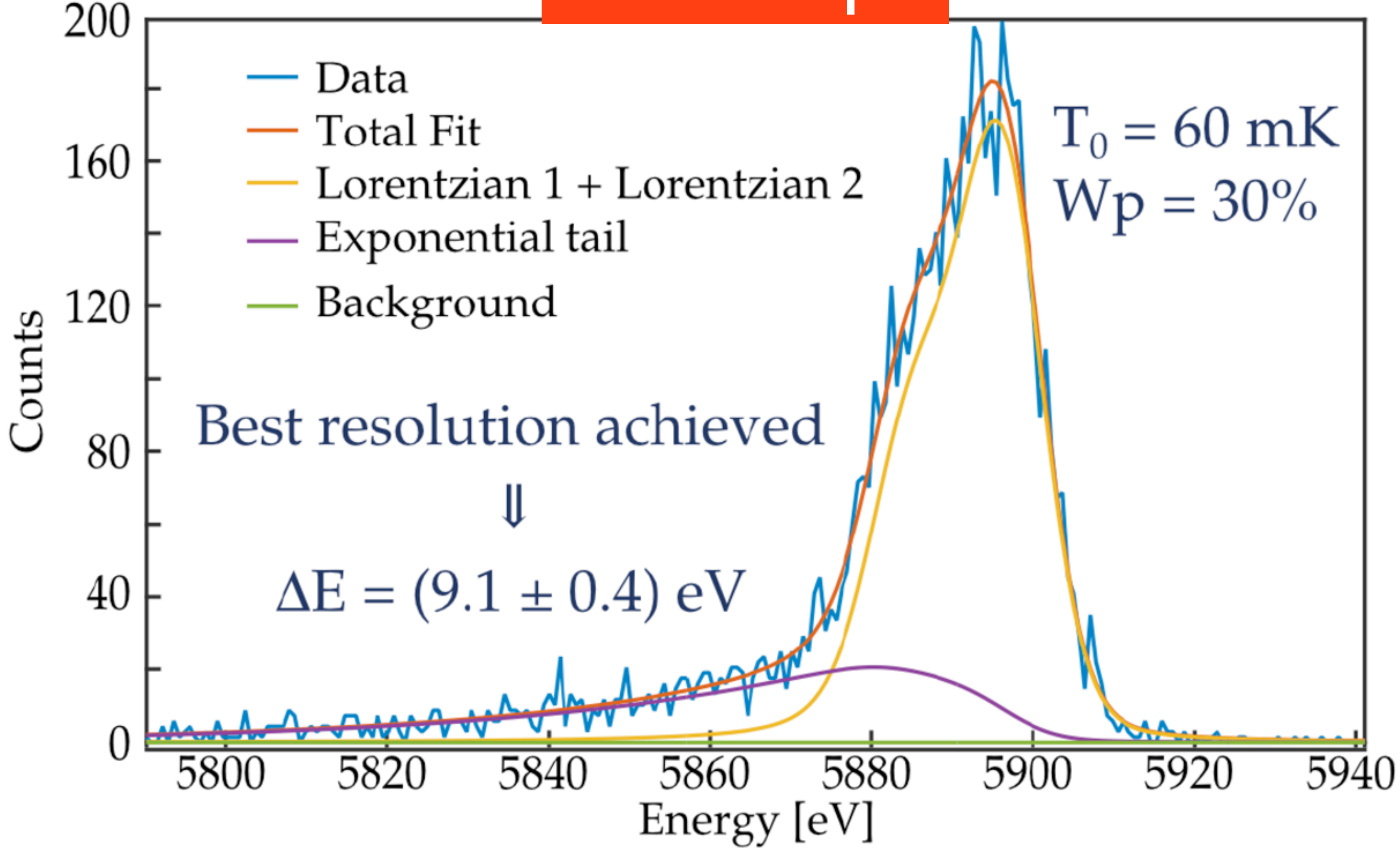
- Use film of material undergoing transition at roughly ^{163}Ho Q (2-3 keV)
    - linear with Ec ⟺T
    - generates current read-out
- Electro-thermic feedback compensates increase of I from T
    - forces TES to stay around linear regime

Absorber: Bi-Au or Au + 6.5x10^{13} ^{163}Ho per detector → **300 dec/sec**
^{163}Ho ion implanted in absorber using dedicated facility at Genoa University

Transition Edge Sensor: MoCu or MoAu superconducting films



An example



$$\Delta E \simeq 2.35\sqrt{4kTE_{\max}}$$

- Q = 2.5 keV
- operating at <150 mK heat capacity
- gives **Δ**E ~ 0.8 eV FWHM

- HOLMES claims **Δ**E = 1 eV

Pile-up

http://lanl.arxiv.org/pdf/1412.5060v3

- TES have a relaxation time of ~several ms
- Two decays can happen close in time within the same TES element and not be discriminated

➡ bias on the E_c measured by summing two processes

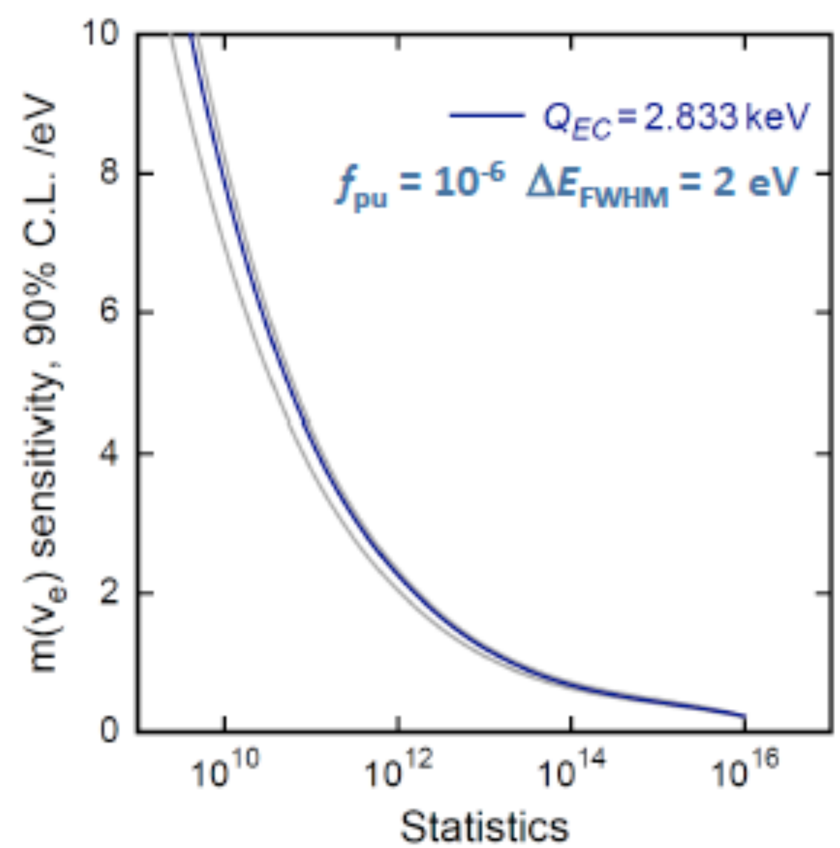
> spectrum is given by the two event pile-up probability $f_{pp} = \tau_R A_{EC}$, where τ_R is the time resolution and A_{EC} is the EC activity in each detector. This kind of statis-

Statistics in the end point region
- $N_{ev} > 10^{14}$ → **A ≈ 1 MBq**

Unresolved pile-up ($f_{pu} \sim a \cdot \tau_r$)
- $f_{pu} < 10^{-5}$ requirement
- $\tau_r < 1\ \mu s$ → $a \sim 10$ Bq
- 10^5 pixels

Precision characterization of the endpoint region
- $\Delta E_{FWHM} < 3$ eV

Background level
- $< 10^{-5}$ events/eV/det/day



Loredana Gastaldo
Kirchhoff Institute for Physics, Heidelberg University

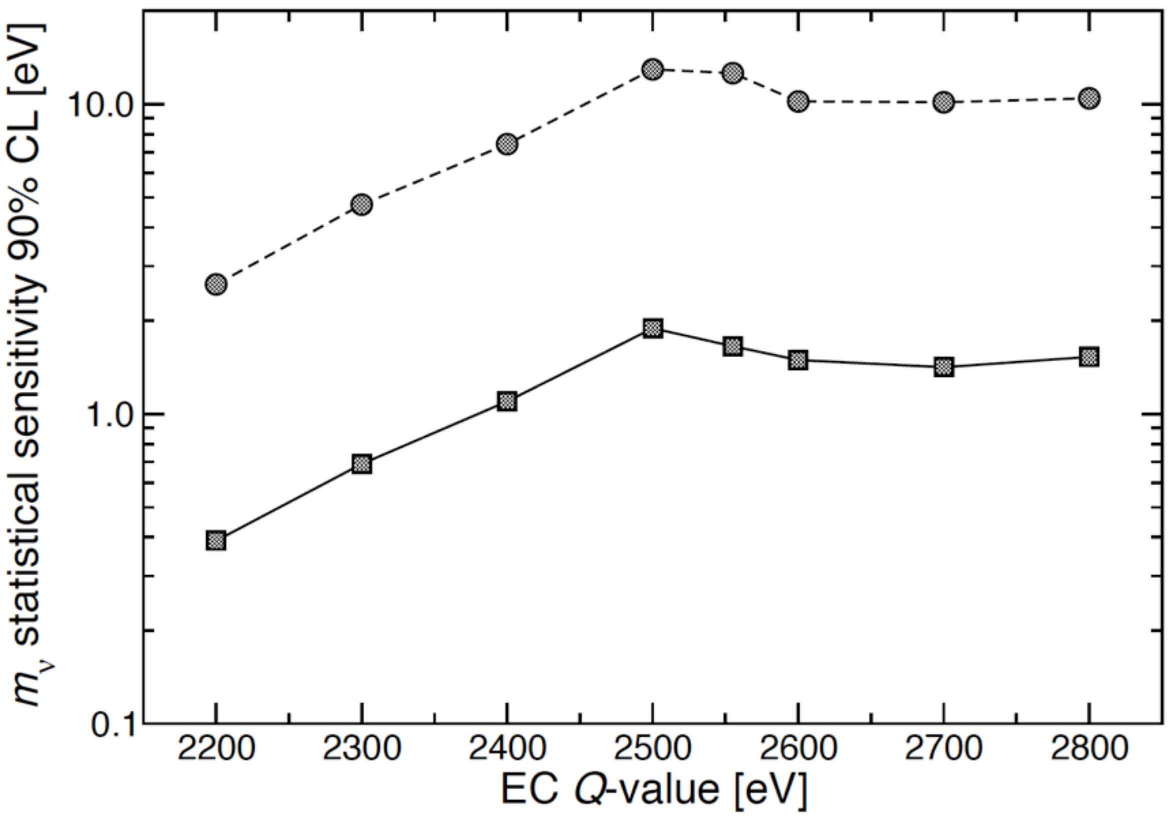


Fig. 4 Monte Carlo estimate of HOLMES neutrino mass statistical sensitivity for $N_{ev} = 3 \times 10^{13}$ (lower curve) or 10^{10} (upper curve) and with $f_{pp} = 3 \times 10^{-4}$, $\Delta E_{\mathrm{FWHM}} = 1\,\mathrm{eV}$, and no background.

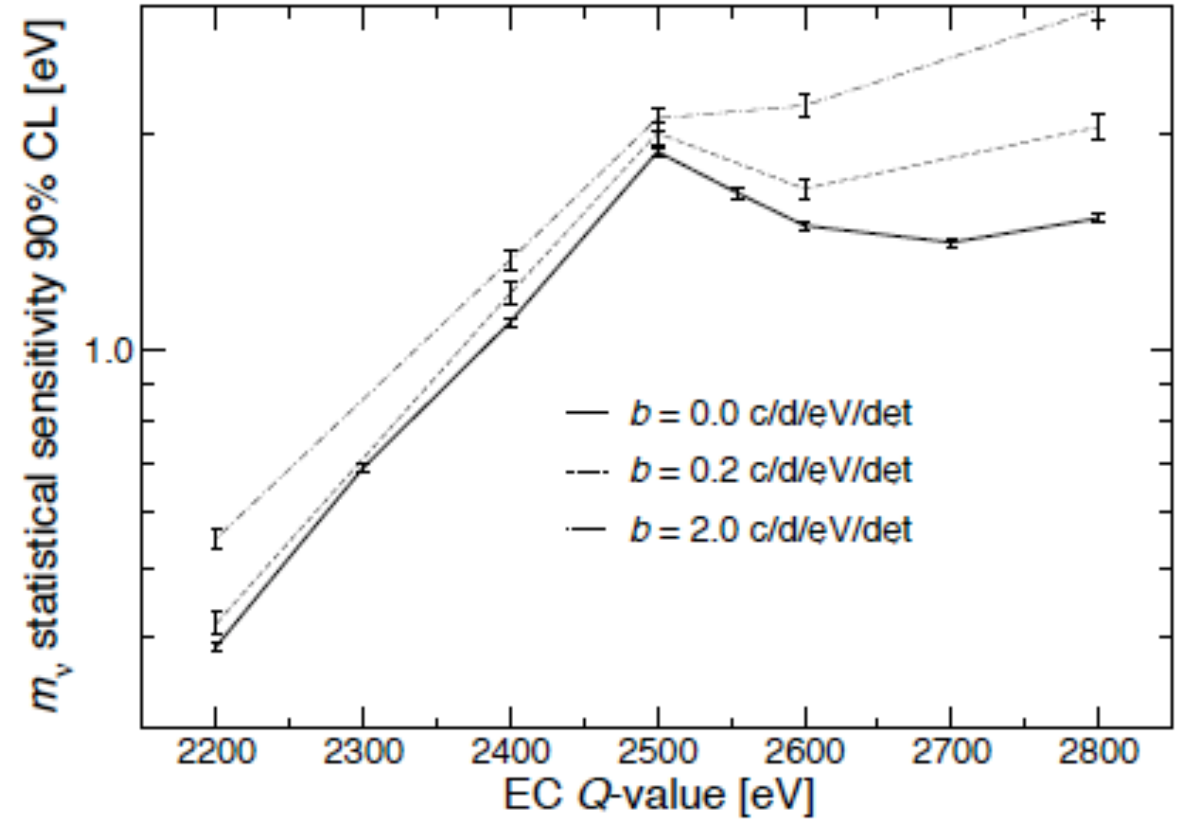


Fig. 9 Monte Carlo estimates of the effect of various background levels on HOLMES baseline statistical sensitivity.

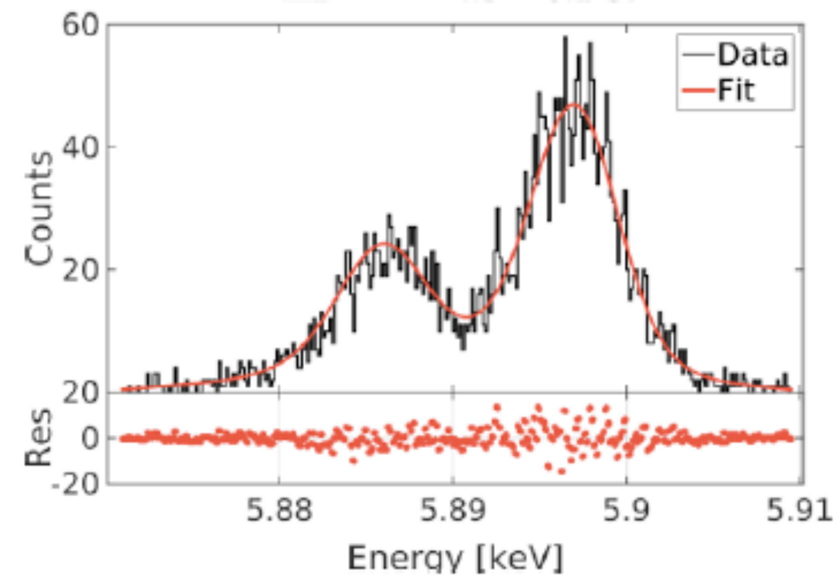
Effect of statistics

Effect of bkg

Overall this technique expected to currently attain
~0.2-0.4 eV neutrino mass sensitivity
⇒ competitive with spectrometers

Source production and purification:

130 MBq available for tests and experiments ✓

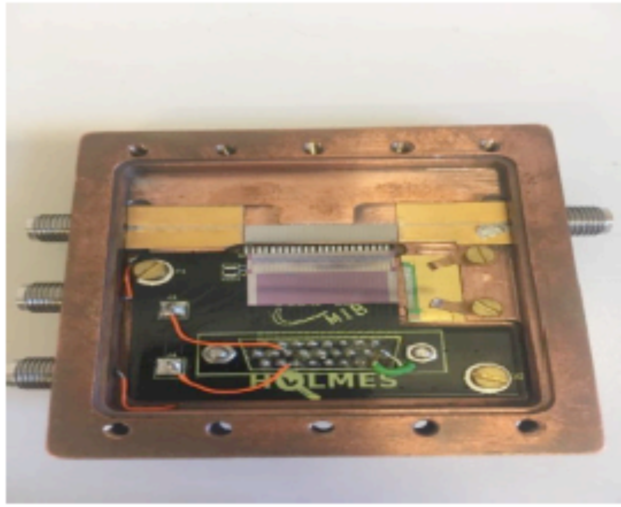


Detector arrays characterization:

very good single pixel performance ✓

operating microwave SQUID multiplexing

next challenge → load TES arrays with ^{163}Ho !



Dedicated mass separator:

facility installed

tests of the ion source on-going

commissioning on-going !

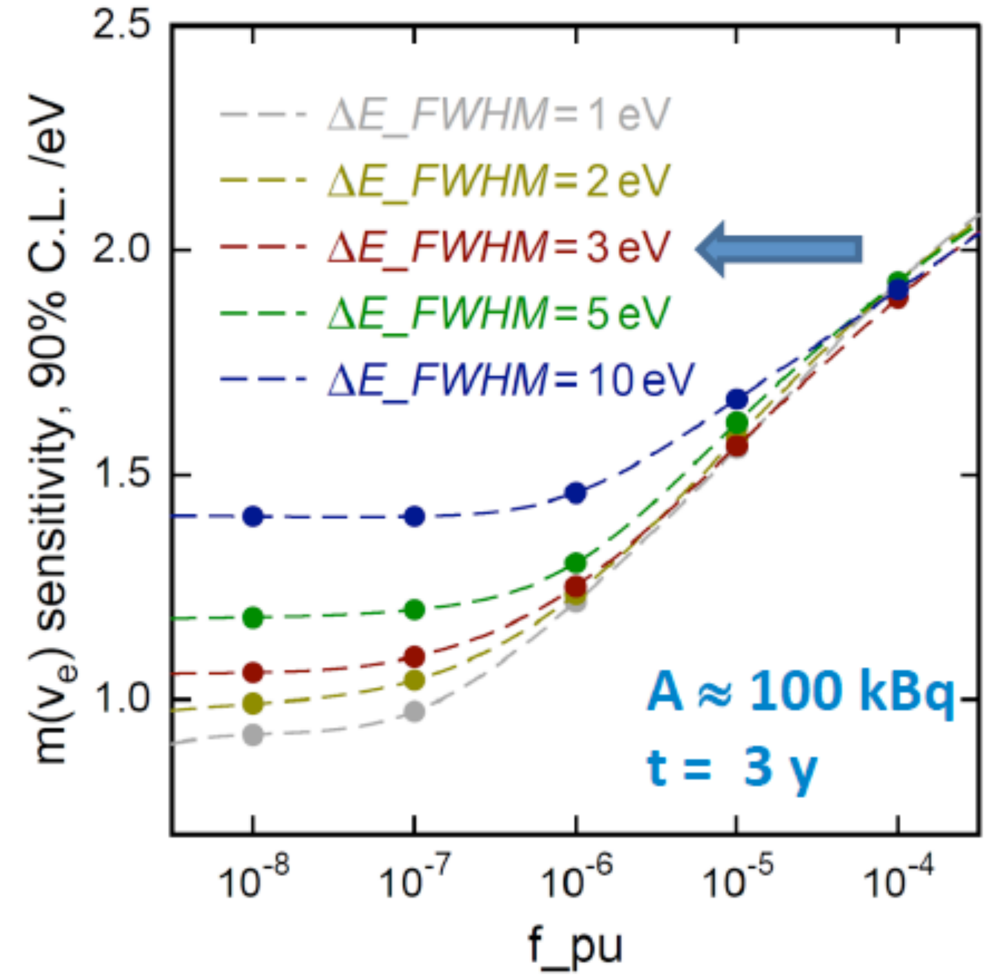
- full scale (starting 2019):
  1000 channels,
  t_M= 3 years (3×10^{13} events)
  → $m(\nu_e) < 1$ eV

Alternative observable: the ECHo experiment

Magnetization of paramagnetic material, MMC

A.Fleischmann, C. Enss and G. M. Seidel,
Topics in Applied Physics 99 (2005) 63

ECHo-100k (2018 – 2021)

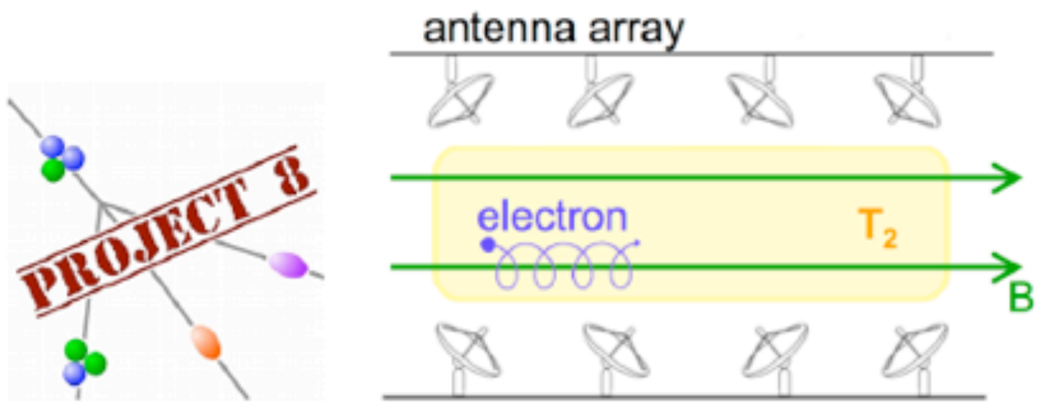


Activity per pixel 10 Bq
Number of detectors 12000
Readout: microwave SQUID multiplexing

Broadly same challenges and figures of merit to attain sensitivity

PROJECT 8

Alternative to KATRIN/ spectroscopy



- Enclosed volume filled with **tritium molecular gas**
- Add a **magnetic field →**
  Decay electrons spiral around field lines
- Add **antennas** to detect the cyclotron radiation

Cyclotron Radiation Emission Spectroscopy (CRES)

- Non-destructive measurement of electron energy

$$\omega_\gamma = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$

@ 1 Tesla	ω(18 keV) ~ 26 GHz P(18 keV) = 1.2 fW

$$Q = M_{^3\mathrm{H}} - M_{^3\mathrm{He}} - m_e = 18.58\,\mathrm{keV}$$

B. Monreal & J. Formaggio PRD 80 (2009) 051301

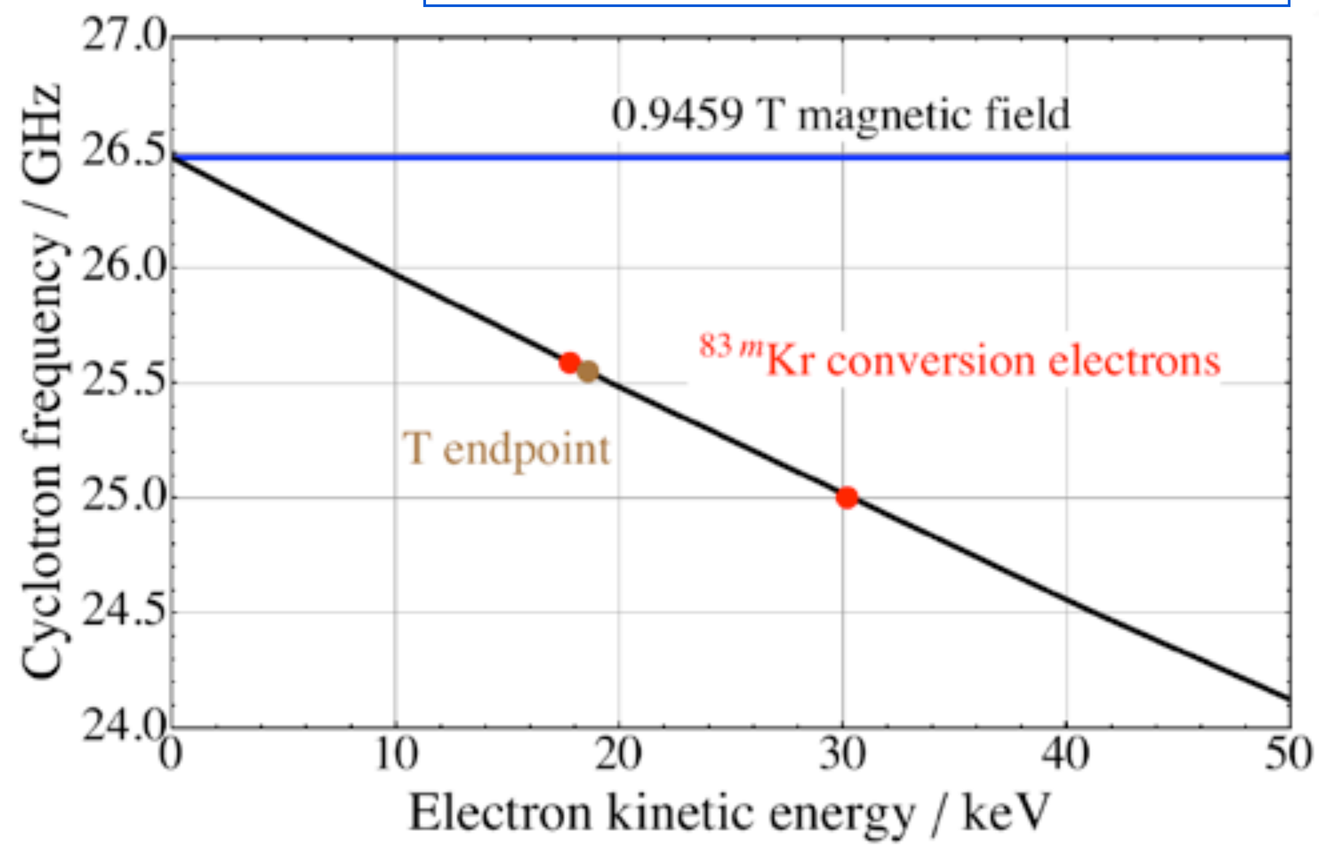
Cyclotron Radiation Emission Spectroscopy (CRES)

$Q = M_{^3\mathrm{H}} - M_{^3\mathrm{He}} - m_e = 18.58\,\mathrm{keV}$

The electron cyclotron frequency is related to the electron kinetic energy

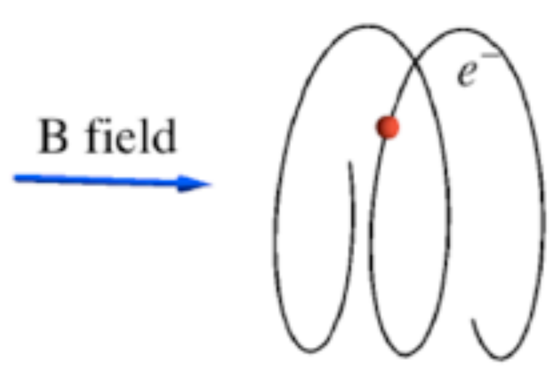
The cyclotron frequency is encoded in cyclotron radiation

Frequency is something we can measure very precisely – eV resolution demonstrated, sub-eV resolution expected



Cyclotron motion:

$$f_\gamma = \frac{f_c}{\gamma} B = \frac{1}{2\pi} \frac{eB}{m_e + E_{kin}/c^2}$$



$f_c = 27\,992.491\,10(6)\,\mathrm{MHz\,T^{-1}}$

Project 8 Goals

- Demonstrate that CRES can be used to measure the tritium endpoint in a small prototype

  (right now, "Phase II")

- Scale to a large-volume system that has sufficient statistics to contribute to the global neutrino mass effort and serve as a intermediate step for an atomic experiment

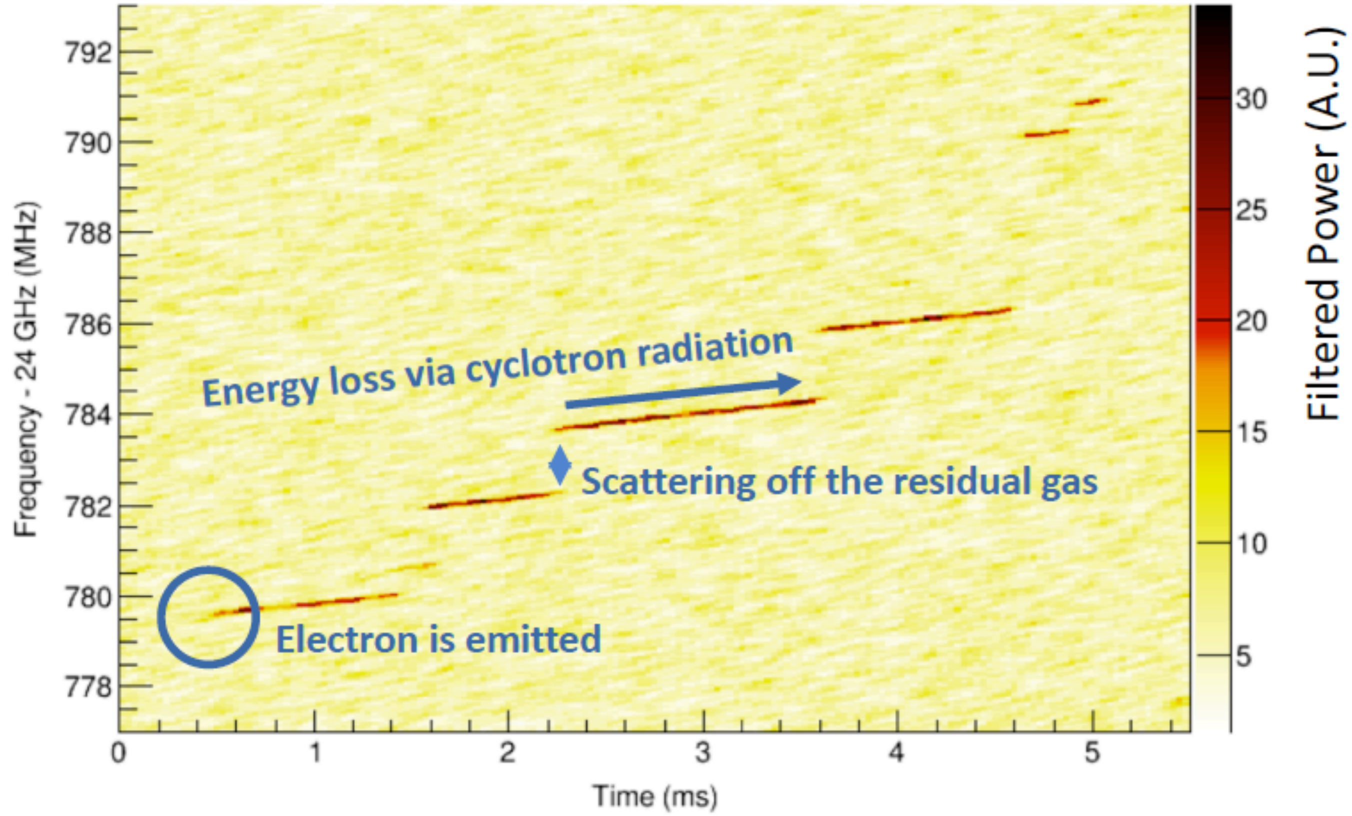
  (near future, "Phase III")

- Transition to an atomic tritium measurement and make the most sensitive measurement of the neutrino mass possible

  (future, "Phase IV")

Project 8 Electron Event with Energy 18 keV

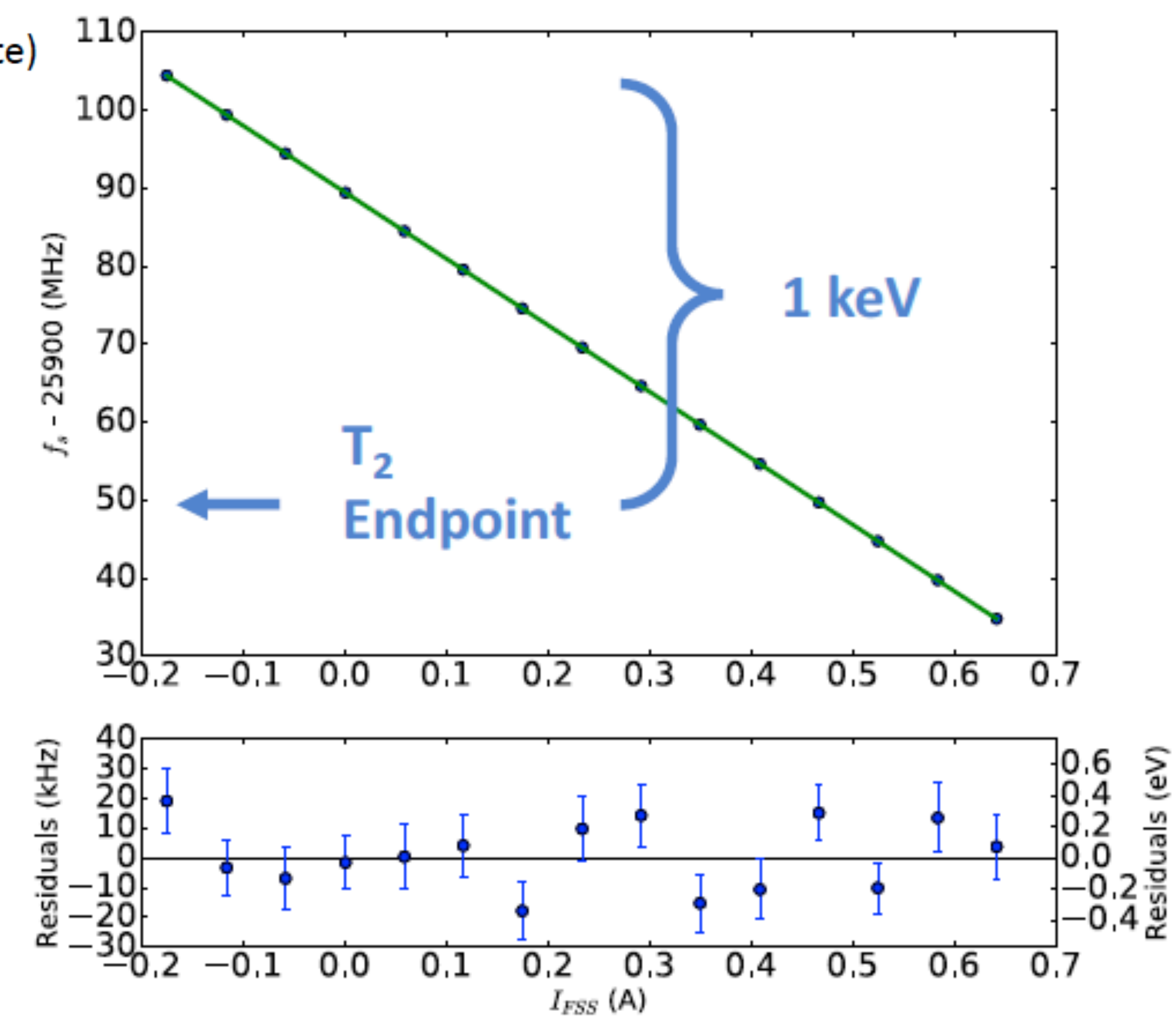
https://zenodo.org/record/1286954#.XEiHGvx7mi4



Frequency increases as energy is lost due to radiation (continuous) and collisions (discrete)

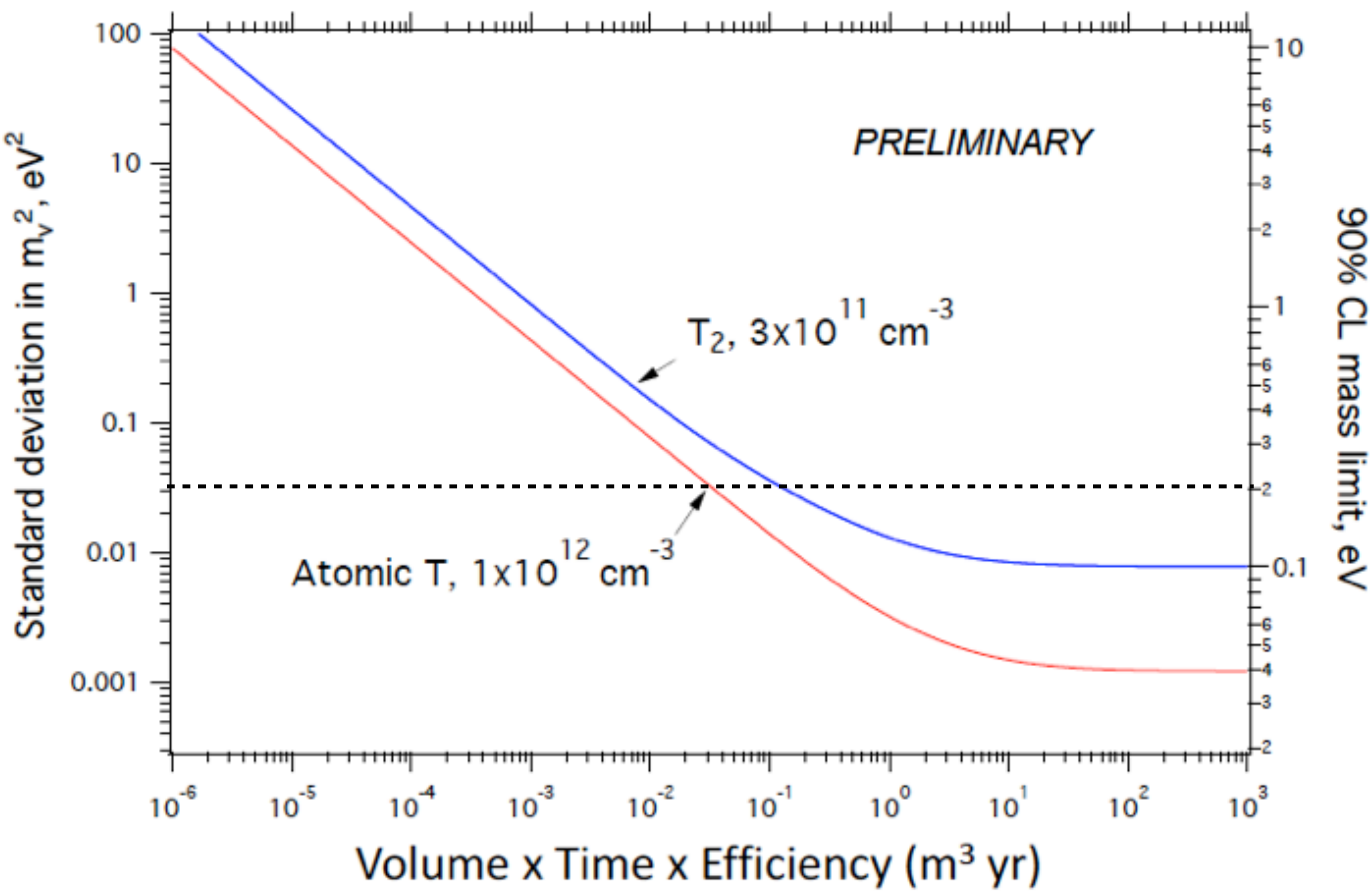
Cyclotron motion:

$$f_\gamma = \frac{f_c}{\gamma} = \frac{1}{2\pi} \frac{eB}{m_e + E_{kin}/c^2}$$



Linearity test completed, good to 0.2 eV
Efficiency data being taken right now

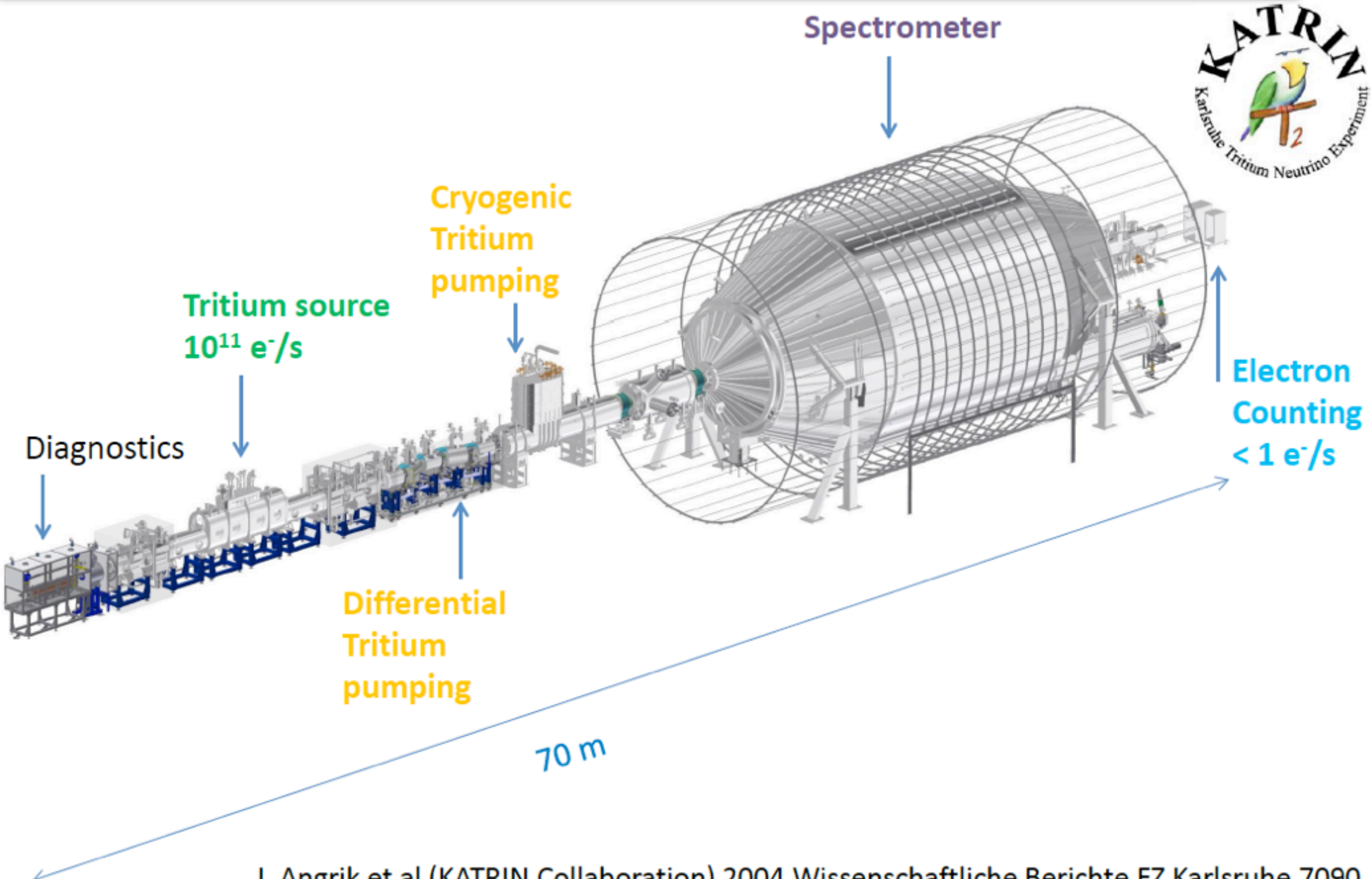
Project 8 Projected Capabilities with Atomic Tritium



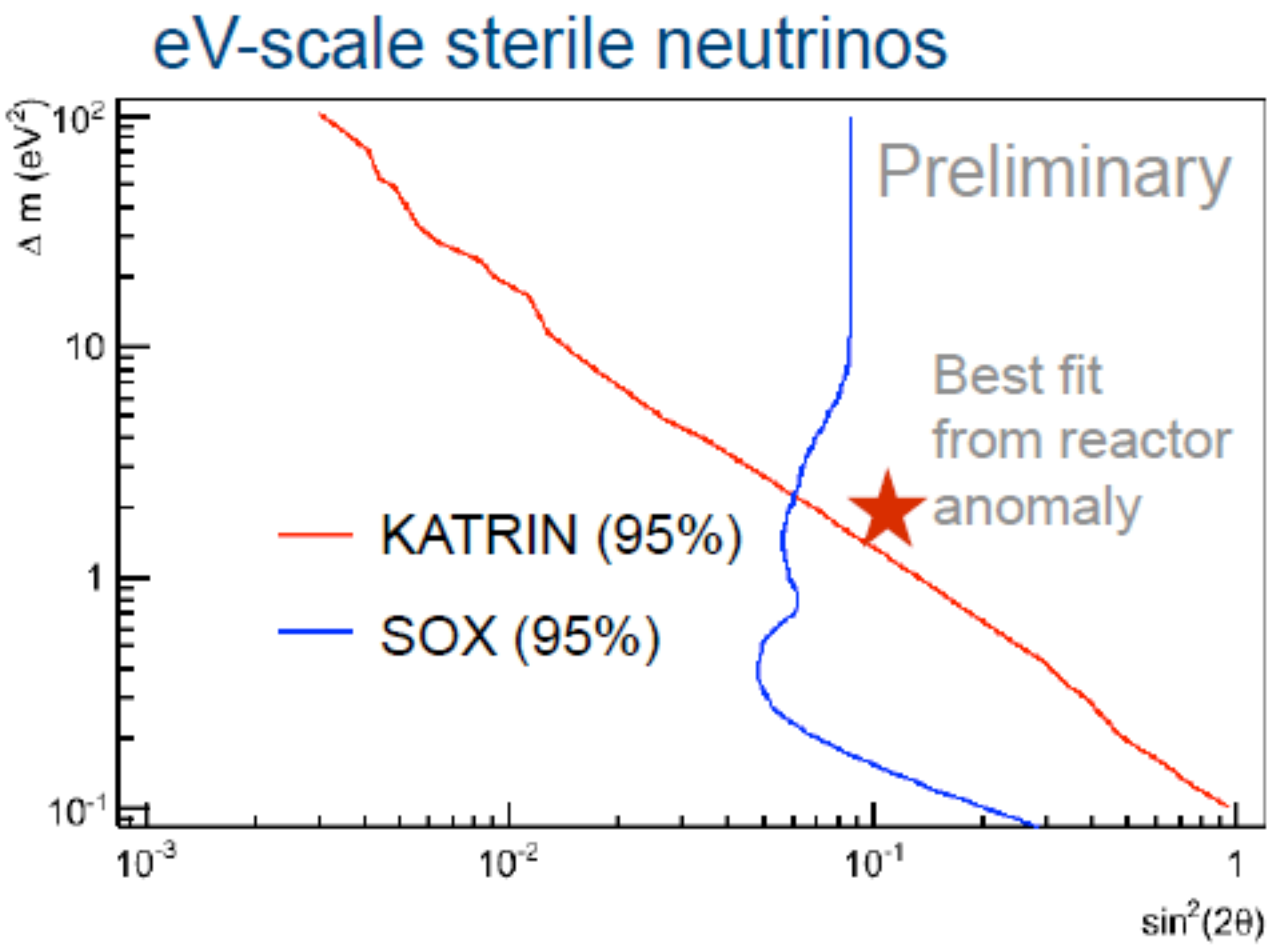
Atomic tritium should allow us sensitivity to a 40 meV mass scale with 10-100 m^3-years exposure

Back up

The KATRIN experiment



J. Angrik et al (KATRIN Collaboration) 2004 Wissenschaftliche Berichte FZ Karlsruhe 7090

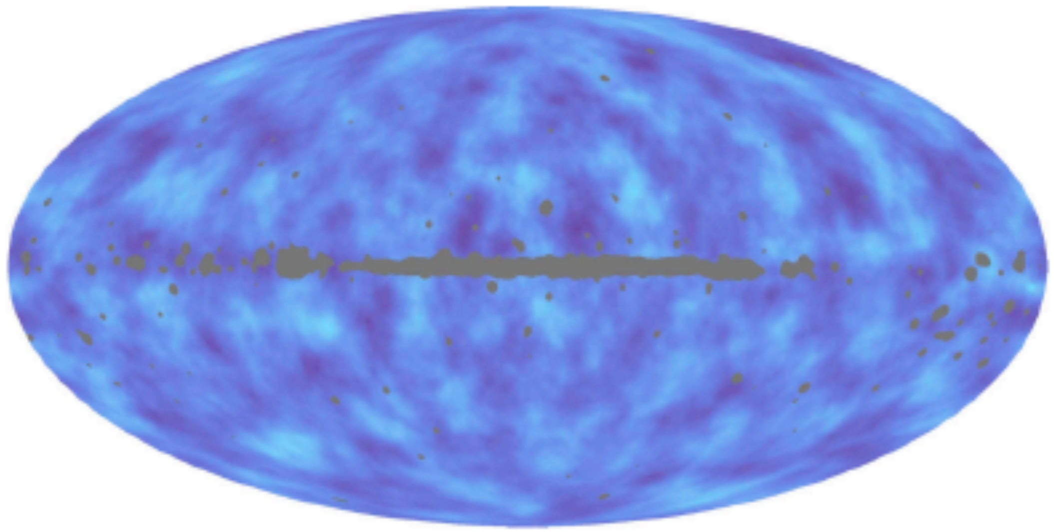


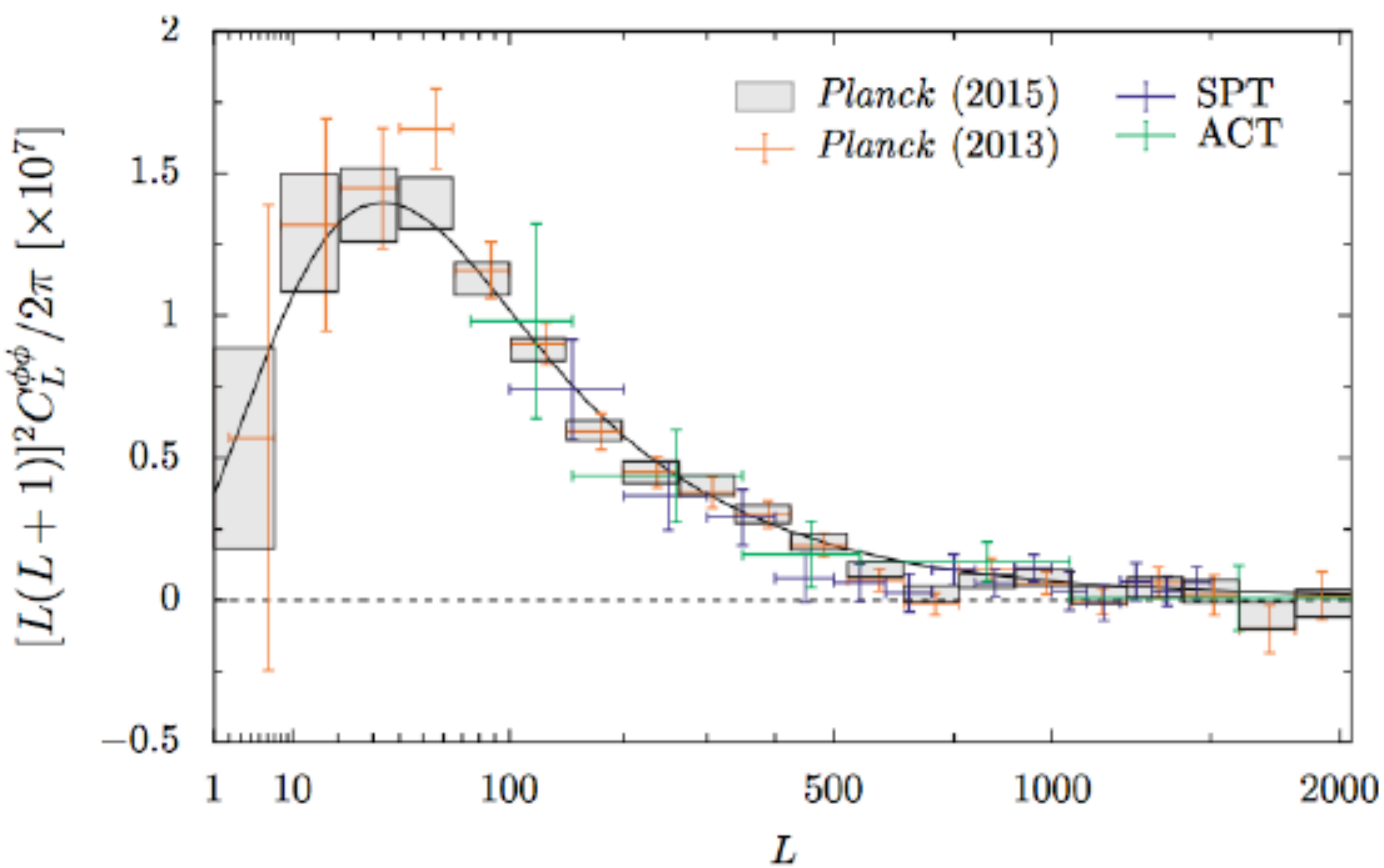
KATRIN **as is** probes the favored parameter space for light sterile neutrinos and is complementary to oscillation experiments

SOX at Borexino

first proposed (AFAIK...): Phys.Rev. D73 (2006) 045021

CMB lensing potential





LARGE SCALE STRUCTURES

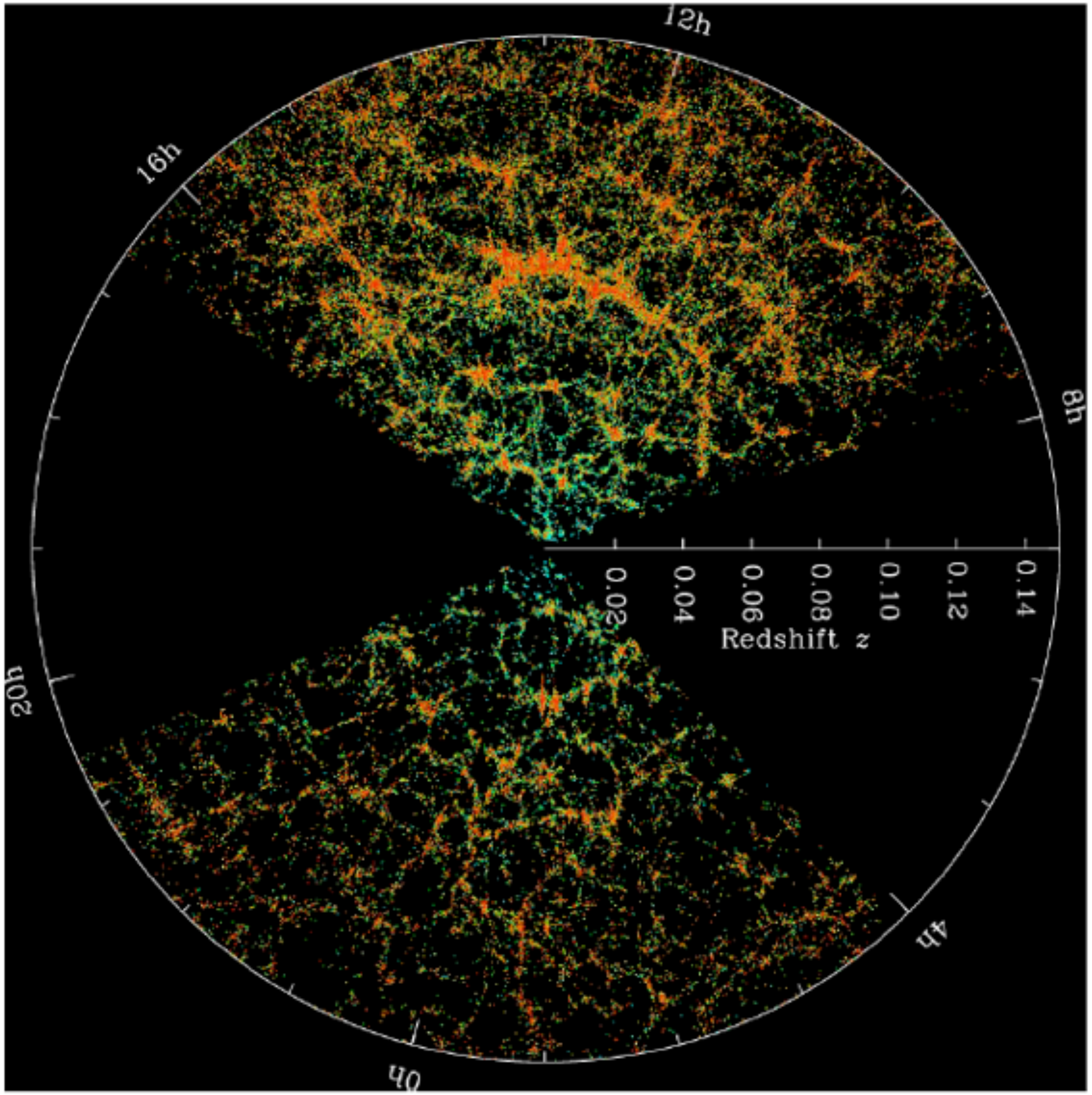
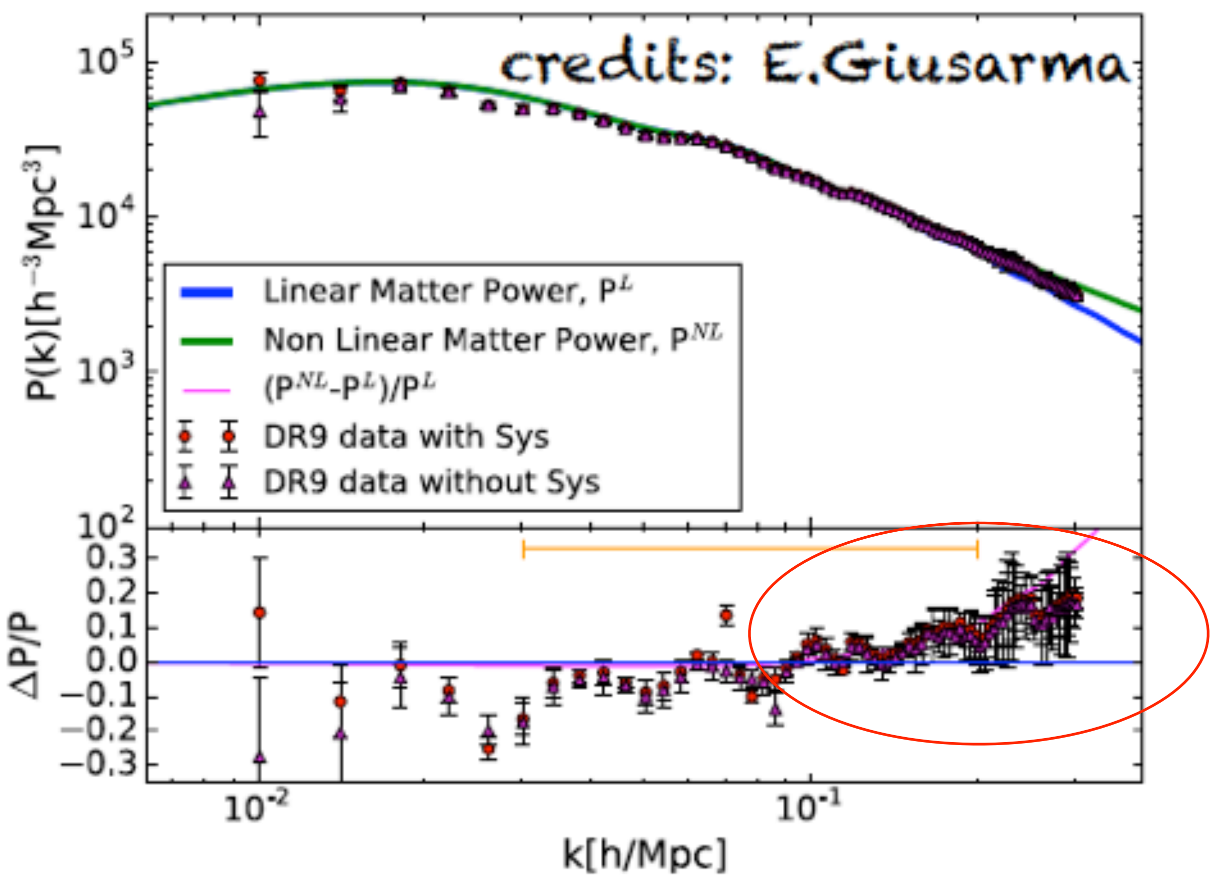
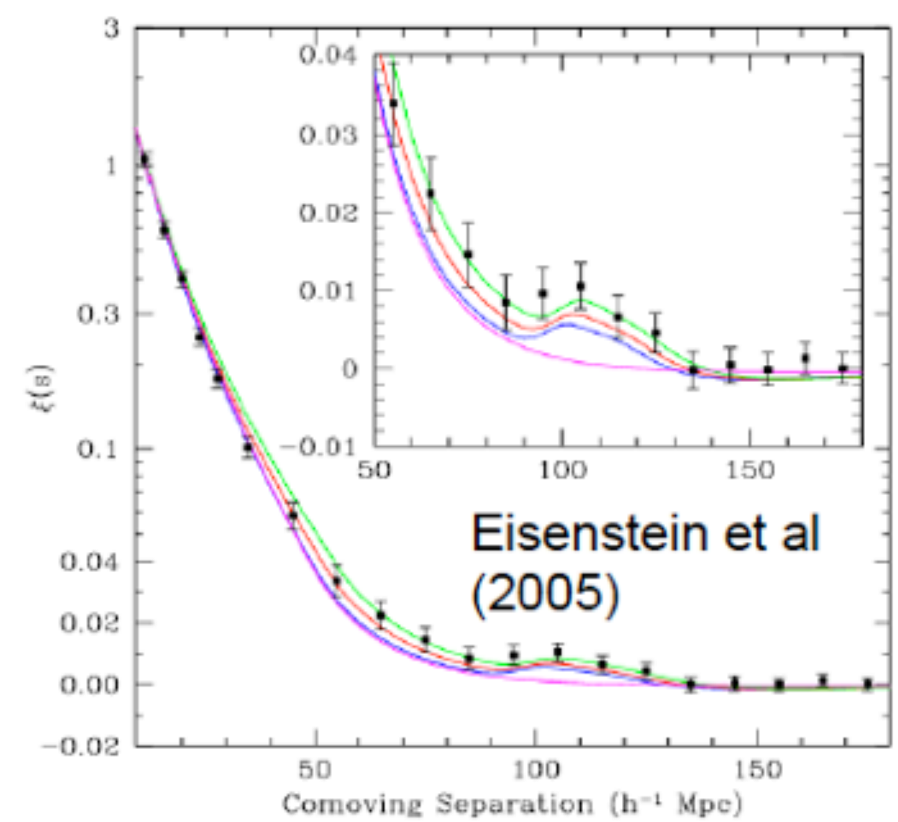


Image Credit: M. Blanton and the Sloan Digital Sky Survey.



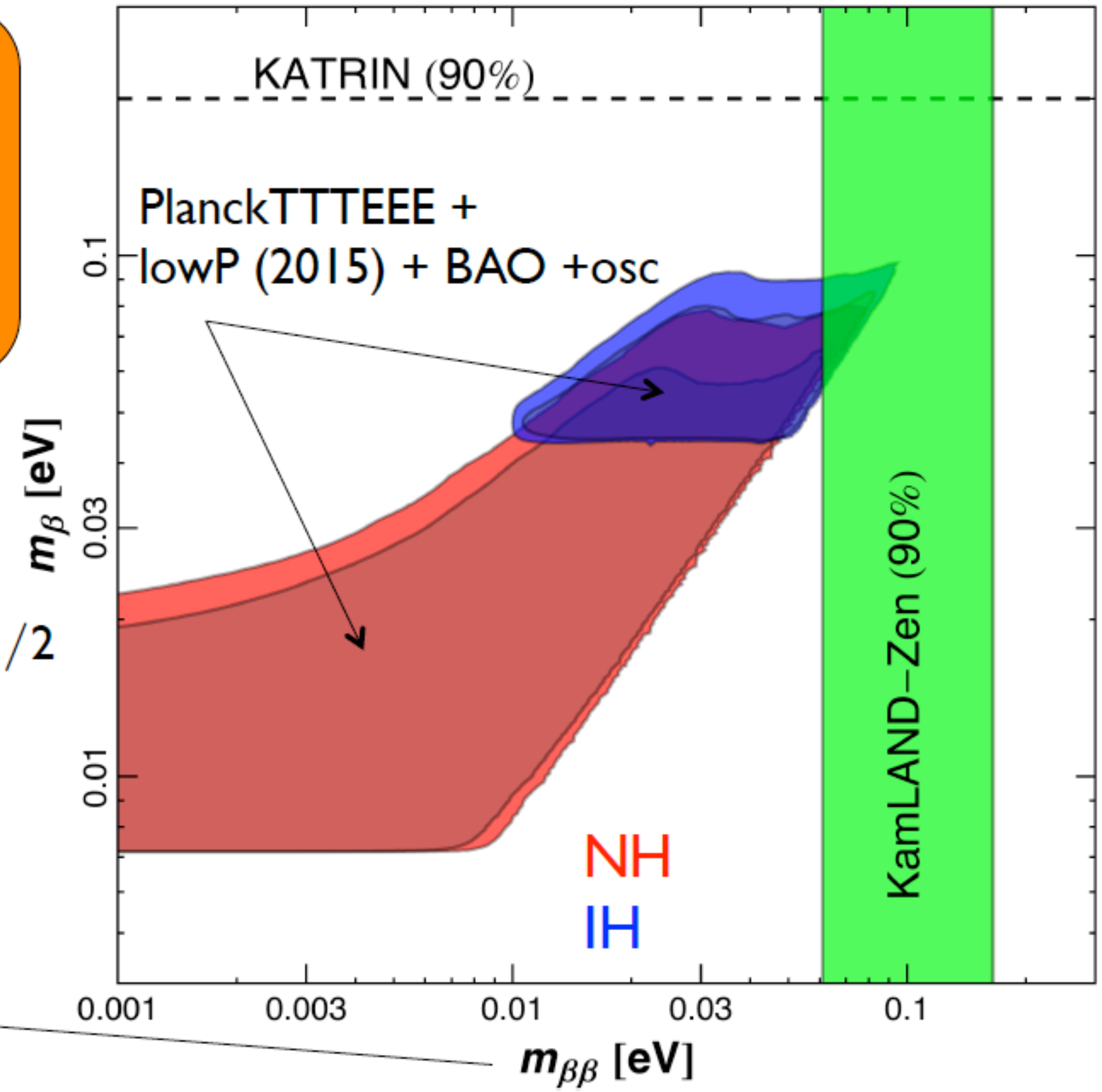


95% constraints on total mass	*PlanckTT*	*PlanckTTTEEE*
+lowP	<0.72 eV	<0.49 eV
+lowP+lensing	<0.68 eV	<0.59 eV
+lowP+BAO	<0.21 eV	<0.17 eV
+lowP+ext	<0.20 eV	<0.15 eV
+lowP+lensing+ext	<0.23 eV	<0.19 eV

Cosmology constraints can be combined with data from oscillation experiments

$$m_\beta \equiv \left[\sum |U_{ei}|^2 \, m_i^2\right]^{1/2}$$

$$m_{\beta\beta} \equiv \left|\sum U_{ei}^2 m_i\right|$$

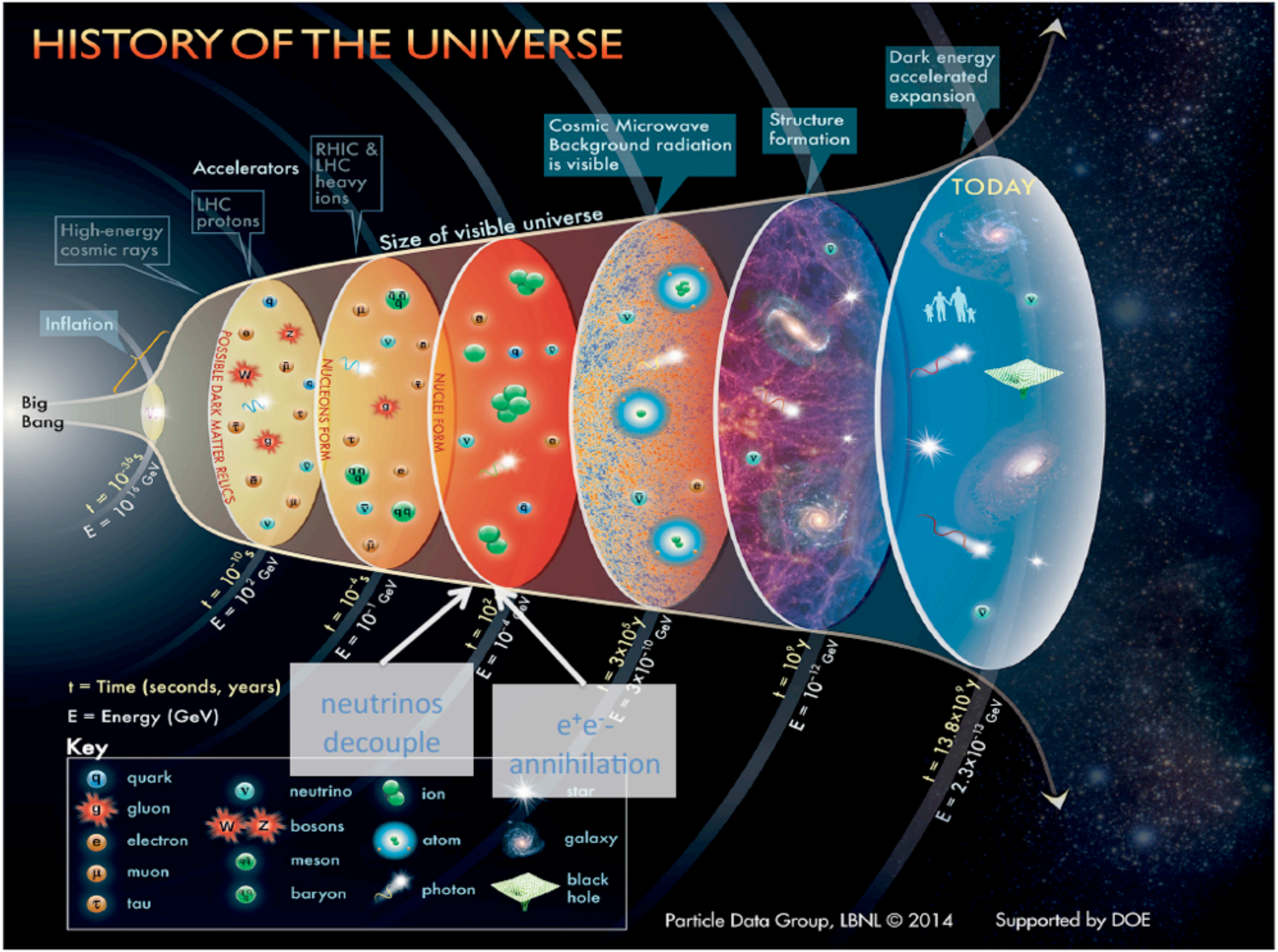


Cosmological constraints

(for more see e.g. seminar by Lattanzi at Roma Tre, 10th Jan 2017)

- cosmological observations

$$\sum m_\nu \equiv \sum_i m_i \qquad \text{(0.2 – 0.7 eV @ 95\%CL)}$$



HISTORY OF THE UNIVERSE
Dark energy accelerated expansion
Cosmic Microwave Background radiation is visible
Structure formation
RHIC & LHC heavy ions
Accelerators
LHC protons
High-energy cosmic rays
Inflation
Size of visible universe
TODAY
Big Bang
POSSIBLE DARK MATTER RELICS
NUCLEONS FORM
NUCLEI FORM
t = 10⁻³⁶ s
E = 10¹⁶ GeV
t = 10⁻¹⁰ s
E = 10² GeV
t = 10⁻⁴ s
E = 10⁻¹ GeV
t = 10² s
E = 10⁻⁴ GeV
t = 3x10⁵ y
E = 3x10⁻¹⁰ GeV
t = 10⁹ y
E = 10⁻¹² GeV
t = 13.8x10⁹ y
E = 2.3x10⁻¹³ GeV
neutrinos decouple
e⁺e⁻-annihilation
t = Time (seconds, years)
E = Energy (GeV)
Key
quark
gluon
electron
muon
tau
neutrino
bosons
meson
baryon
ion
atom
photon
star
galaxy
black hole
Particle Data Group, LBNL © 2014
Supported by DOE

The Cosmic Neutrino Background - Properties

At freeze out the neutrinos had a thermal velocity distribution.

Since then the neutrinos have continued to move along geodesics with a velocity which has red-shifted as a result of the expansion of the universe.

This geodesic movement is called *free streaming*.

These free-streaming neutrinos make up the *cosmic neutrino background*.

- Below $T \sim 1$ MeV, neutrino free stream keeping an equilibrium spectrum:

$$f_\nu(p) = \frac{1}{e^{p/T} + 1}$$

- Today $T_\nu = 1.9$ K and $n_\nu = 113$ part/cm^3 per species

Constraints from the Neutrino Background

In standard cosmologies, the cosmological neutrino background only interacts gravitationally after freeze-out and *all cosmological bounds on neutrino masses arise from gravitational interactions of the cosmic neutrino background.*

The gravitational interaction depends on the sum of the gravity from all of the neutrinos, which is proportional to the sum of the masses once the neutrinos have become nonrelativistic.
This is why cosmology constrains the sum of the masses.

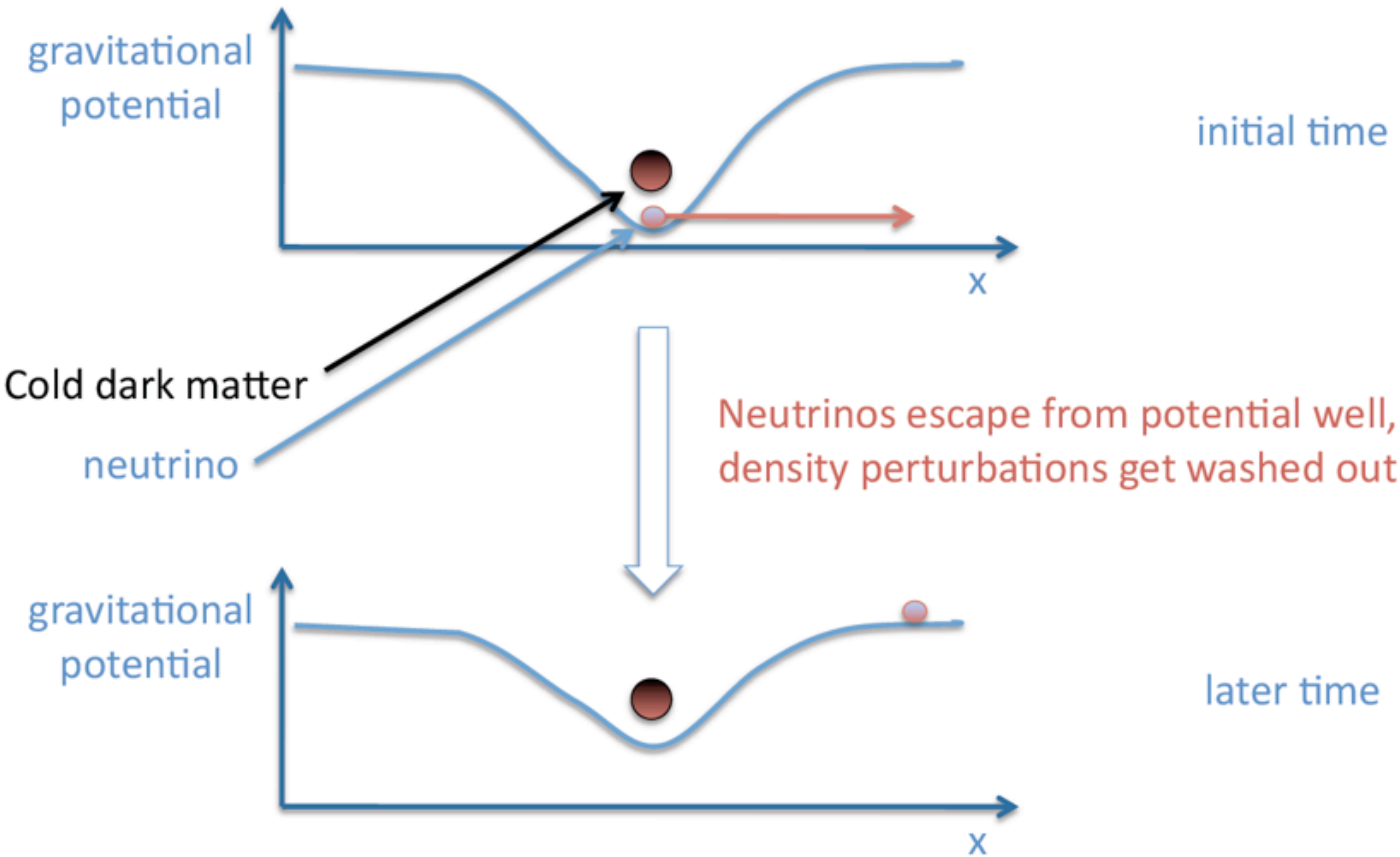
➡ Effects:

I) They contribute to the recent expansion of the universe identically to dark matter.

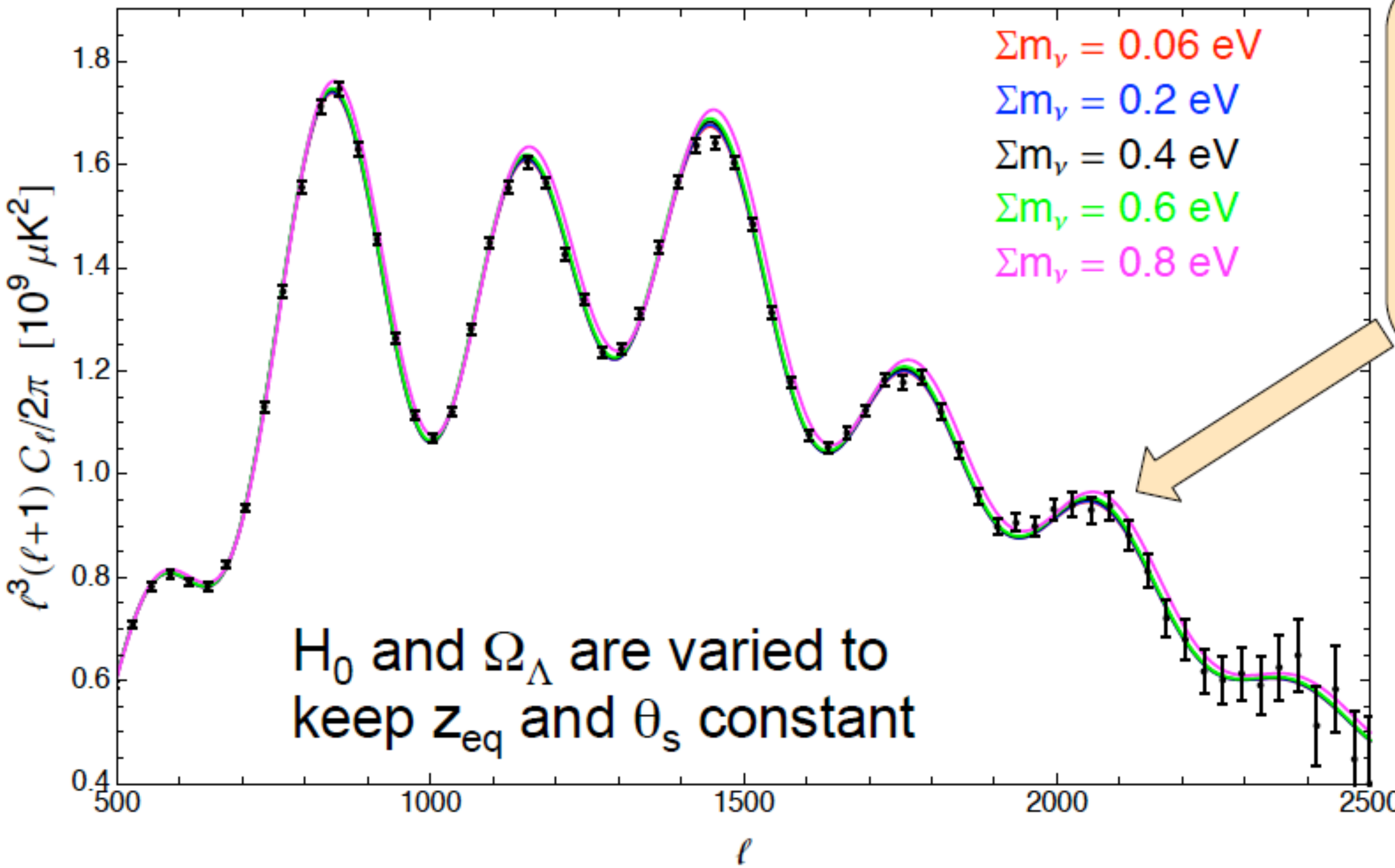
II) Since freeze out they free stream a distance called the *free-streaming length.* This disrupts structure formation on scales below the free-streaming length.

Free streaming

Velocity dispersion large wrt size of potential well



HOW HEAVY?



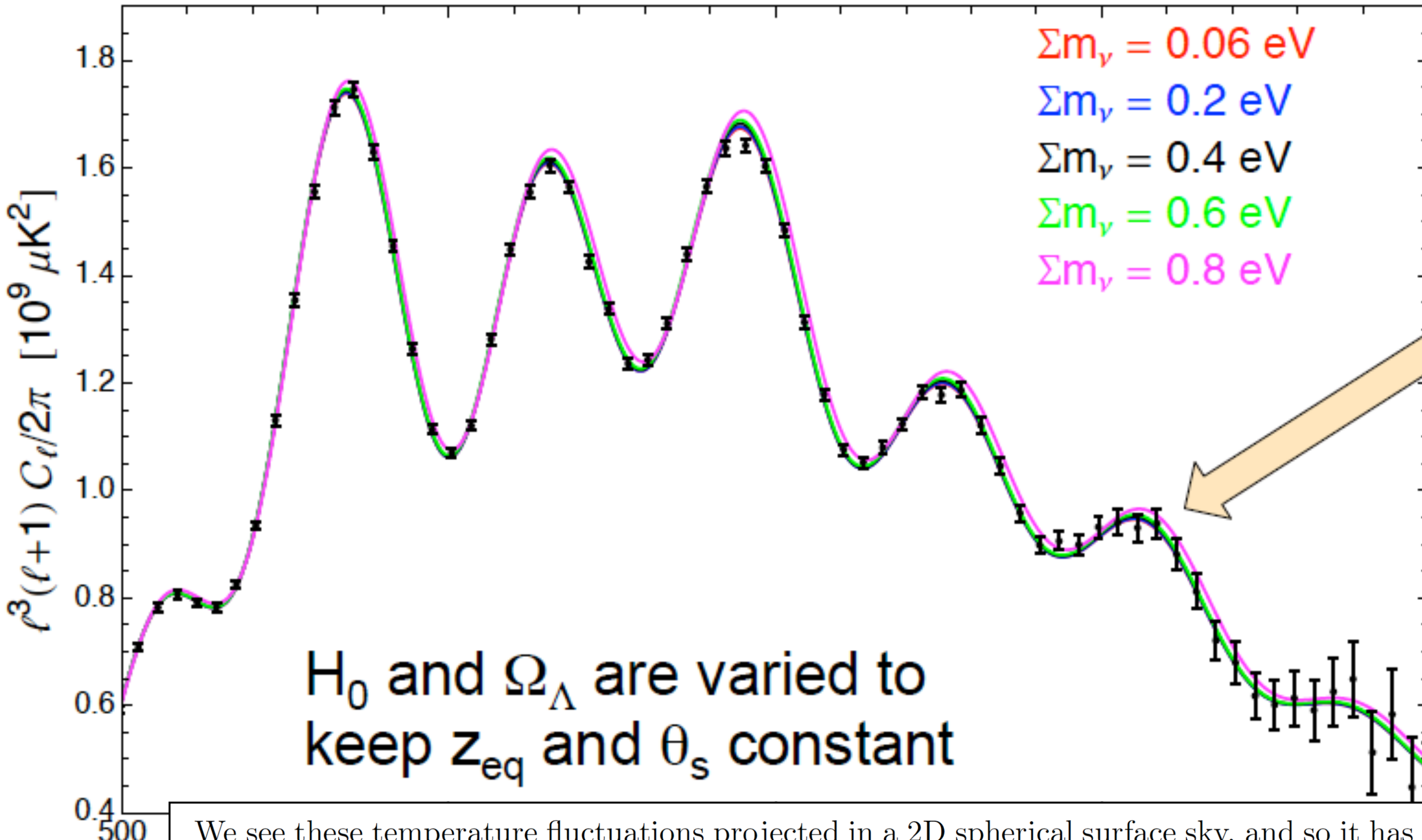
Net effect is to decrease lensing

Perturbations: free streaming, damping of small-scale perturbations

- proportional to the neutrino energy density
  - the effect is larger for larger masses

Model-dependent: interplay with Λ_{CDM}, H_0

How heavy?



Net effect is to decrease lensing

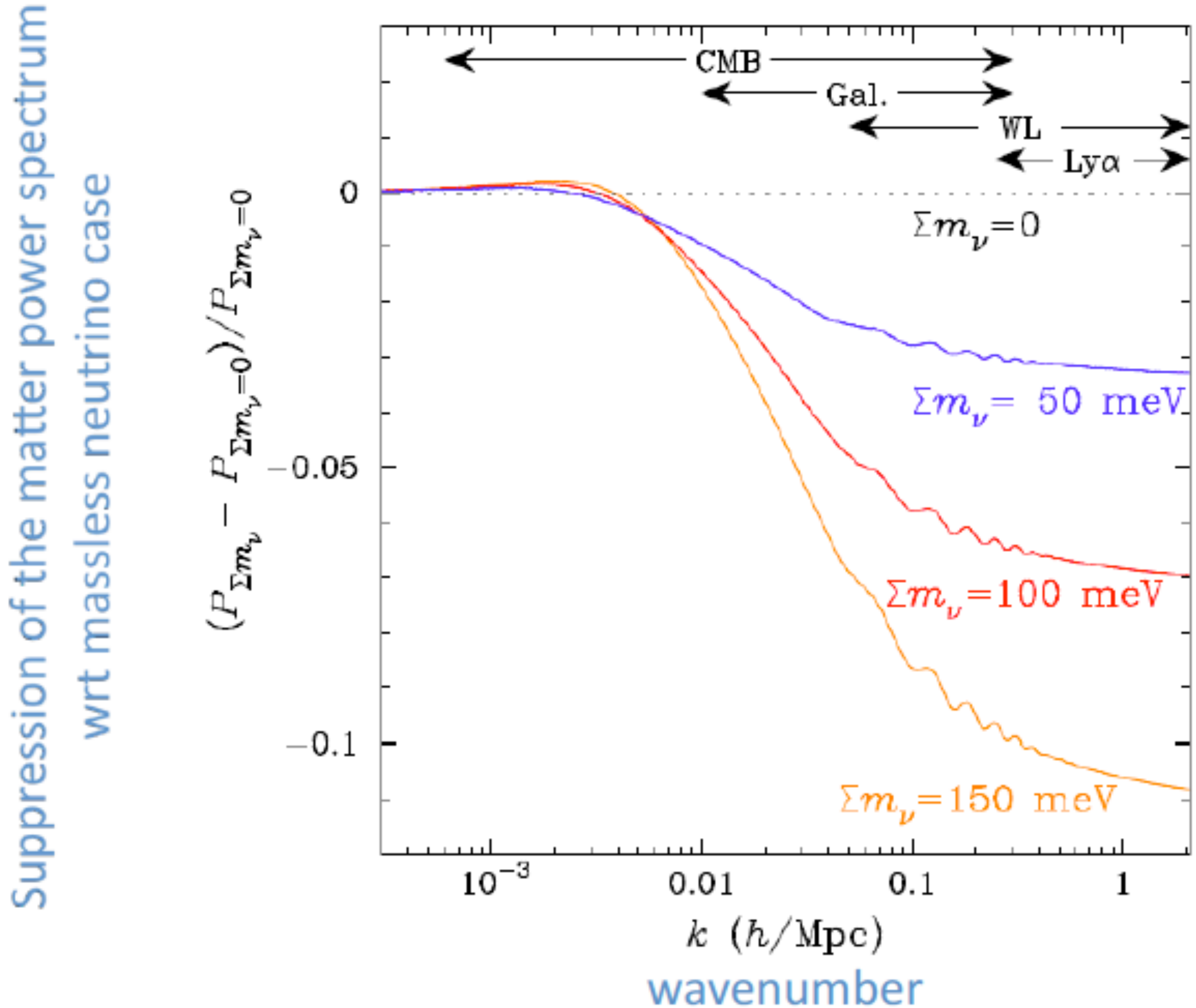
P ... ons

We see these temperature fluctuations projected in a 2D spherical surface sky, and so it has become common in the literature to expand the temperature field using spherical harmonics. The spherical harmonics form a complete orthonormal set on the unit sphere and are defined as

$$Y_{lm} = \sqrt{\frac{2\ell+1}{4\pi}\frac{(\ell-m)!}{(\ell+m)!}}P_\ell^m(cos\theta)e^{im\phi} \tag{2}$$

where the indices $\ell = 0, ..., \infty$ and $-\ell \leq m \leq \ell$ and P_ℓ^m are the Legendre polynomials. ℓ is called the multipole and represents a given angular scale in the sky α, given approximately by $\alpha = \pi/\ell$ (in degrees).

Matter power spectrum with massive neutrinos (at low redshifts)



$$\Delta\mathcal{P}/\mathcal{P} \approx 8\,\frac{\omega_\nu}{\omega_\mathrm{m}}$$

ratio of density of **ν** to matter (=non-relativistic particles) in the present universe

[Figure from Abazajian+ 2013]

FUTURE PROSPECTS FROM THE LAB

The absolute mass scale can be measured through:
(numbers on the right are forecast for future sensitivities)

- tritium beta decay

$$m_\beta \equiv \left[\sum |U_{ei}|^2 m_i^2\right]^{1/2}$$

(200 meV @ 68%CL)
(Katrin)

- neutrinoless double beta decay

$$m_{\beta\beta} \equiv \left|\sum U_{ei}^2 m_i\right|$$

(8 – 20 meV @ 90%CL)
(nEXO, 5-year exposure)

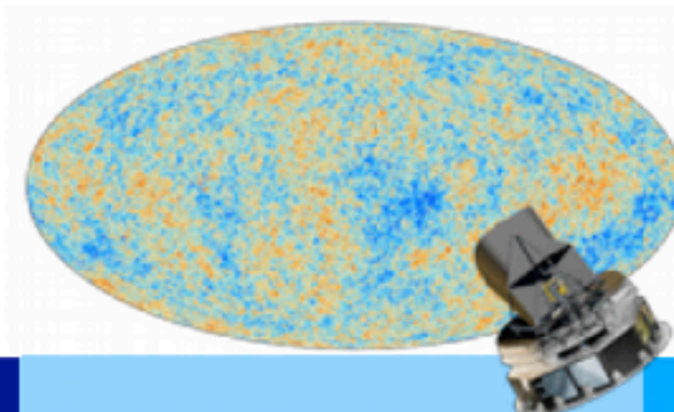
- cosmological observations

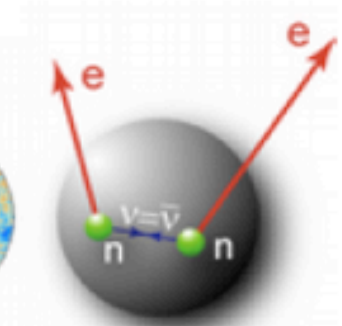
$$\sum m_\nu \equiv \sum_i m_i$$

(16 – 45 meV @ 68%CL)
(CORE, CORE+LSS)

Mass scale: experimental tools / 1

three complementary tools available
→ low temperature detectors play key role
(E. Fiorini and T. Niinikoski, Nucl. Instrum. and Meth. 224, p.83 (1984))





tool	Cosmology CMB+LSS+...	Neutrinoless Double Beta decay	Beta decay end-point
observable	$m_\Sigma=\Sigma_k m_{\nu_k}$	$m_{\beta\beta}=\|\Sigma_k m_{\nu_k} U_{ek}^2\|$	$m_\beta=(\Sigma_k m_{\nu_k}^2 \|U_{ek}\|^2)^{1/2}$
present sensitivity	≈0.1 eV	≈0.1 eV	2 eV
future sensitivity	**0.05 eV**	**0.05 eV**	**0.2 eV**
model dependency	yes ☹	yes ☹	no ☺
systematics	large ☹	yes 😐	large ☹

Summary & outlook

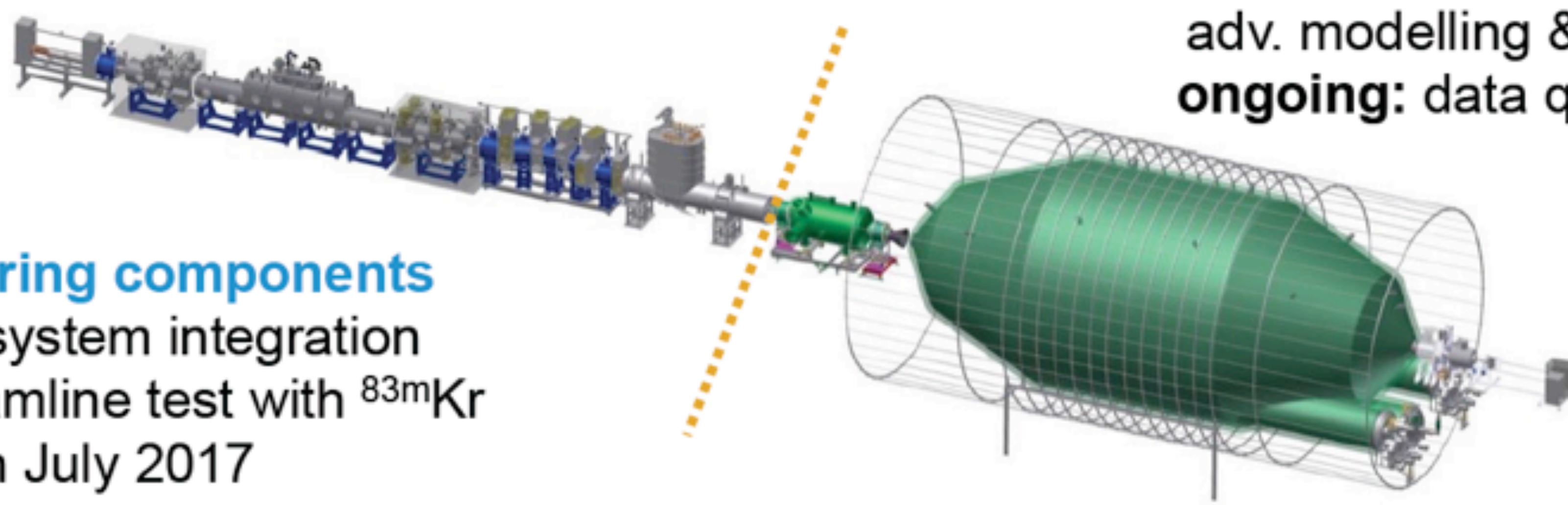
- β decay allows model-independent, **direct** access to neutrino mass scale
- KATRIN will exhaust degenerate mass regime: **200 meV** (90% CL for 5 yrs of running); reaching sub-eV sensitivity with first few weeks of data
- Interesting physics potential beyond m_ν: eV and keV scale sterile ν, RH currents, LIV, …

Preparing KATRIN for neutrino-mass measurements:

Analysis chain

adv. modelling & analysis framework

ongoing: data quality filters, blinding



Tritium-bearing components

- now: final system integration
- overall beamline test with ^{83m}Kr achieved in July 2017
- **next:** inactive commissioning with D_2, then $D_2(T_2)$

Spectrometer & detector section

2 successful commissioning phases already done

ongoing: background investigations

➜ **First tritium runs starting in 2018, inauguration ceremony: 11 June 2018**