



# *The Discovery of the Higgs at LHC*

Year 2023

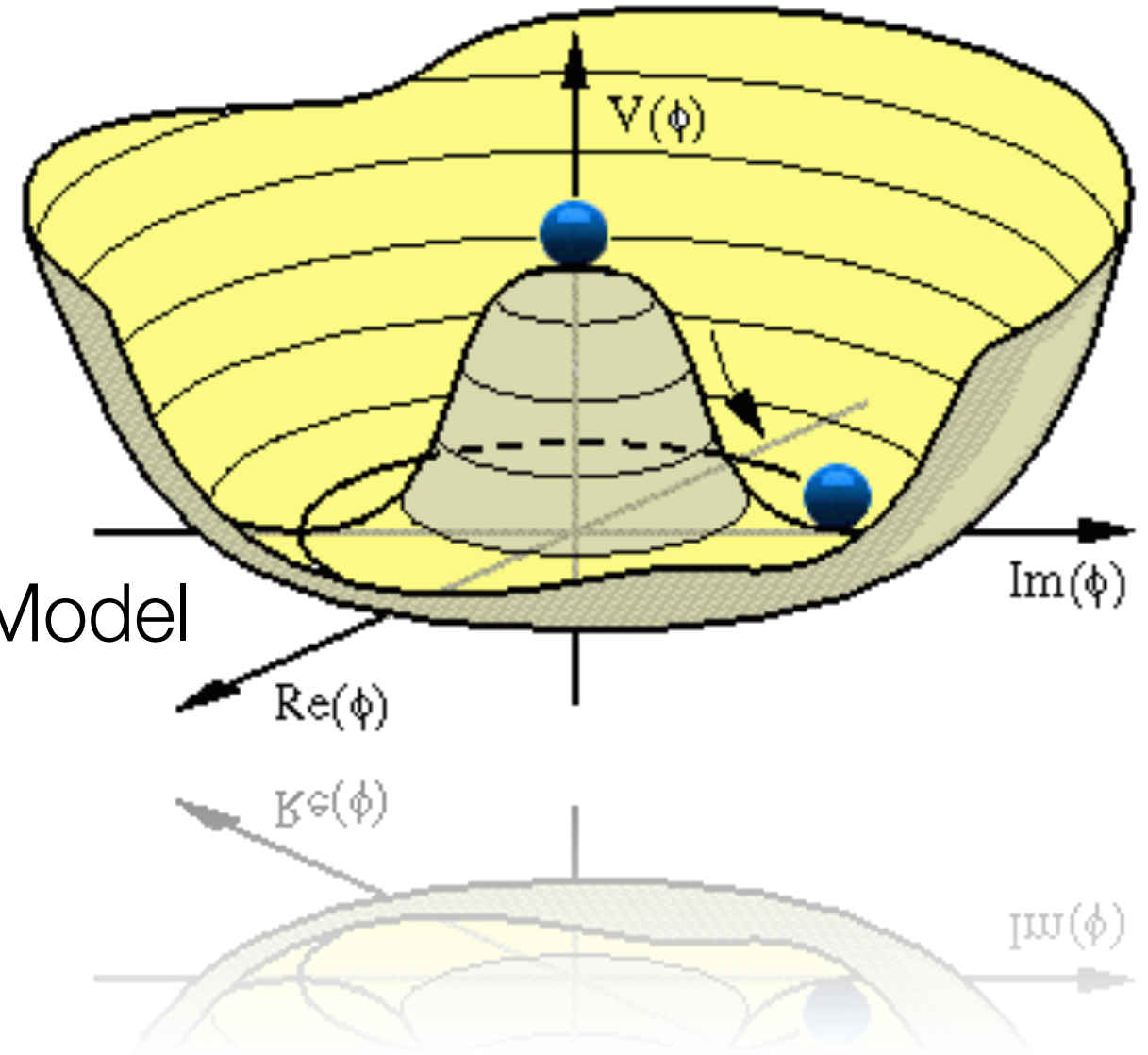
*Collider Physics*  
*Toni Baroncelli*



# The Discovery of the Higgs at LHC

## The Higgs

(The “once”)  
Missing Piece in the Standard Model



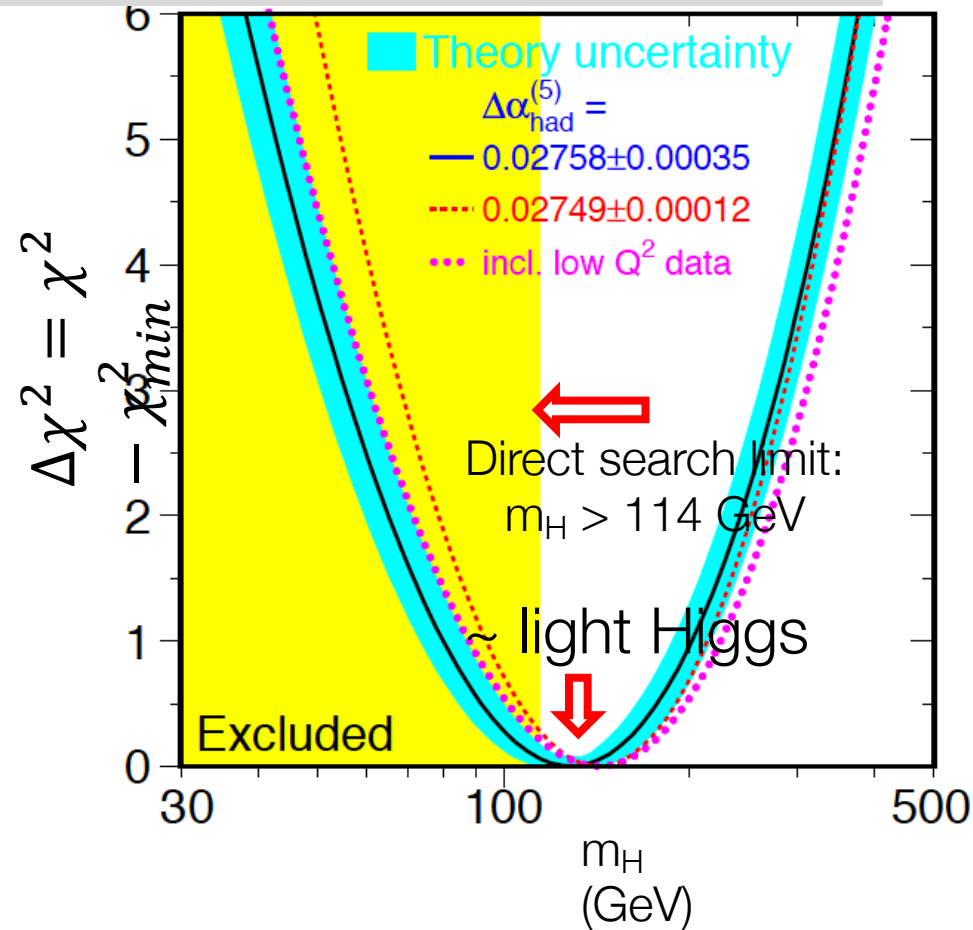
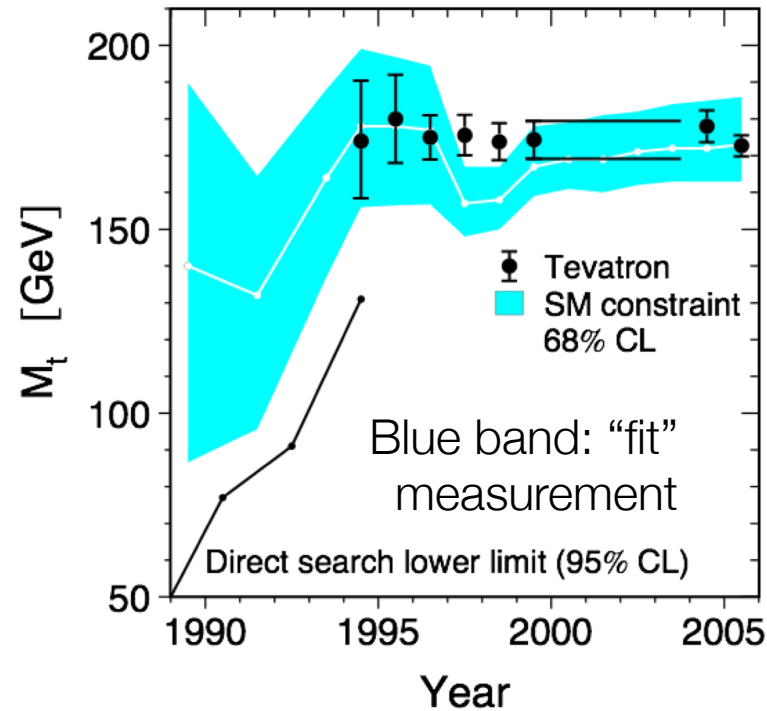
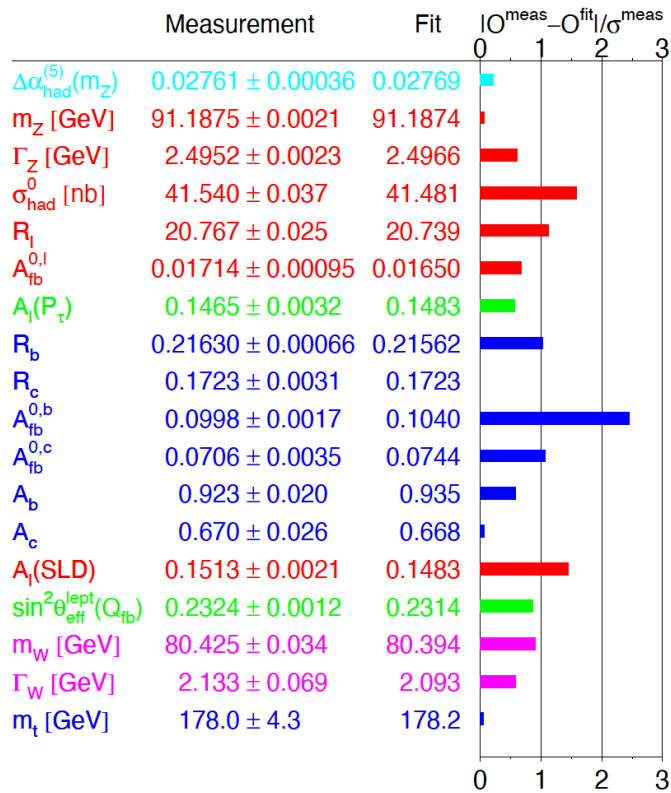


# $m_t$ and $m_H \rightarrow EW$ fits

There are very many EW measurements: cross sections, asymmetries and many others.

Some diagrams contain loops where top and Higgs bosons circulate  $\rightarrow$  slightly modify observables  $\rightarrow$  indication on top mass and Higgs mass before their discoveries

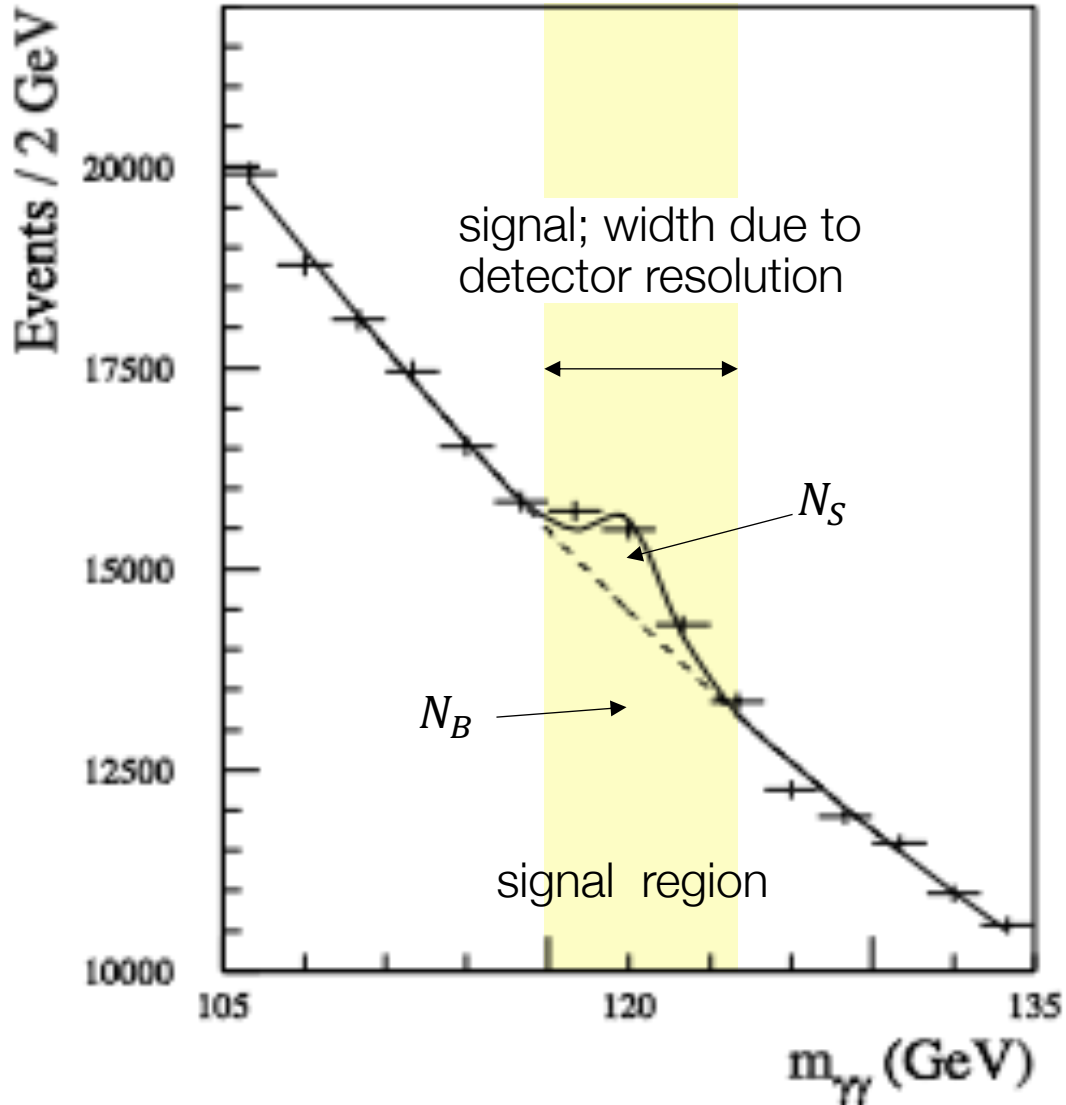
Summer 2004





# Higgs and the Tevatron: Reminder on Discoveries

The Higgs Boson was also searched for at the Tevatron. Without success! Why?



Signal significance:

$$S = \frac{N_S}{\sqrt{N_B + N_S}}$$

$N_S$ : # signal events

$N_B$ : # background events

In the “signal region”

Kinematically out of reach?

By “convention” a discovery is claimed when the significance

$$S > 5:$$

This means that the signal

$$N_S = N_{\text{tot}} - N_B$$

is 5 times larger than statistical uncertainty on  $N_B + N_S$  → the probability of a fluctuation is very small: the Gaussian probability that upward fluctuation by more than  $5\sigma$  is observed is

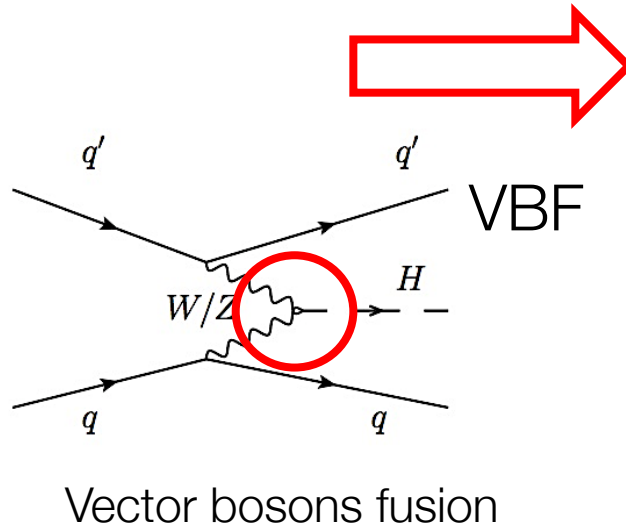
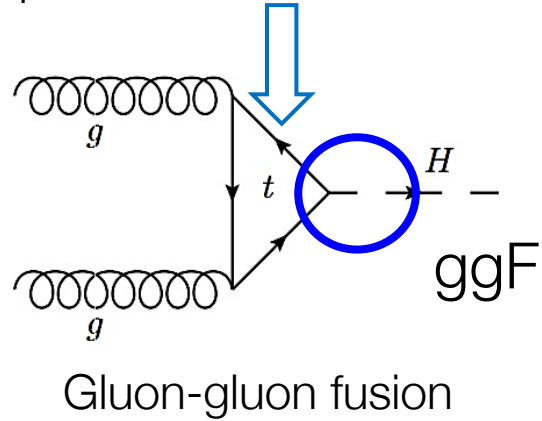
$$P_{5\sigma} = 10^{-7}$$

The sensitivity to a signal increases with increasing statistics

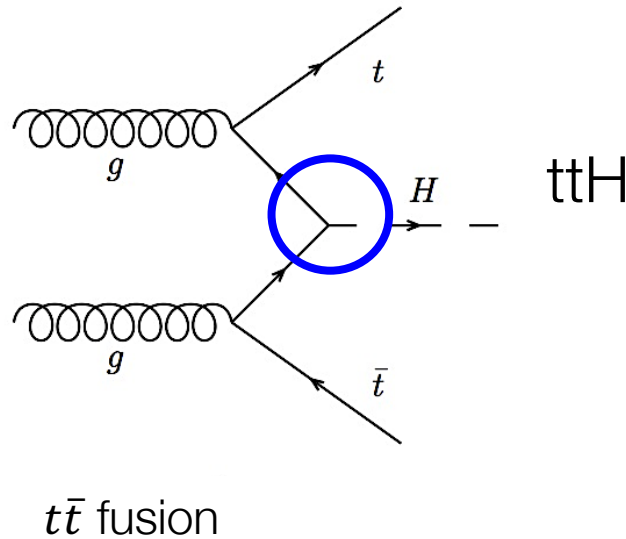
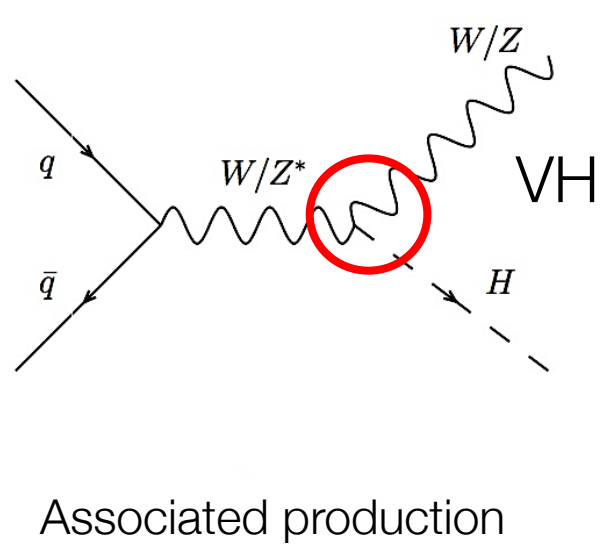
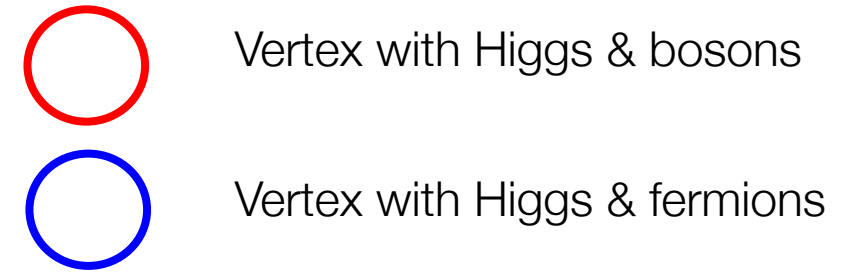


# Higgs Production Mechanisms at Hadron Colliders

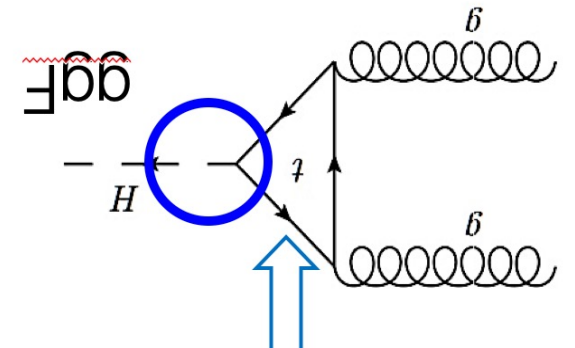
Loop with circulating **tops**, other quarks contribute much less



Higgs couples to massive particles, cannot couple directly to gluons and photons

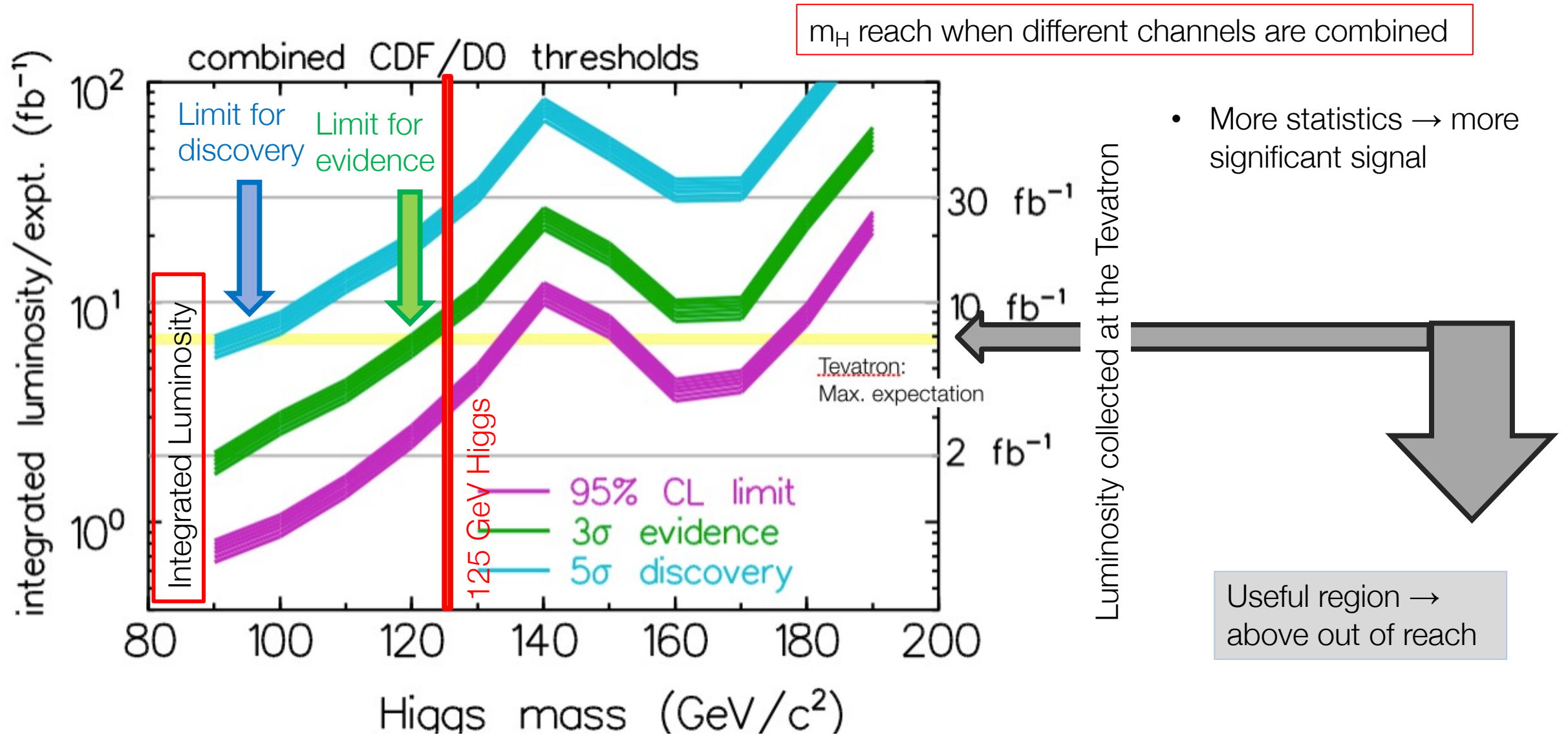


We will find the same vertices also in the Higgs decay!





# Higgs Search at the Tevatron (forecast)





# Higgs Decay Channels Investigated by CDF & D0

CDF

Channel	Luminosity (fb <sup>-1</sup> )	$m_H$ range (GeV/c <sup>2</sup> )	Reference
$WH \rightarrow \ell\nu b\bar{b}$ (ST,DT,2,3 jet)	5.3	100-150	[14]
$VH \rightarrow \tau^+\tau^- b\bar{b}/q\bar{q}\tau^+\tau^-$	4.9	105-145	[15, 16]
$ZH \rightarrow \nu\bar{\nu} b\bar{b}$ (ST,TLDT)	5.2-6.4	100-150	[17, 18]
$ZH \rightarrow \ell^+\ell^- b\bar{b}$ (ST,DT,ee,μμ,ee <sub>ICR</sub> ,μμ <sub>trk</sub> )	4.2-6.2	100-150	[19]
$VH \rightarrow \ell^\pm\ell^\pm + X$	5.3	115-200	[20]
$H \rightarrow W^+W^- \rightarrow e^\pm\nu e^\mp\nu, \mu^\pm\nu\mu^\mp\nu$	5.4	115-200	[21]
$H \rightarrow W^+W^- \rightarrow e^\pm\nu\mu^\mp\nu$ (0,1,2+ jet)	6.7	115-200	[22]
$H \rightarrow W^+W^- \rightarrow \ell\bar{\nu}jj$	5.4	130-200	[23]
$H \rightarrow \gamma\gamma$	4.2	100-150	[24]
$t\bar{t}H \rightarrow t\bar{t}b\bar{b}$ (ST,DT,TT,4,5+ jets)	2.1	105-155	[25]

Sensitivity range of the channel

At high  $m_H \rightarrow$  decay to WW

D0

Channel	Luminosity (fb <sup>-1</sup> )	$m_H$ range (GeV/c <sup>2</sup> )	Reference
$WH \rightarrow \ell\nu b\bar{b}$ 2-jet channels $4 \times$ (TDT,LDT,ST,LDTX)	5.7	100-150	[5]
$WH \rightarrow \ell\nu b\bar{b}$ 3-jet channels $2 \times$ (TDT,LDT,ST)	5.6	100-150	[6]
$ZH \rightarrow \nu\bar{\nu} b\bar{b}$ (TDT,LDT,ST)	5.7	100-150	[7]
$ZH \rightarrow \ell^+\ell^- b\bar{b}$ $4 \times$ (TDT,LDT,ST)	5.7	100-150	[8, 9]
$H \rightarrow W^+W^-$ $2 \times$ (0,1 jets)+(2+ jets)+(low- $m_{\ell\ell}$ )+(e- $\tau_{had}$ )+(μ- $\tau_{had}$ )	5.9	110-200	[10]
$WH \rightarrow WW^+W^-$ (same-sign leptons 1+ jets)+(tri-leptons)	5.9	110-200	[10]
$ZH \rightarrow ZW^+W^-$ (tri-leptons 1 jet)+(tri-leptons 2+ jets)	5.9	110-200	[10]
$H + X \rightarrow \tau^+\tau^-$ (1 jet)+(2 jets)	2.3	100-150	[11]
$WH + ZH \rightarrow jjb\bar{b}$ $2 \times$ (TDT,LDT)	4.0	100-150	[12]
$H \rightarrow \gamma\gamma$	5.4	100-150	[13]

At high  $m_H \rightarrow$  decay to WW

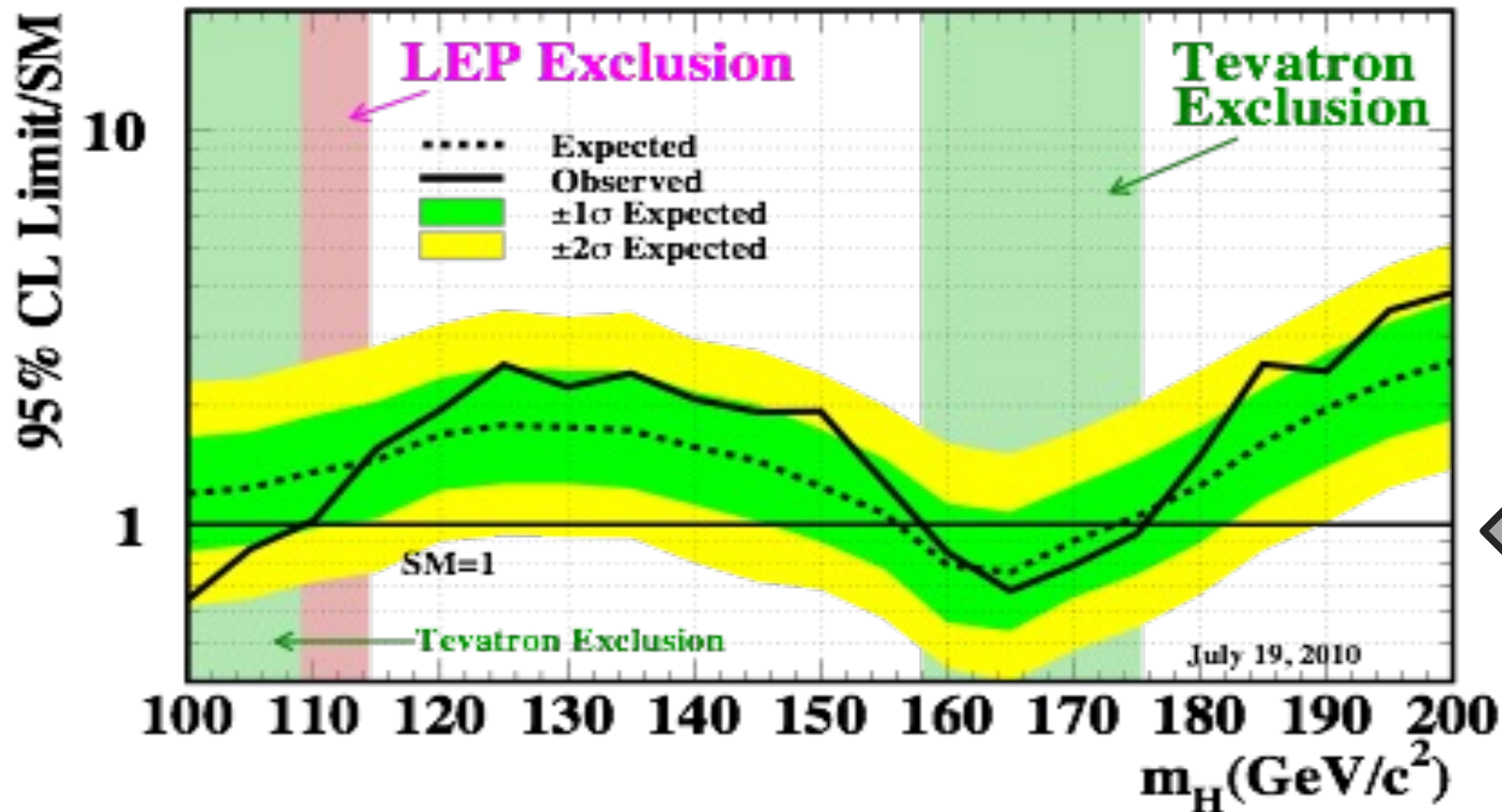


# Tevatron Upper Limits on Higgs Mass

Observed and expected 95% C.L. upper limits on the ratios to the SM cross section, as functions of the Higgs boson mass for the combined CDF and D0 analyses :

If the experiments could be repeated 100 times, 95% of times they would get the same 'exclusion' result

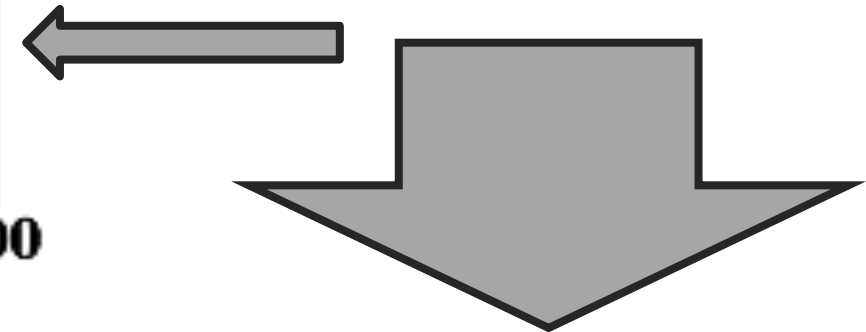
Tevatron Run II Preliminary,  $\langle L \rangle = 5.9 \text{ fb}^{-1}$



Ups  $\uparrow$  and downs  $\downarrow$  depend on the decay channel used in the search

Some channels are more sensitive than others

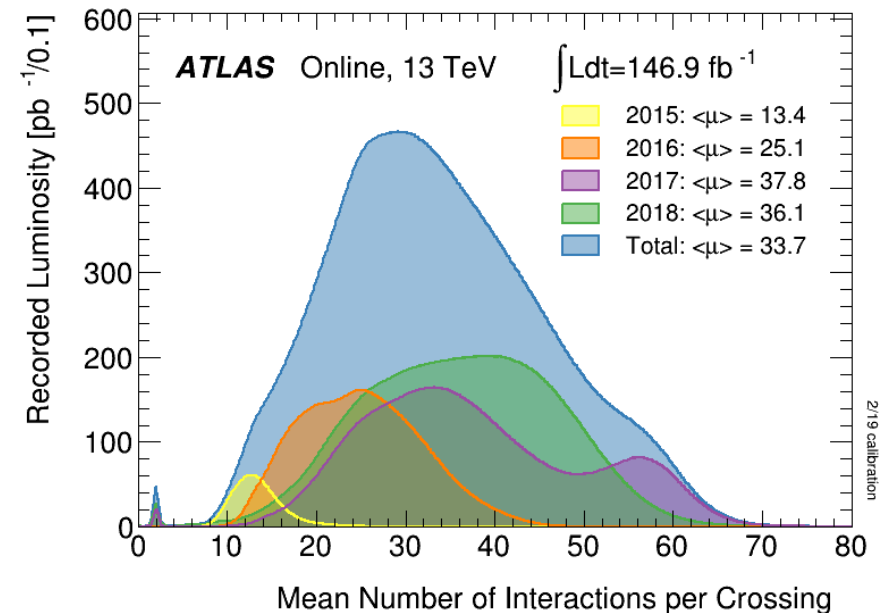
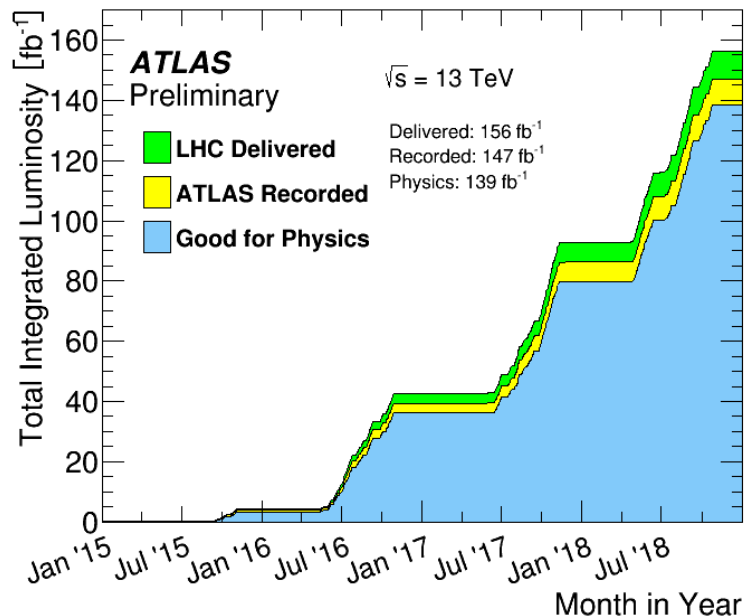
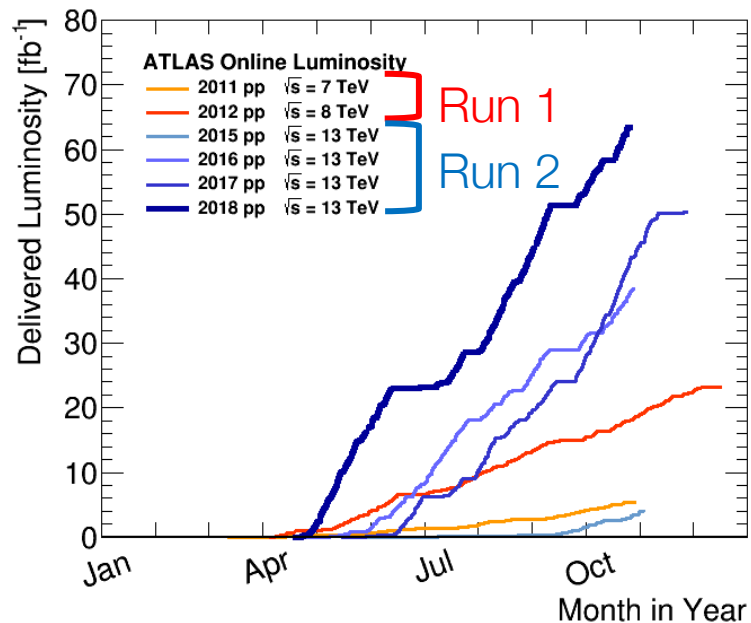
Exclusion region  $\rightarrow$  above out of reach







# Delivered Luminosity by LHC in Run 1 & 2



Integrated luminosity in LHC (fb<sup>-1</sup>)

- Run 1 (7 and 8 TeV)
- Run 2 (13 TeV)

Delivered by LHC in Run 2: 156 fb<sup>-1</sup>  
Recorded by ATLAS: 147 fb<sup>-1</sup>

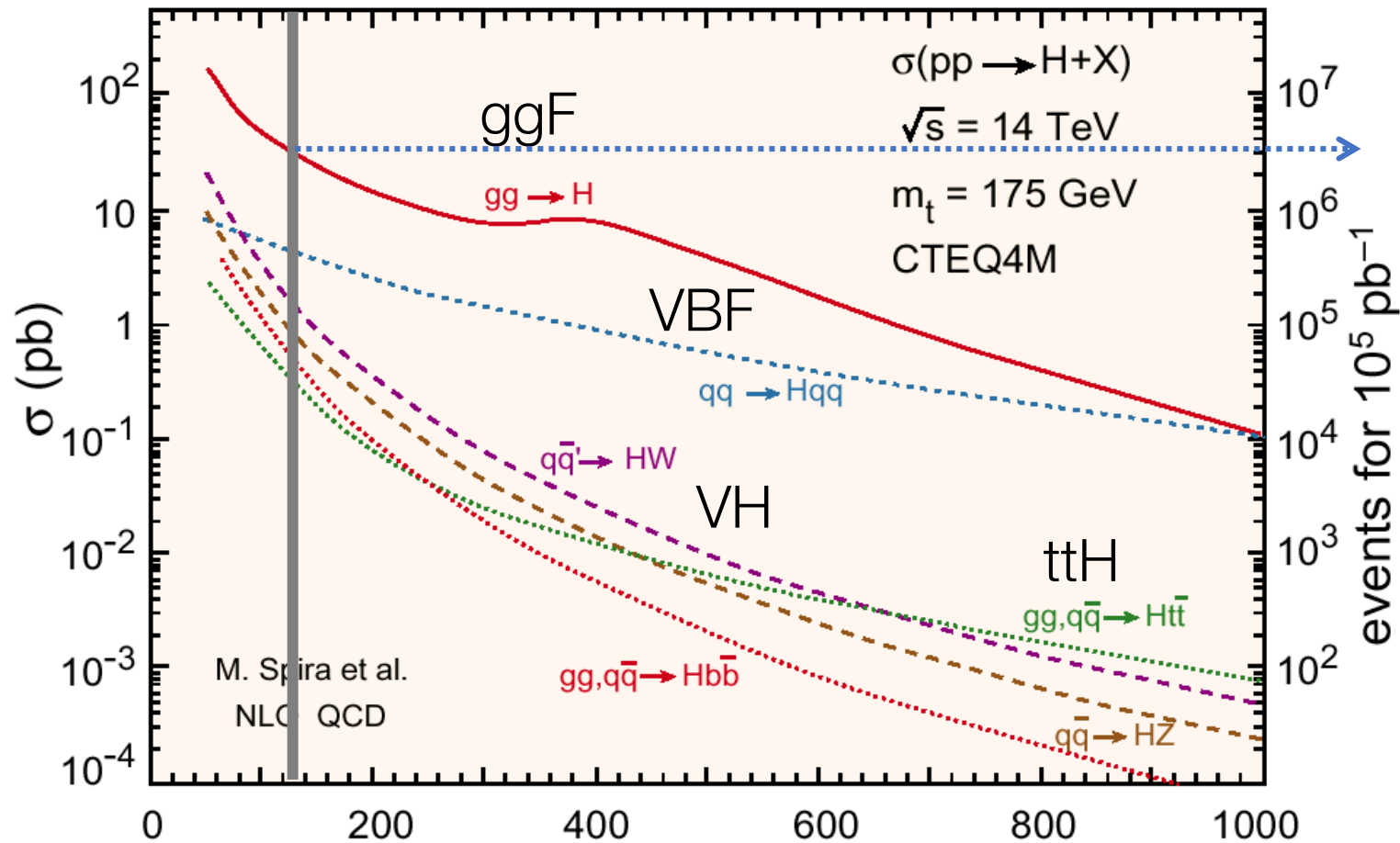
Year	2010	2011	2012	2015	2016	2017	2018
Luminosity delivered (fb <sup>-1</sup> )	0.05	6.1	23.3	4.2	41	50	68
CMS Energy	7	7	8	13			



# Higgs Search at the LHC

LHC: < 13 TeV!

Higgs production cross-section at centre-of-mass energies of 14 TeV, as a function of  $m_H$ .



Different production mechanisms are shown

$10^5 \text{ pb}^{-1} = 100 \text{ fb}^{-1}$

'Millions' of 125 GeV Higgs produced by the ggF mechanism at 14 TeV for  $100 \text{ fb}^{-1}$  total luminosity



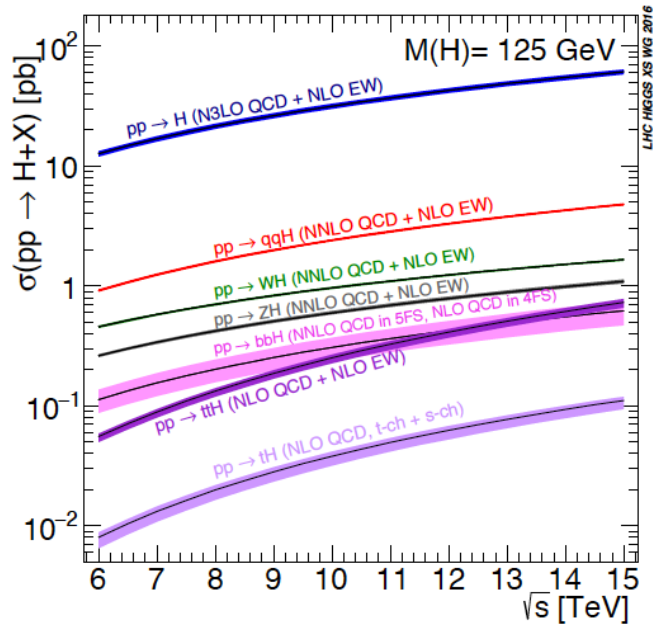
Acceptance and efficiency reduces this number drastically!

LHC cms 14 TeV

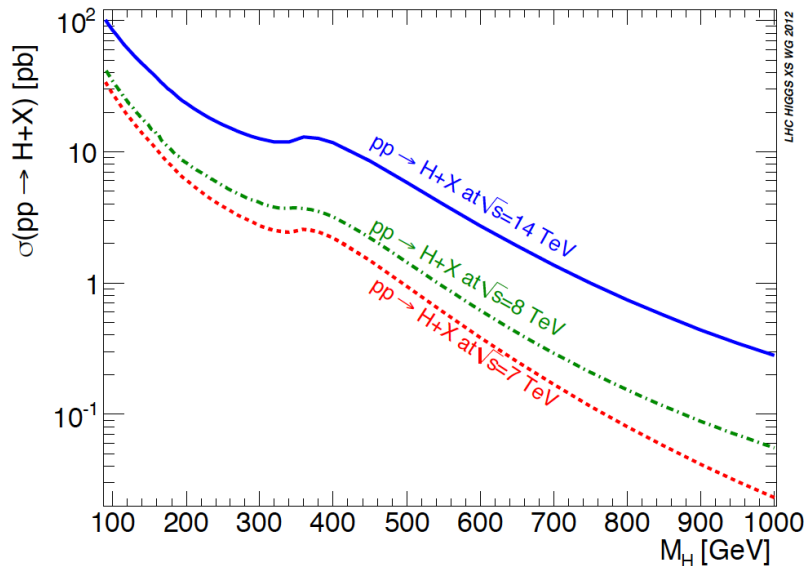
$M_H$  (GeV)



# Cross Section vs $\sqrt{s}$ for a 125 GeV Higgs



$\sqrt{s}$ (TeV)	Production cross section (in pb) for $m_H = 125$ GeV						total	Calculation
	ggF	VBF	WH	ZH	$t\bar{t}H$			
1.96	0.95 <sup>+17%</sup> <sub>-17%</sub>	0.065 <sup>+8%</sup> <sub>-7%</sub>	0.13 <sup>+8%</sup> <sub>-8%</sub>	0.079 <sup>+8%</sup> <sub>-8%</sub>	0.004 <sup>+10%</sup> <sub>-10%</sub>	1.23	Tevatron	
7	16.9 <sup>+5%</sup> <sub>-5%</sub>	1.24 <sup>+2%</sup> <sub>-2%</sub>	0.58 <sup>+3%</sup> <sub>-3%</sub>	0.34 <sup>+4%</sup> <sub>-4%</sub>	0.09 <sup>+8%</sup> <sub>-14%</sub>	19.1	LHC Run 1	
8	21.4 <sup>+5%</sup> <sub>-5%</sub>	1.60 <sup>+2%</sup> <sub>-2%</sub>	0.70 <sup>+3%</sup> <sub>-3%</sub>	0.42 <sup>+5%</sup> <sub>-5%</sub>	0.13 <sup>+8%</sup> <sub>-13%</sub>	24.2		
13	48.6 <sup>+5%</sup> <sub>-5%</sub>	3.78 <sup>+2%</sup> <sub>-2%</sub>	1.37 <sup>+2%</sup> <sub>-2%</sub>	0.88 <sup>+5%</sup> <sub>-5%</sub>	0.50 <sup>+9%</sup> <sub>-13%</sub>	55.1	LHC Run 2	
14	54.7 <sup>+5%</sup> <sub>-5%</sub>	4.28 <sup>+2%</sup> <sub>-2%</sub>	1.51 <sup>+2%</sup> <sub>-2%</sub>	0.99 <sup>+5%</sup> <sub>-5%</sub>	0.60 <sup>+9%</sup> <sub>-13%</sub>	62.1	LHC Run HL	



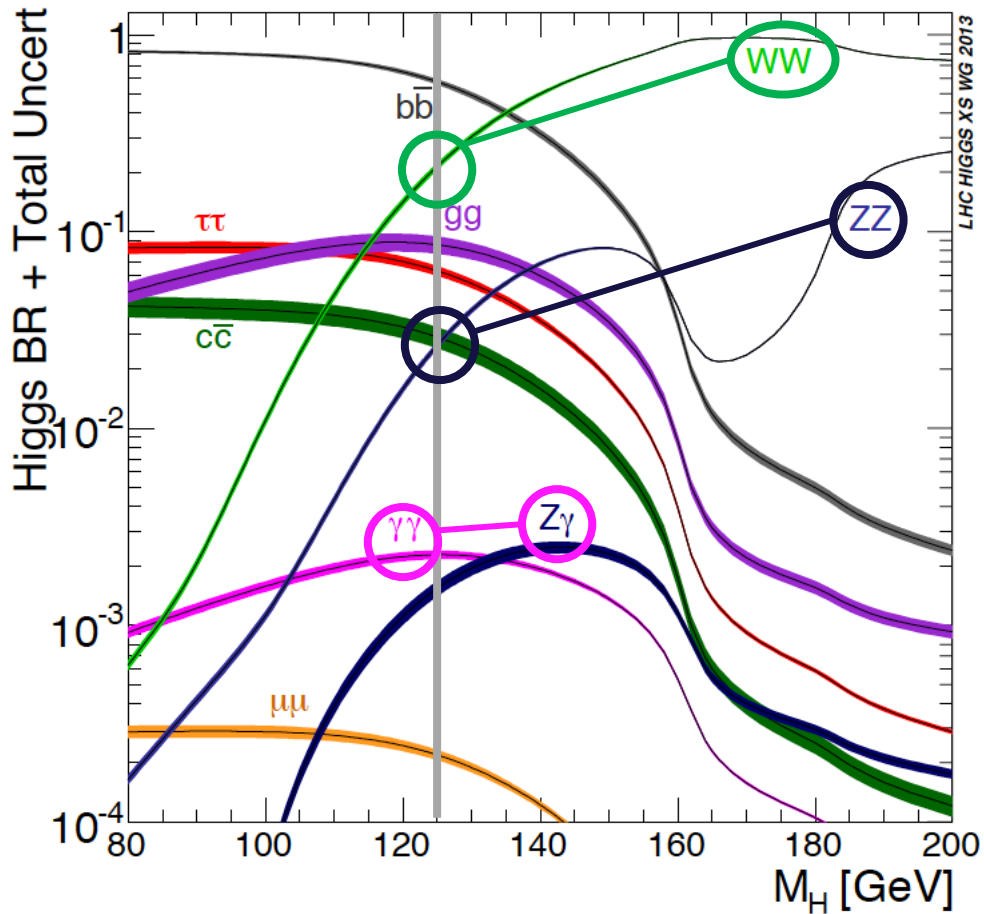
- Relative ratios between different production modes is ~ constant at LHC energies: ggF  $\rightarrow$  VBF  $\rightarrow$  VH  $\rightarrow$  ttH
- The total cross section for the production of an Higgs of 125 GeV increases significantly with centre-of-mass energy & decreases with mass

$$\frac{\sigma_{Higgs(125\text{ GeV})}^{\sqrt{s}=1.96\text{ TeV}}}{\sigma_{Higgs(125\text{ GeV})}^{\sqrt{s}=13\text{ TeV}}} = 2.2\%$$

Tevatron: ggF then WH  
LHC: ggF then VBF



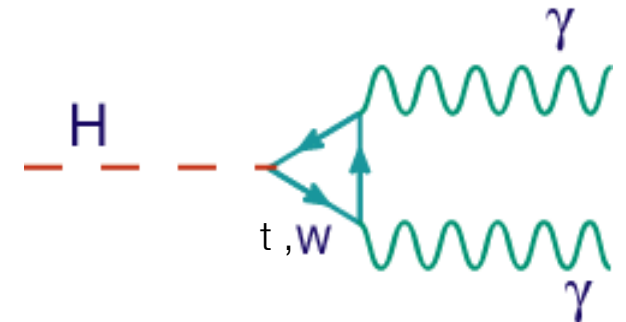
# Higgs Decay Branching Fractions



Discovery driven by significance of different channels

What really counts is how well one can distinguish an Higgs signal from the background

Example: the decay of the Higgs into a pair of photons is very small ( $H \rightarrow WW$  is  $\sim 100$  times larger than  $H \rightarrow \gamma\gamma$ ), however the distinct topology it generates made it very important in the Higgs discovery

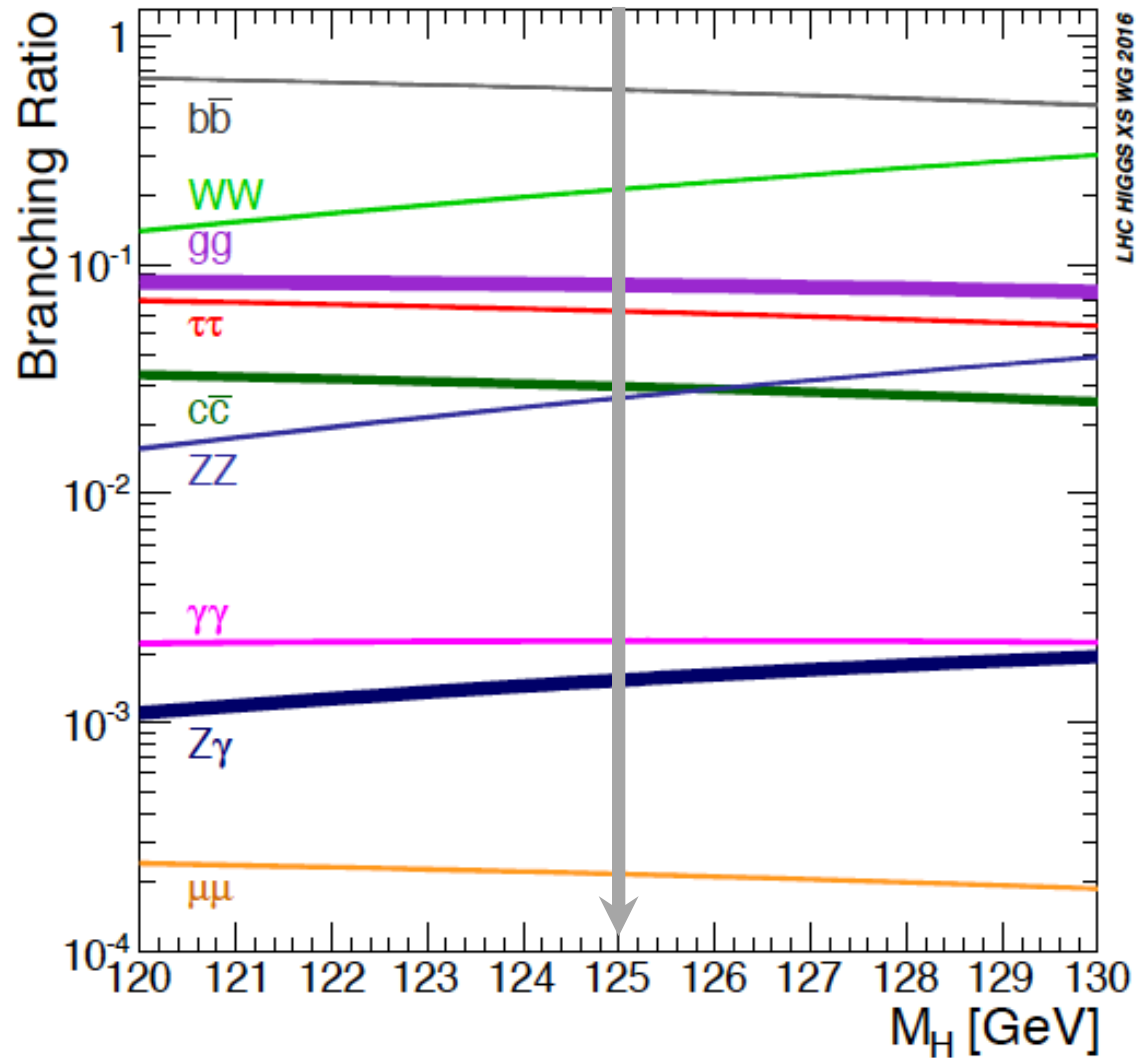


- Channels used for the Higgs discovery:
  - $ZZ$  to 4 leptons
  - Two photons
  - Two  $WW$  to leptons and neutrinos (a bit late!)

For  $m_H = 125$  GeV:  $H \rightarrow bb, WW, gg, \tau\tau$   
 For  $m_H > 160$  GeV:  $H \rightarrow WW, ZZ$  dominant



# Zooming the Branching Ratios of the Higgs



$m_H = 125 \text{ GeV}$

Decay channel	Branching ratio	Rel. uncertainty
$H \rightarrow \gamma\gamma$	$2.27 \times 10^{-3}$	+5.0% -4.9%
$H \rightarrow ZZ$	$2.62 \times 10^{-2}$	+4.3% -4.1%
$H \rightarrow W^+W^-$	$2.14 \times 10^{-1}$	+4.3% -4.2%
$H \rightarrow \tau^+\tau^-$	$6.27 \times 10^{-2}$	+5.7% -5.7%
$H \rightarrow b\bar{b}$	$5.84 \times 10^{-1}$	+3.2% -3.3%
$H \rightarrow Z\gamma$	$1.53 \times 10^{-3}$	+9.0% -8.9%
$H \rightarrow \mu^+\mu^-$	$2.18 \times 10^{-4}$	+6.0% -5.9%



# The Story of the Higgs (Discovery) in one Slide

- Indirect bounds on  $m_H$  from global EW fits : two decades at LEP, SLC, Tevatron suggest a ~light Higgs

$$m_H = 89^{+35}_{-26} \text{ GeV}$$

- Direct and model-independent search at LEP up to 209 GeV cms gave a 95% CL lower bound on  $m_H$

$$m_H > 114.4 \text{ GeV } 95\% \text{ CL}$$

- Direct search after LEP shutdown in 2000 at Tevatron ppbar collider using 10fb-1 gave



a] excluded intervals 90-109 GeV and 149-182 GeV

b] broad excess at the level of  $3 \sigma$  in the interval  $115 < m_H < 140$  GeV with a maximum at 125 GeV

- LHC run in 2011 (7 TeV, 5 fb<sup>-1</sup>), 2012 (8 TeV, 20 fb<sup>-1</sup>) gave evidence for a new particle decaying to  $\gamma\gamma$  and ZZ with rates as predicted by SM. Evidence for decays to  $W^+W^-$  but no evidence for  $b\bar{b}$  and  $\tau^+\tau^-$
- LHC July 2012 : ATLAS & CMS claim a discovery of a new particle with a mass of about 125 GeV



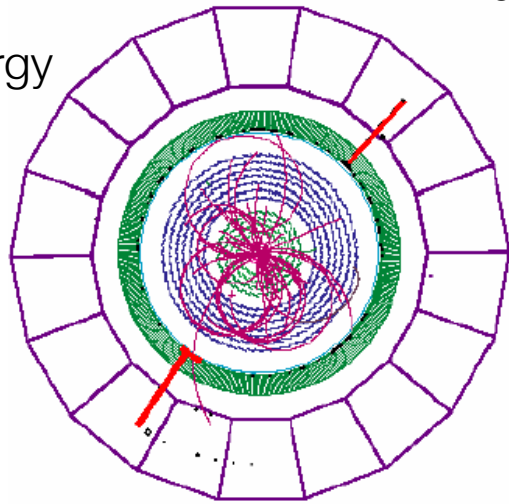
# Most Promising Higgs Decay Channels

Channel	LHC Potential
$gg \rightarrow H \rightarrow bb$	Huge QCD background ( $gg \rightarrow bb$ ); extremely difficult
$gg \rightarrow H \rightarrow \tau\tau$	Higgs with low $p_T$ , hard to discriminate from background; problematic
$gg \rightarrow H \rightarrow \gamma\gamma$	Small rate, large combinatorial background, but excellent determination of $m_H$ (CMS: crystal calorimeter)
$gg \rightarrow H \rightarrow WW$	Large rate, but 2 neutrinos in leptonic decay, Higgs spin accessible via lepton angular correlations
$gg \rightarrow H \rightarrow ZZ$	$ZZ \rightarrow 4\mu$ : "gold-plated" channel for high-mass Higgs (ATLAS: muon spectrometer)

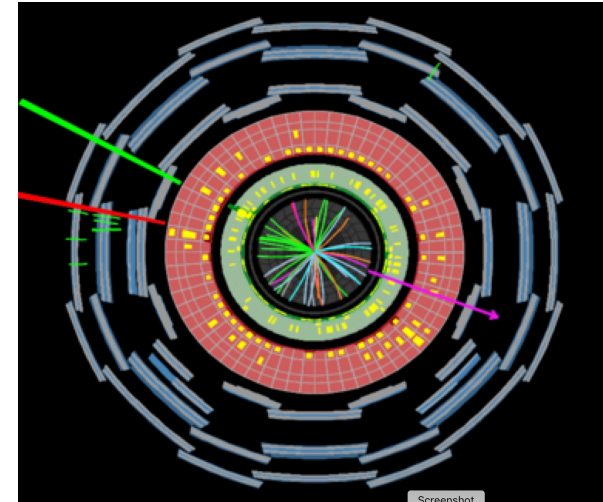
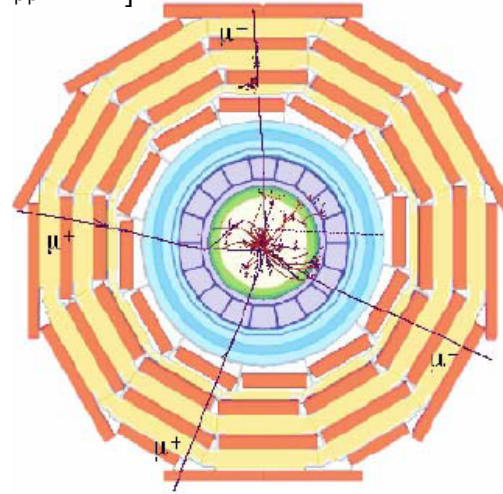


# Topologies!

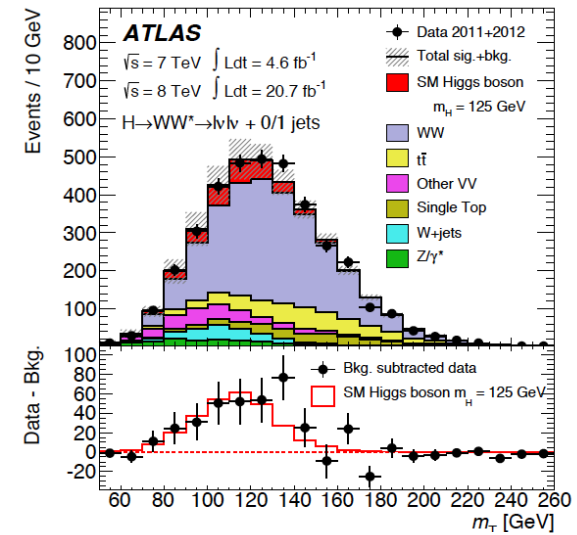
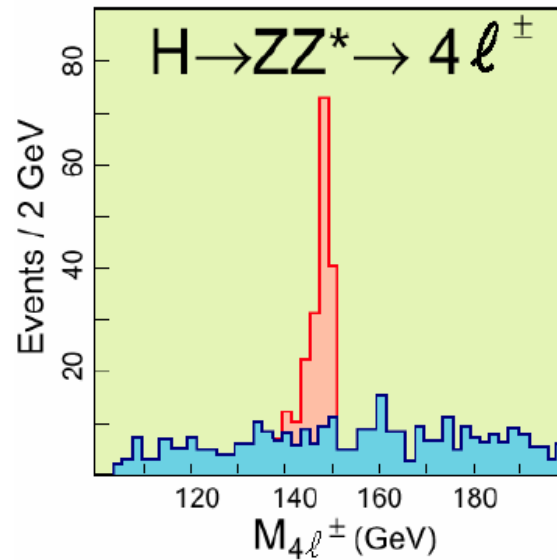
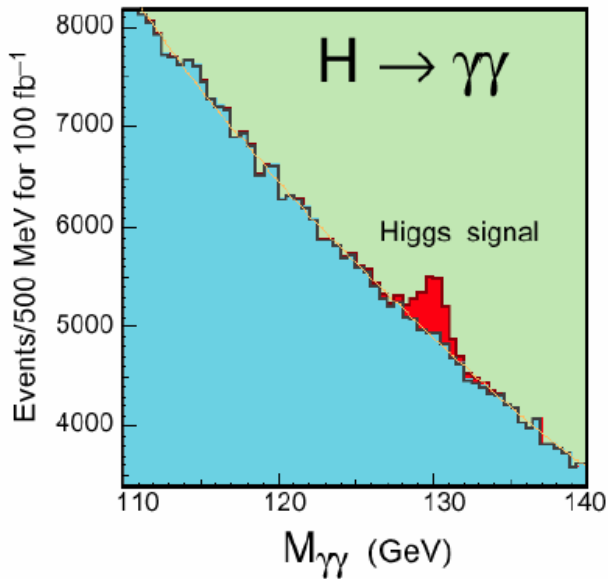
Two high-energy photons



4 muons [ $M_{\mu\mu} = M_Z$ ]



Electron + muon + MET







# Complications of Real Life: Background

1. Choose channels with low SM background

not possible:  $H \rightarrow bb$  ... without associated production  
 ...(VH mode,  $V=W,Z$ )

possible:  $H \rightarrow \gamma\gamma$  ... despite of small branching ratio .  
 $H \rightarrow ZZ$  ... with one Z decaying leptonically  
 $H \rightarrow WW$  ... large signal and large background

2. Optimize detector resolution

Example: mass resolution  $\sigma_m$  worsen by a factor of 2;

$N_s$  doesn't change; signal region has to be increased by a factor 2

$\rightarrow$  number  $N_B$  of background events increases by factor of 2

$S = N_s/\sqrt{N_B}$  decreases by  $\sqrt{2}$

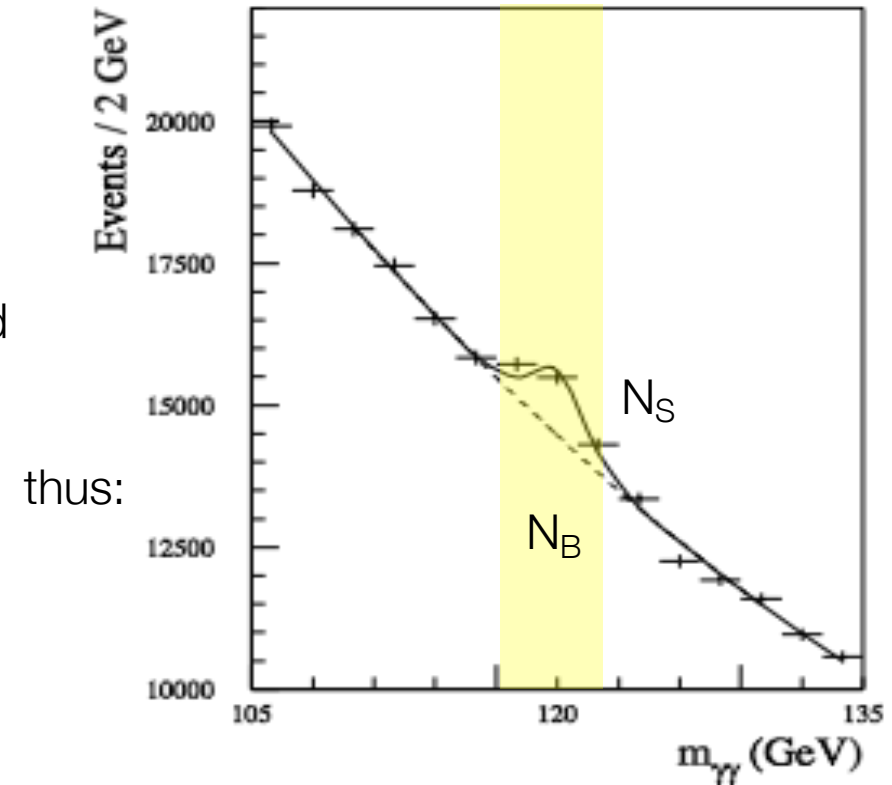
$$S \sim \frac{1}{\sqrt{\sigma_m}}$$

3. Recorded luminosity  $\mathcal{L}$

Signal:  $N_s \sim \mathcal{L}$

Background:  $N_B \sim \mathcal{L}$

$$S \sim \sqrt{\mathcal{L}}$$

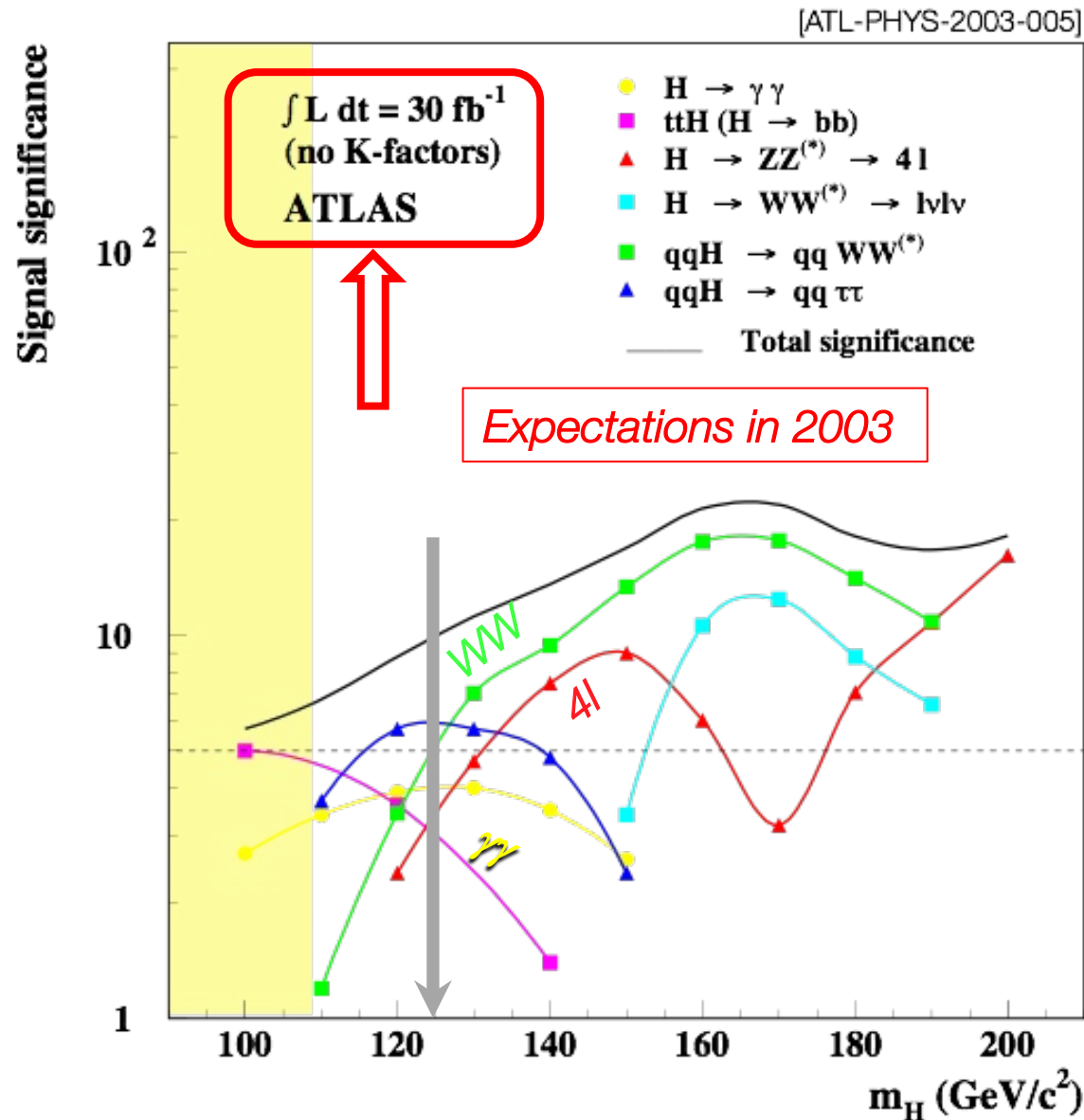


thus:

Decay channel	Mass resolution
$H \rightarrow \gamma\gamma$	1–2%
$H \rightarrow ZZ \rightarrow l^+l^-l'^+l'^-$	1–2%
$H \rightarrow W^+W^- \rightarrow l^+\nu_l l'^-\bar{\nu}_{l'}$	20%
$H \rightarrow b\bar{b}$	10%
$H \rightarrow \tau^+\tau^-$	15%



# (Pre-Discovery) Discovery Potential



## Statement in 2003

- Full mass range can already be covered after a few years at low luminosity
- Several channels available over a large range of masses

• *Low mass discovery requires combination of three of the most demanding channels*

- Comparable situation for the CMS experiment

← Significance =  $5\sigma$   
→ Discovery

Prediction almost correct:

- $\gamma\gamma$ ,  $ZZ$  to 4 leptons,  $WW$  to  $l\nu l\nu$  (higgs to  $qq\tau\tau$  not used)
- Combination of channels



# Higgs Terms in the Lagrangian

$$\mathcal{L} = -g_{Hf\bar{f}}\bar{f}fH + \frac{g_{HHH}}{6}H^3 + \frac{g_{HHHH}}{24}H^4 + \delta_V V_\mu V^\mu \left( g_{HVV}H + \frac{g_{HHVV}}{2}H^2 \right)$$

linear

quadratic

$$g_{Hf\bar{f}} = \frac{m_f}{v}$$

$$g_{HVV} = \frac{2m_V^2}{v}$$

$$g_{HHVV} = \frac{2m_V^2}{v^2}$$

$$g_{HHH} = \frac{3m_H^2}{v}, \quad g_{HHHH} = \frac{3m_H^2}{v^2}$$

$V = W^\pm$  or  $Z$

$\delta W = 1, \delta Z = 1/2$

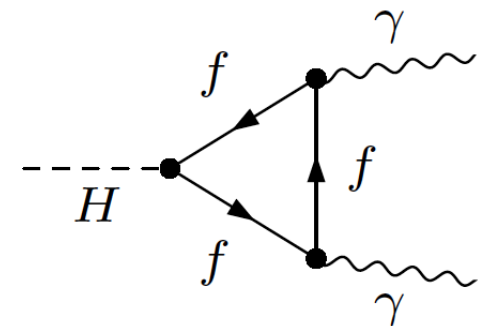
- The dominant mechanisms for Higgs boson production and decay involve the coupling of H to W, Z and/or the third generation quarks and leptons.

Coupling to bosons (W or Z)  $\propto m_V^2$

Coupling to fermions  $\propto m_f$

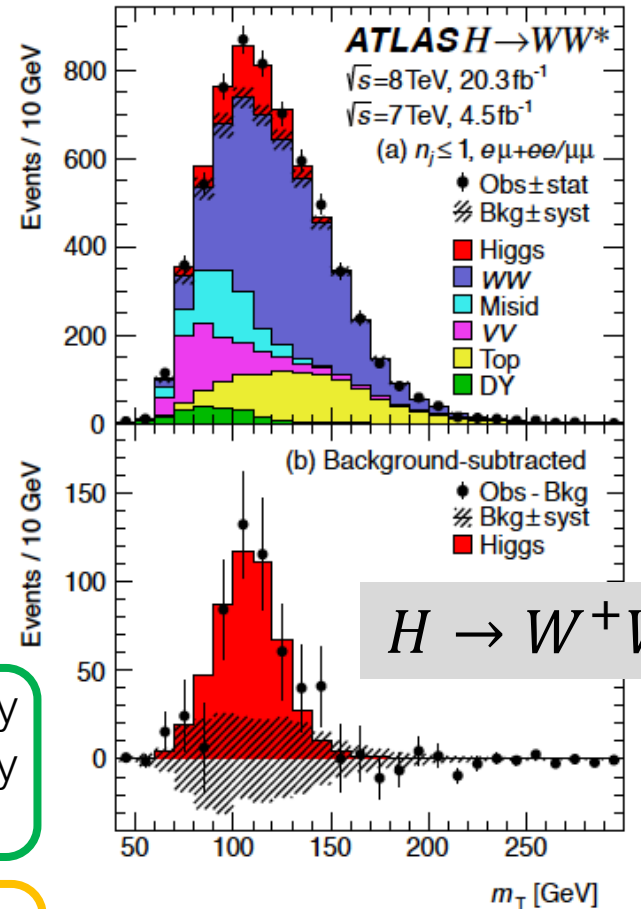
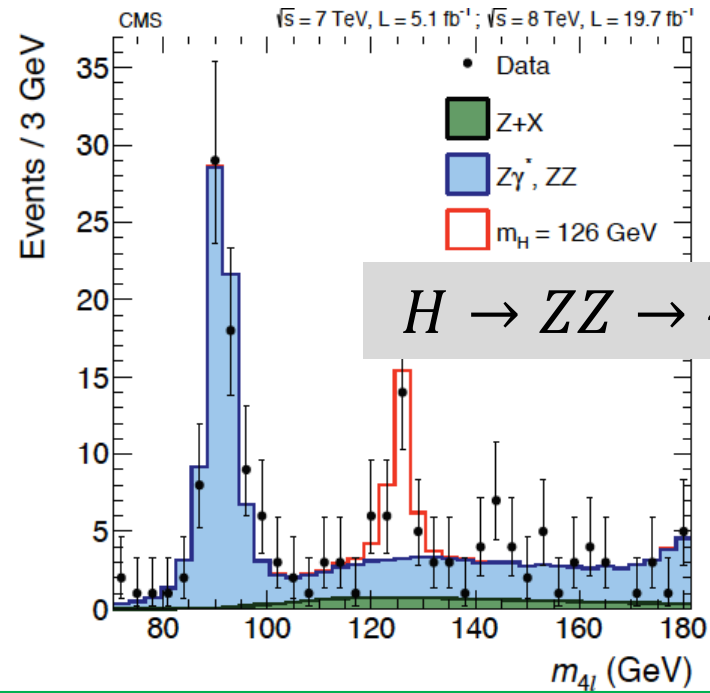
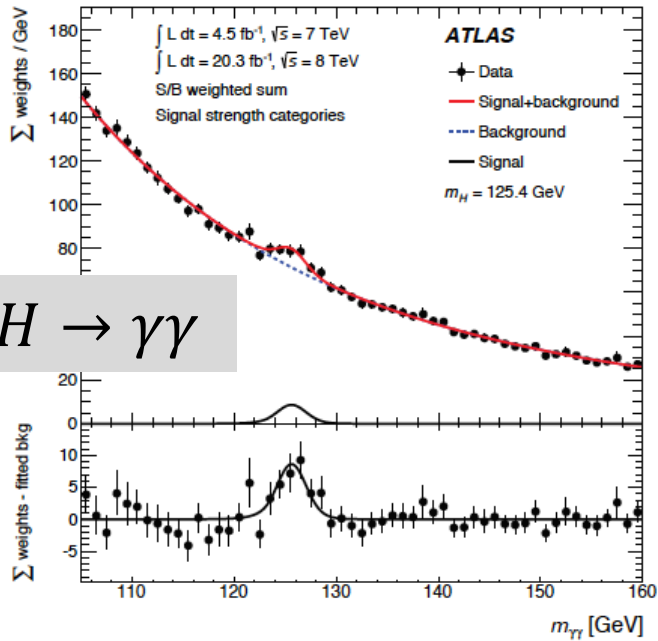
- The Higgs coupling to photons is generated by loops
  - virtual  $W^+W^-$  pair provides the dominant contribution
  - virtual  $tt$  pair is subdominant.

- The Higgs coupling to gluons, is induced by a one-loop graph: H couples to a virtual  $tt$  pair.





# Summary of LHC Run-1 Results (7 TeV + 8 TeV)



• In the  $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ \rightarrow 4l$  channels, all final state particles can be very precisely measured and the reconstructed  $m_H$  resolution is excellent (typically 1-2%).

• the  $H \rightarrow W^+W^- \rightarrow l^+\nu_l l^-\bar{\nu}_l$  channel has relatively large branching fraction, but the  $m_H$  resolution is poor (approximately 20%) due to the presence of neutrinos.

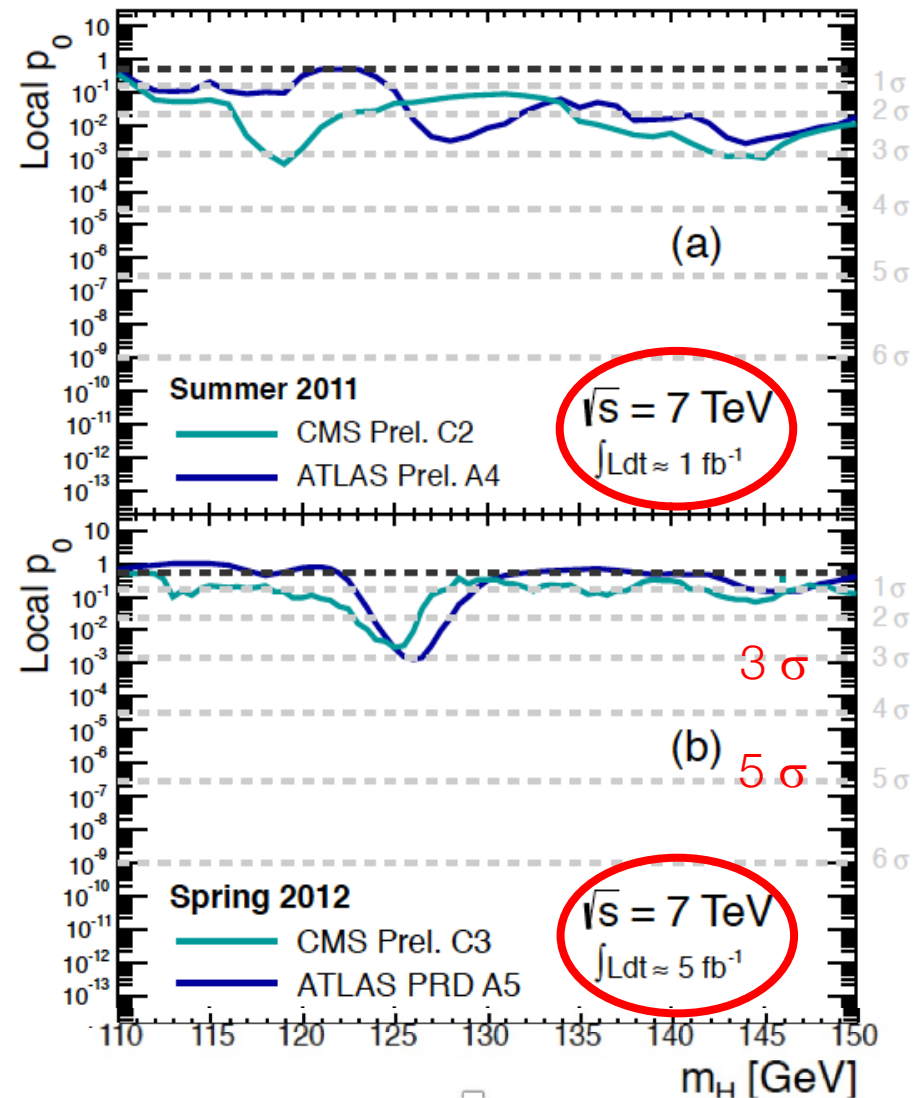


# Past History: a Few Important Facts

- On July 4, 2012 the observation at the LHC of a narrow resonance with a mass of about 125 GeV was announced.
- **Initial strategy: integrate all production modes for one decay channel**
- The observed decay channels indicated  $\rightarrow$  is a boson.
- **Decays rates to  $\gamma\gamma$  and  $ZZ$  consistent with the Standard Model (SM) Higgs boson.**
- There were indications that the new particle also decays to  $W^+W^-$ .

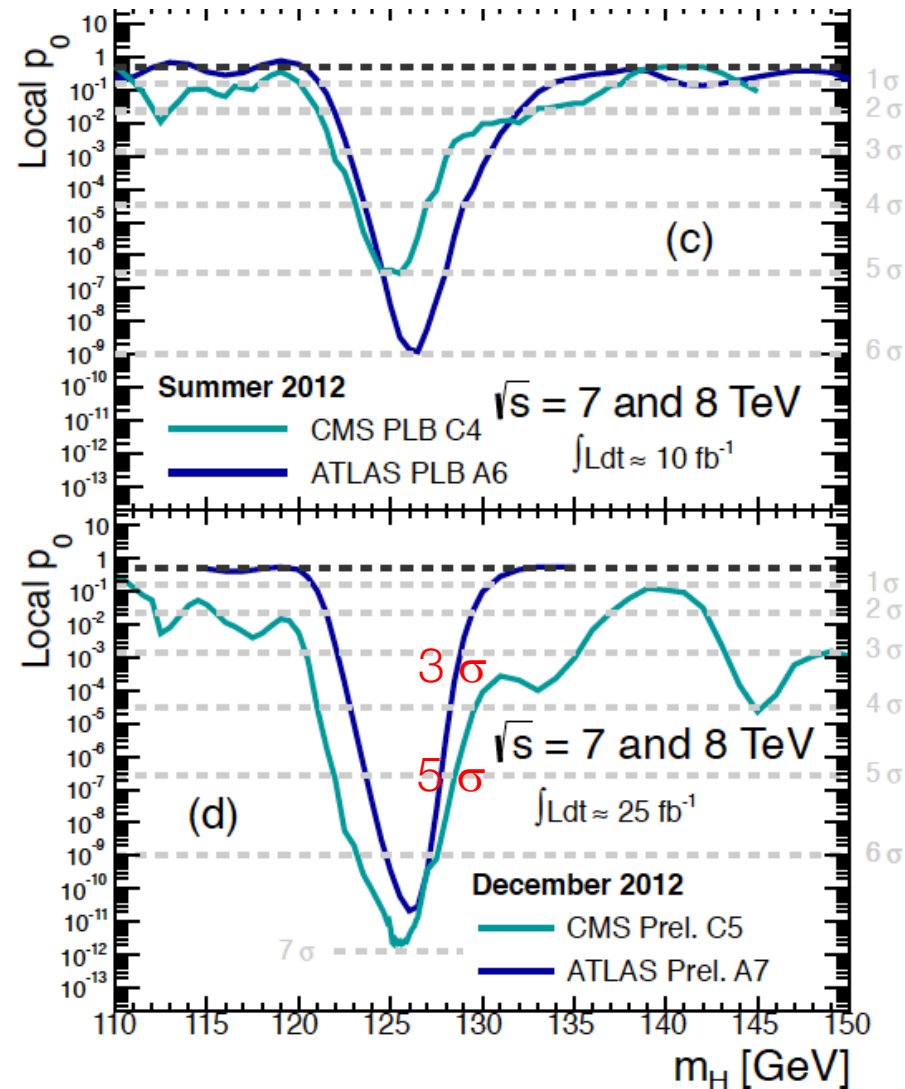
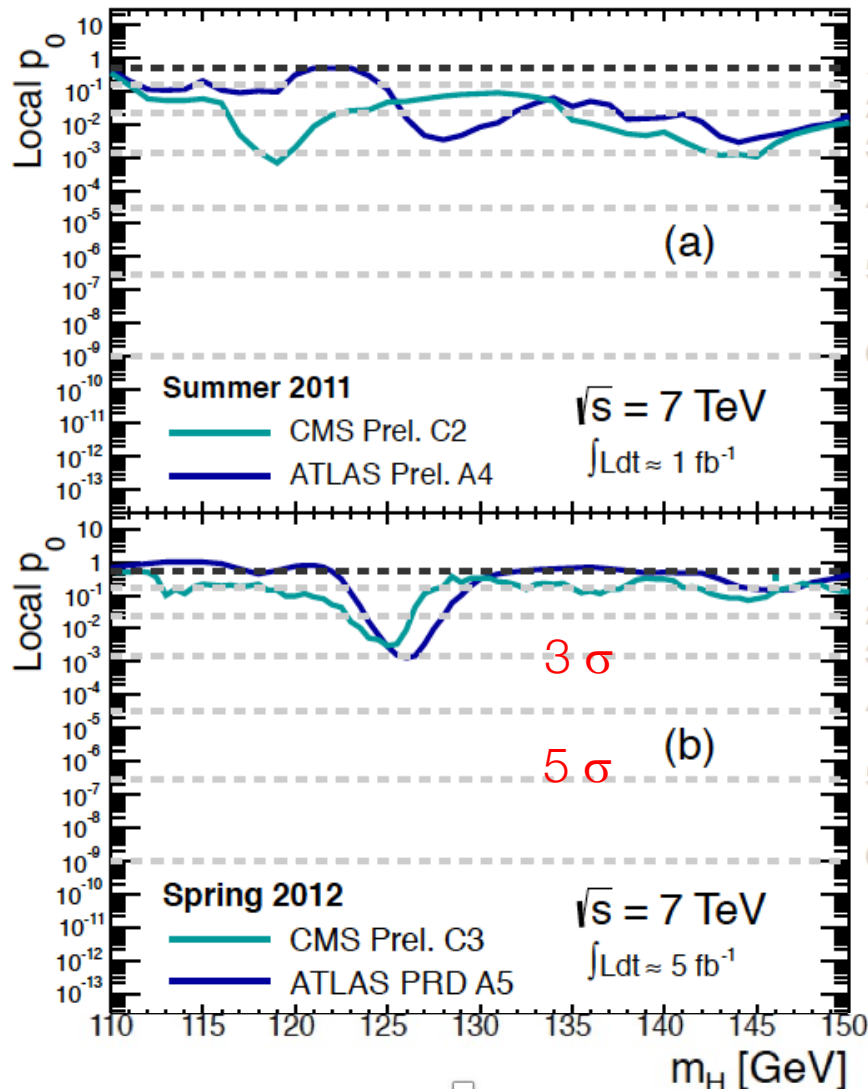
The significance of these observations are quantified by a p-value, the probability for an experiment to give a result compatible with background only.

ATLAS / CMS observed a significance of 5.9 / 4.9  $\sigma$  at a mass  $m_H = 126.5\text{GeV} / 125.5$





# Recent History: a Few Important Facts



$p_0$  = probability that the excess can be described by background only

A  $p_0$  of  $2.87 \times 10^{-7}$  corresponds to  $5\sigma$  excess over the background-only prediction.



# Topologies of Production Mechanisms

ggf, gluon-fusion process :

- largest cross section
- Loop with heavy top quark. *No very distinctive feature in the topology!*

VBF, vector boson fusion :

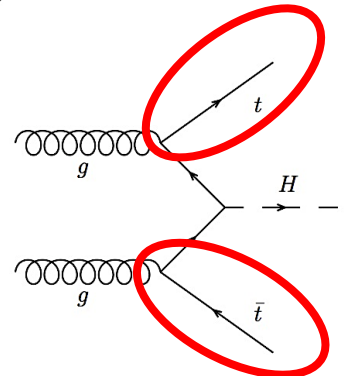
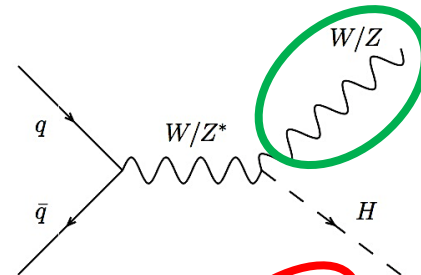
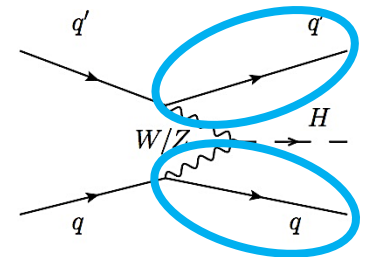
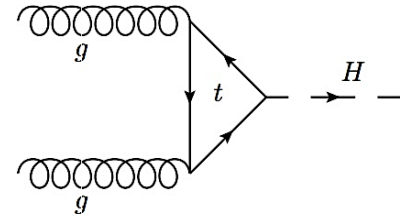
- second-largest cross section
- scattering of  $qq'$  ( $q\bar{q}$ ), mediated by the exchange of a W or Z boson.
- The scattered quarks give two hard jets in the forward and backward regions with a large dijet mass ( $\geq 400\text{GeV}$ ) and separated by  $\Delta\eta_{jj} \geq 3.5 \rightarrow$  one jet very forward + 1 jet very backward.

VH, associated production with W and Z gauge bosons :

- Third cross-section
- W and Z leptonic decay(s)  $\rightarrow$  MET & high  $p_T$  leptons  $\rightarrow$  clean signatures.

ttH: Higgs radiation off top quarks

- High  $p_T$  leptons, MET, b-tagged jets. Complex topology with many decay channels

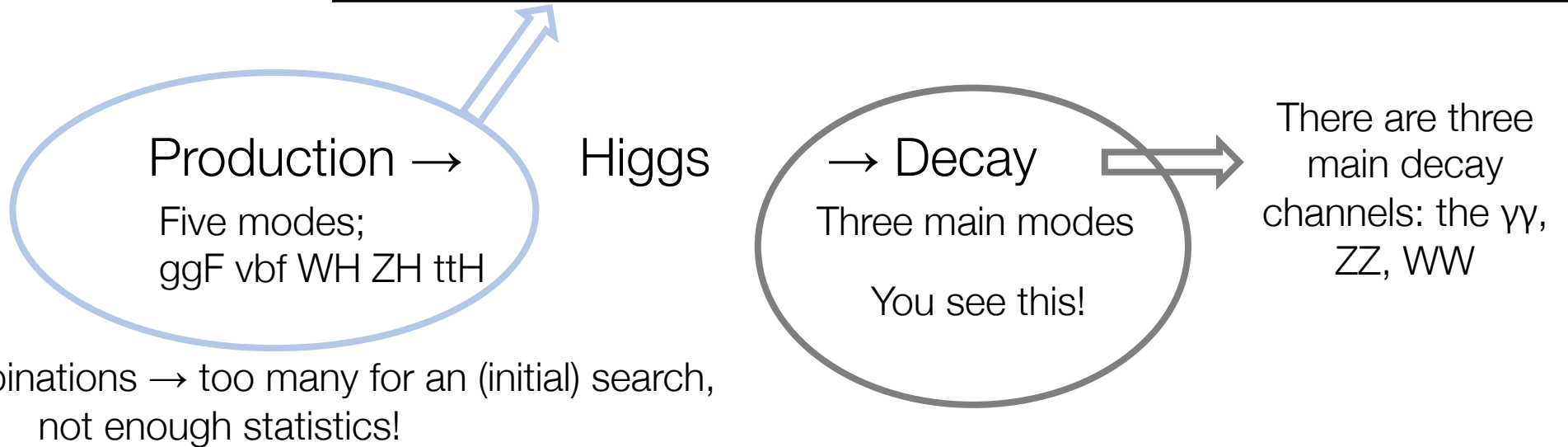




# Improve Initial Discovery

Is the observed resonance the Higgs boson predicted by the SM?  
→ more analysis

Higgses <b>produced</b> in ATLAS & CMS in LHC Run1				
Production Mode	ggF	VBF	WH+ZH	ttH
Higgs events	500.000	40.000	20.000	3.000



Strategy was:

- define categories for each production mode and for each decay mode (→ **check with SM!**)
- optimise the search using characteristics of the (production-decay) topology of the event
- *then integrate over all production modes and give results for 3 decay channels only (WW arrived a bit later)*

Events selected for the discovery were very few: between 1 and 100 for each production-decay category.





# Categorisation

Example:  $H \rightarrow \gamma\gamma$

Several production modes:

1. Consider one production mode at a time (from the most distinct,  $ttH$ , to the least distinct,  $ggf$ )
2. Is the topology / kinematics of that event compatible with that production?

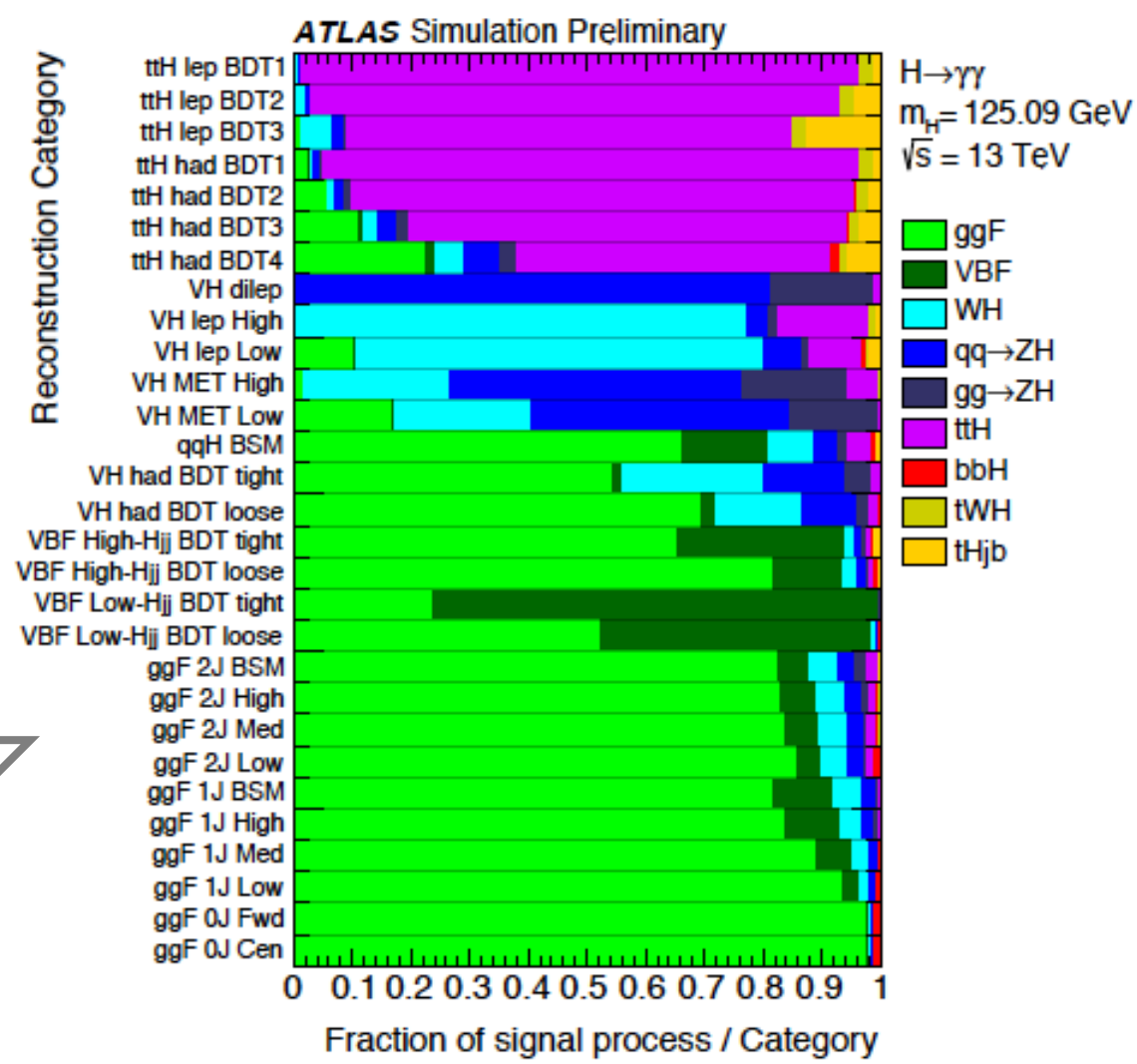
No

Go to next production

Yes

Go to next event

- Attribution to one category is exclusive!
- There is some wrong attribution
- Composition of one category studied by MC





# Of the Categorisation

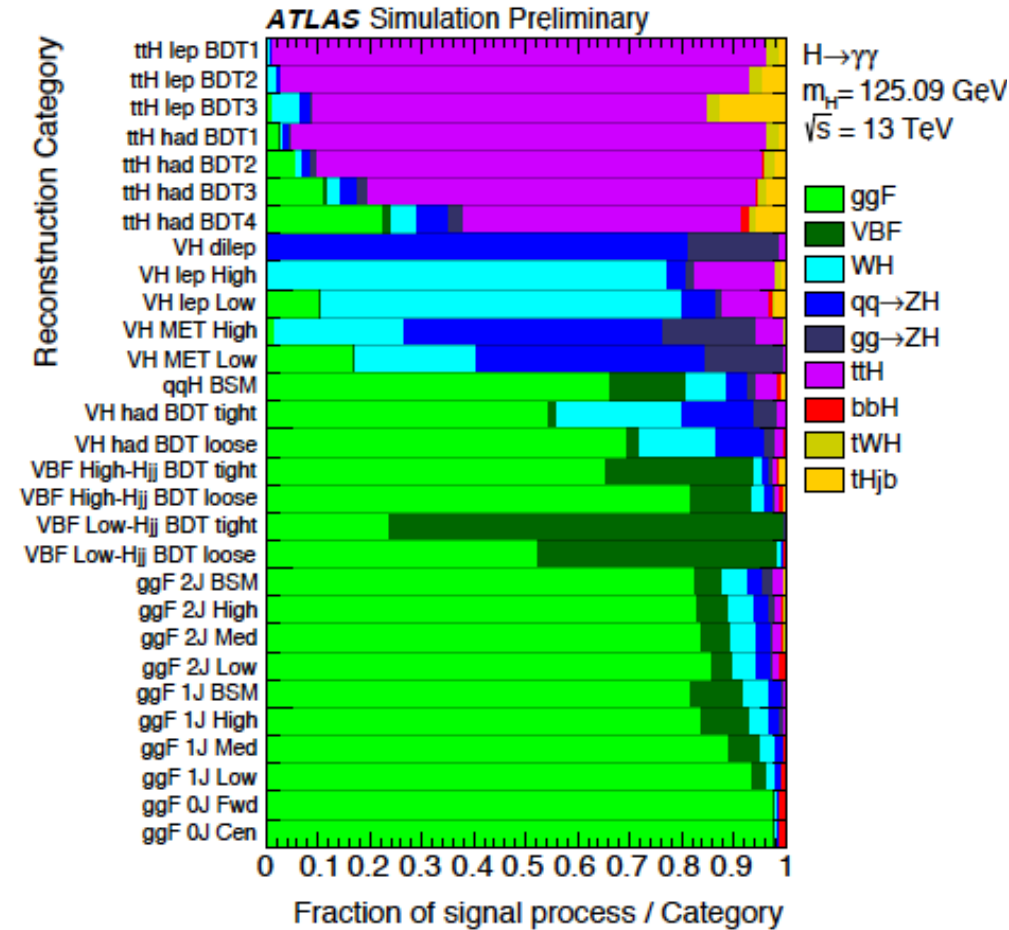
Initial discovery was based on a low sample of events → limited sensitivity to SM Higgs predictions

- Selection: decay modes only, integrate production modes
- **Peak in the mass → a resonance, not necessarily a Higgs**
- Later in time → more statistics → Need to check if also production modes are in agreement with SM

Start with most distinctive and finish with least distinctive

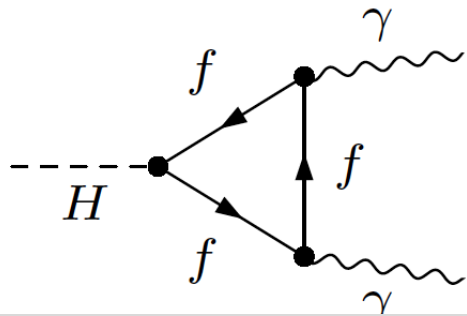
decision ↓

- Separate production processes with topological characteristics → categorisation.
- One category ~ mostly one production mode (but also others) → not a measurement of that production cross-section.
- Simulations are used to determine the relative contributions of the various Higgs production modes in a particular category.

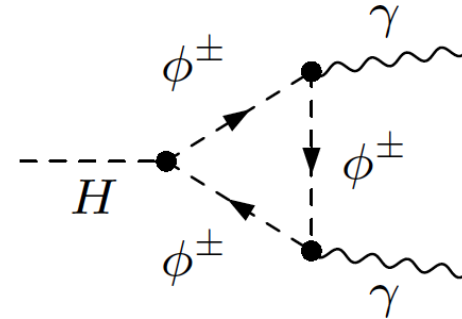
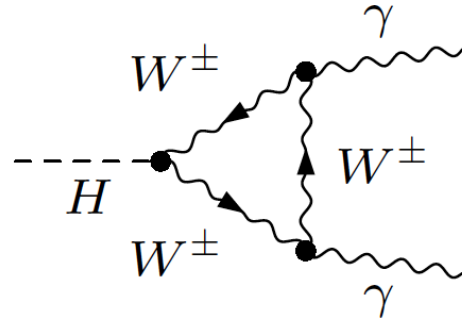




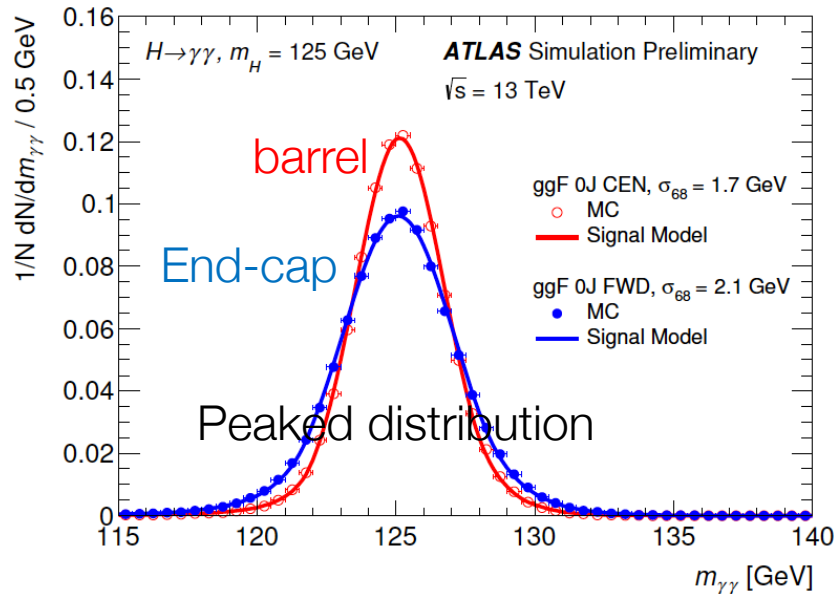
# Higgs Decay to Two Photons



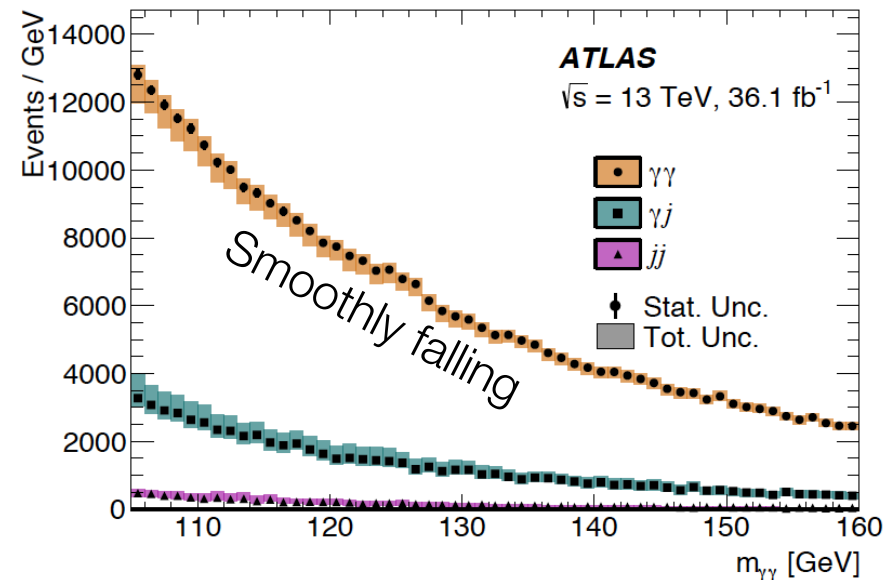
~Only top quarks contribute, contributions from light fermions negligible



**Method:** look for a peak in the invariant mass of two high  $p_T$  photons over a smoothly falling background distribution.



Gaussian central part + power law tails on both sides.



Fits to large control samples of data or simulated background events

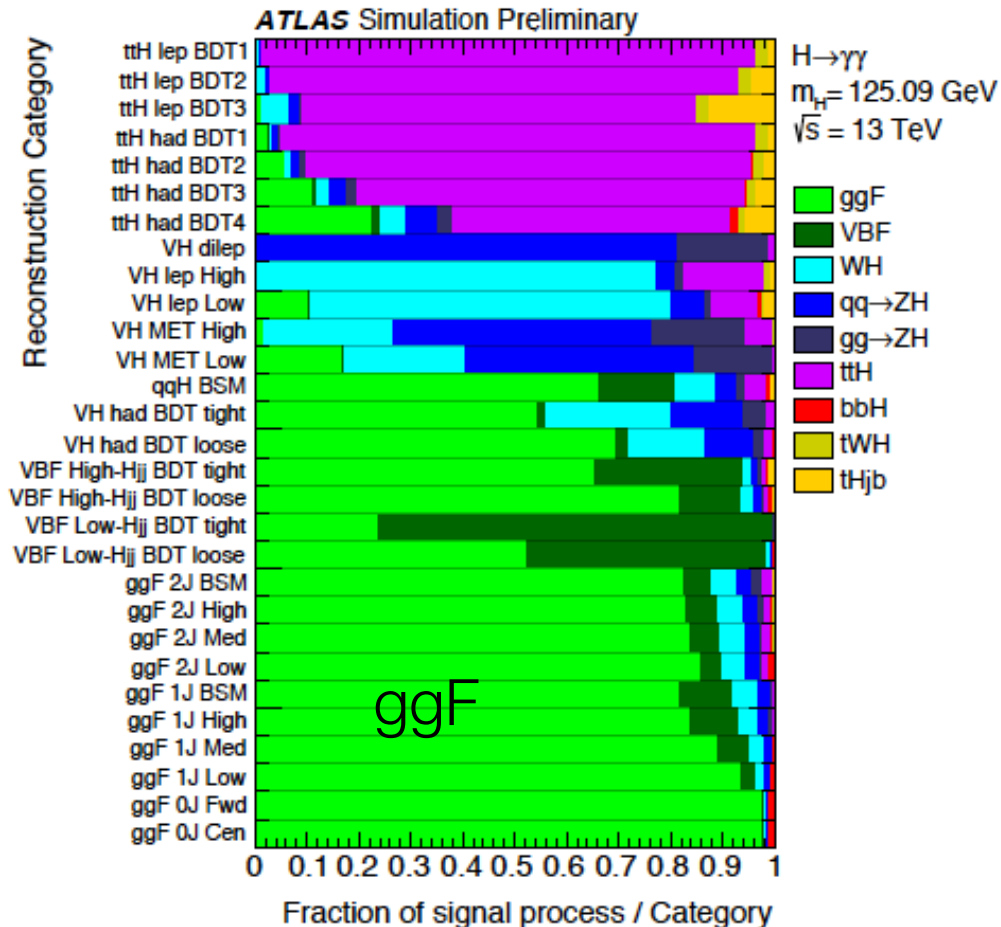


# 29 Categories in Higgs $\rightarrow \gamma\gamma$

Summary of the 29 event reconstruction categories for the measurement of production mode cross sections.

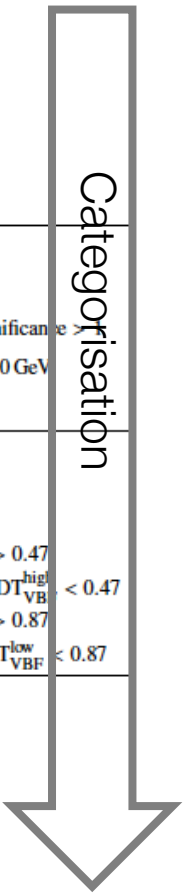
Each event is assigned to the first category whose requirements are satisfied, using the descending order given in the table.

As a result, the event populations of categories are mutually exclusive.



Category label	Selection
ttH lep BDT1	$N_{lep} \geq 1, N_{b-jet} \geq 1, BDT_{ttHlep} > 0.987$
ttH lep BDT2	$N_{lep} \geq 1, N_{b-jet} \geq 1, 0.942 < BDT_{ttHlep} < 0.987$
ttH lep BDT3	$N_{lep} \geq 1, N_{b-jet} \geq 1, 0.705 < BDT_{ttHlep} < 0.942$
ttH had BDT1	$N_{lep} = 0, N_{jets} \geq 3, N_{b-jet} \geq 1, BDT_{ttHhad} > 0.996$
ttH had BDT2	$N_{lep} = 0, N_{jets} \geq 3, N_{b-jet} \geq 1, 0.991 < BDT_{ttHhad} < 0.996$
ttH had BDT3	$N_{lep} = 0, N_{jets} \geq 3, N_{b-jet} \geq 1, 0.971 < BDT_{ttHhad} < 0.991$
ttH had BDT4	$N_{lep} = 0, N_{jets} \geq 3, N_{b-jet} \geq 1, 0.911 < BDT_{ttHhad} < 0.971$
VH dilep	$N_{lep} \geq 2, 70 \text{ GeV} \leq m_{\ell\ell} \leq 110 \text{ GeV}$
VH lep High	$N_{lep} = 1,  m_{e\gamma} - 89 \text{ GeV}  > 5 \text{ GeV}, p_T^{\ell+E_T^{miss}} > 150 \text{ GeV}$
VH lep Low	$N_{lep} = 1,  m_{e\gamma} - 89 \text{ GeV}  > 5 \text{ GeV}, p_T^{\ell+E_T^{miss}} < 150 \text{ GeV}, E_T^{miss} \text{ significance} > 9$
VH MET High	$150 \text{ GeV} < E_T^{miss} < 250 \text{ GeV}, E_T^{miss} \text{ significance} > 9 \text{ or } E_T^{miss} > 250 \text{ GeV}$
VH MET Low	$80 \text{ GeV} < E_T^{miss} < 150 \text{ GeV}, E_T^{miss} \text{ significance} > 8$
qqH BSM	$N_{jets} \geq 2, p_{T,j1} > 200 \text{ GeV}$
VH had BDT tight	$60 \text{ GeV} < m_{ij} < 120 \text{ GeV}, BDT_{VH} > 0.78$
VH had BDT loose	$60 \text{ GeV} < m_{ij} < 120 \text{ GeV}, 0.35 < BDT_{VH} < 0.78$
VBF high- $p_T^{HJJ}$ BDT tight	$ \Delta\eta_{JJ}  > 2,  \eta_{\gamma\gamma} - 0.5(\eta_{j1} + \eta_{j2})  < 5, p_T^{HJJ} > 25 \text{ GeV}, BDT_{VBF}^{high} > 0.47$
VBF high- $p_T^{HJJ}$ BDT loose	$ \Delta\eta_{JJ}  > 2,  \eta_{\gamma\gamma} - 0.5(\eta_{j1} + \eta_{j2})  < 5, p_T^{HJJ} > 25 \text{ GeV}, -0.32 < BDT_{VBF}^{high} < 0.47$
VBF low- $p_T^{HJJ}$ BDT tight	$ \Delta\eta_{JJ}  > 2,  \eta_{\gamma\gamma} - 0.5(\eta_{j1} + \eta_{j2})  < 5, p_T^{HJJ} < 25 \text{ GeV}, BDT_{VBF}^{low} > 0.87$
VBF low- $p_T^{HJJ}$ BDT loose	$ \Delta\eta_{JJ}  > 2,  \eta_{\gamma\gamma} - 0.5(\eta_{j1} + \eta_{j2})  < 5, p_T^{HJJ} < 25 \text{ GeV}, 0.26 < BDT_{VBF}^{low} < 0.87$
ggF 2J BSM	$N_{jets} \geq 2, p_T^{\gamma\gamma} \geq 200 \text{ GeV}$
ggF 2J High	$N_{jets} \geq 2, p_T^{\gamma\gamma} \in [120, 200] \text{ GeV}$
ggF 2J Med	$N_{jets} \geq 2, p_T^{\gamma\gamma} \in [60, 120] \text{ GeV}$
ggF 2J Low	$N_{jets} \geq 2, p_T^{\gamma\gamma} \in [0, 60] \text{ GeV}$
ggF 1J BSM	$N_{jets} = 1, p_T^{\gamma\gamma} \geq 200 \text{ GeV}$
ggF 1J High	$N_{jets} = 1, p_T^{\gamma\gamma} \in [120, 200] \text{ GeV}$
ggF 1J Med	$N_{jets} = 1, p_T^{\gamma\gamma} \in [60, 120] \text{ GeV}$
ggF 1J Low	$N_{jets} = 1, p_T^{\gamma\gamma} \in [0, 60] \text{ GeV}$
ggF 0J Fwd	$N_{jets} = 0, \text{ one photon with }  \eta  > 0.95$
ggF 0J Cen	$N_{jets} = 0, \text{ two photons with }  \eta  \leq 0.95$

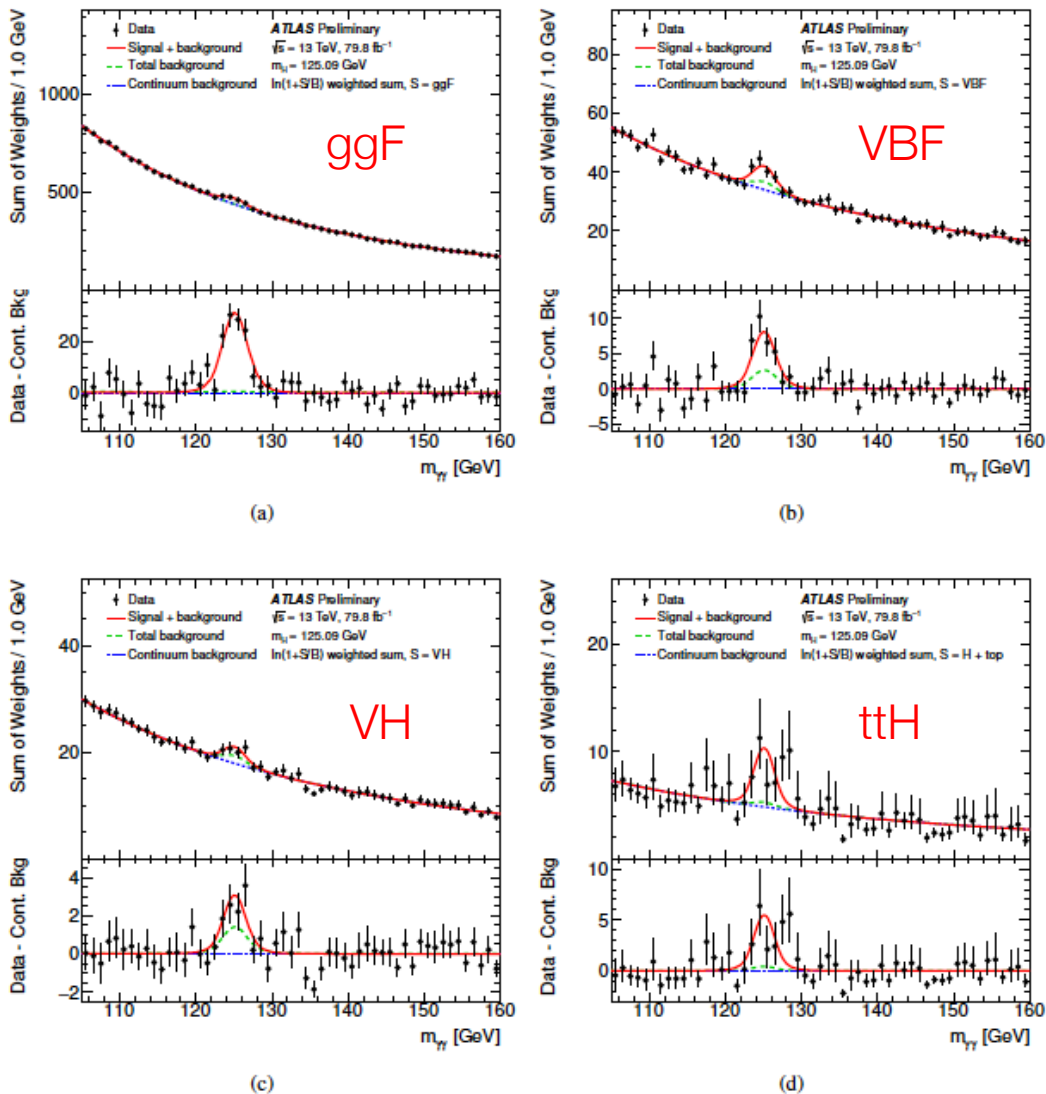
Categorisation





# The $m_{\gamma\gamma}$ distribution with $\sim 80\text{fb}^{-1}$

$m_{\gamma\gamma}$  in the four production categories



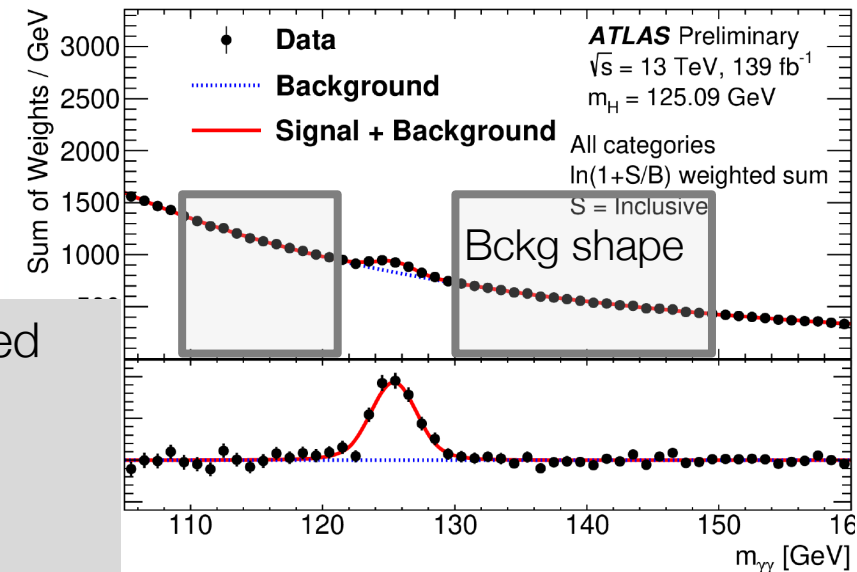
Cumulated  $m_{\gamma\gamma}$  distribution for the  $H \rightarrow \gamma\gamma$  decay

Signal-strength is defined as the ratio between observed and the SM expected cross section

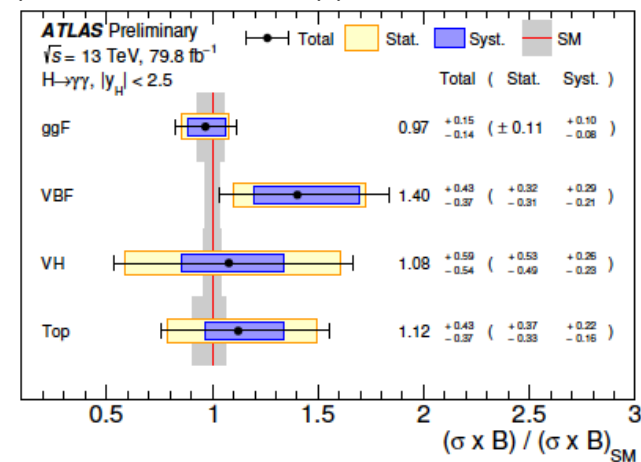
$$\mu = \frac{\sigma^{\text{measured}} \times BR^{\text{measured}}}{\sigma^{\text{SM}} \times BR^{\text{SM}}}$$

This observable tells how well data are in agreement with SM

Signal consists of between 150 and 200  $H \rightarrow \gamma\gamma$  decay candidates



$\mu$  for the  $H \rightarrow \gamma\gamma$  decay channel



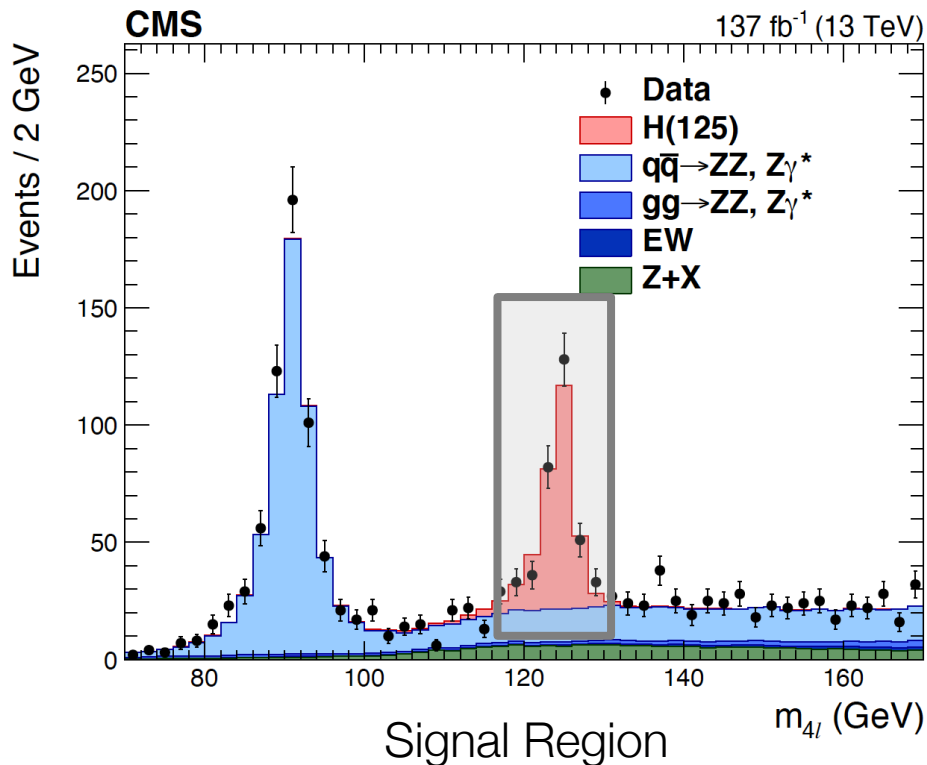


# Higgs decay to $ZZ \rightarrow 4 \text{ leptons}$ ( $80 \text{ fb}^{-1}$ )

Method:  $H \rightarrow ZZ^* \rightarrow l^+ l^- l'^+ l'^-$  look for a narrow mass peak over a continuous background.

$$H \rightarrow ZZ^* \rightarrow 4\ell \text{ decay } (4\mu, 2e2\mu, 2\mu2e, 4e)$$

Different event-observables  $\rightarrow$  the probability for the event to be signal-like or background-like  $\rightarrow$  Neural Network  
 Event selection to magnify the ratio S/B



Number of expected and observed events in the four decay channels after the event selection, in the mass range  $115 \text{ GeV} < m_{4l} < 130 \text{ GeV}$ .

The sum of the expected number of SM Higgs boson events and the estimated background yields is compared to the data.

Final state	Signal	$ZZ^*$ background	Other backgrounds	Total expected	Observed
$4\mu$	$40.5 \pm 1.7$	$19.0 \pm 1.1$	$1.71 \pm 0.10$	$61.2 \pm 2.0$	64
$2e2\mu$	$28.2 \pm 1.2$	$13.3 \pm 0.8$	$1.38 \pm 0.10$	$42.8 \pm 1.4$	64
$2\mu2e$	$22.1 \pm 1.4$	$9.2 \pm 0.9$	$2.99 \pm 0.09$	$34.3 \pm 1.7$	39
$4e$	$21.1 \pm 1.4$	$8.6 \pm 0.8$	$2.90 \pm 0.09$	$32.5 \pm 1.6$	28
<b>Total</b>	<b><math>112 \pm 5</math></b>	<b><math>50 \pm 4</math></b>	<b><math>8.96 \pm 0.12</math></b>	<b><math>171 \pm 6</math></b>	<b>195</b>

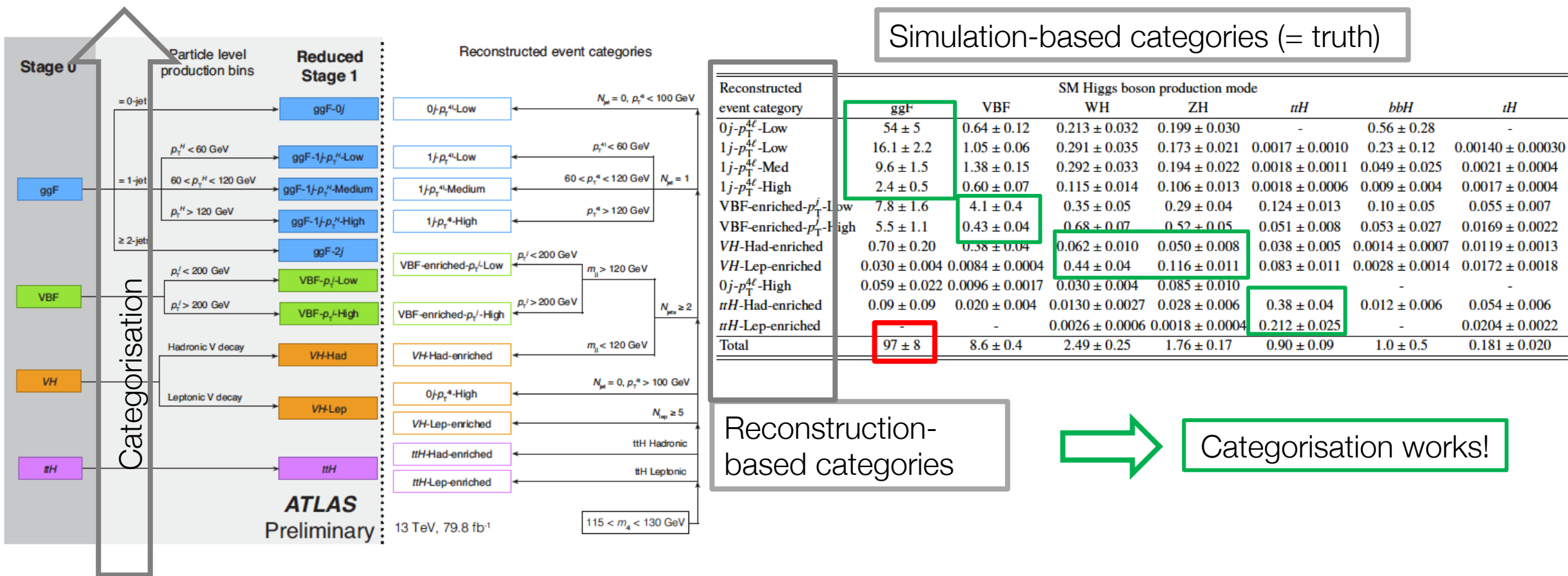
Agreement data/MC to  $1.7 \sigma$



# Composition of the Signal

A categorisation ~similar to what was used for the  $H \rightarrow \gamma\gamma$  decay is also used for the  $H \rightarrow 4l$  decay

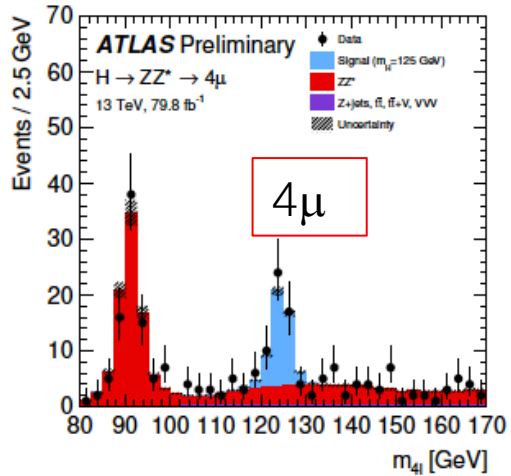
Simulated signal composition for  $\sim 80 \text{ fb}^{-1}$  of luminosity at 13 TeV: The **ggF** component, as expected, dominates.



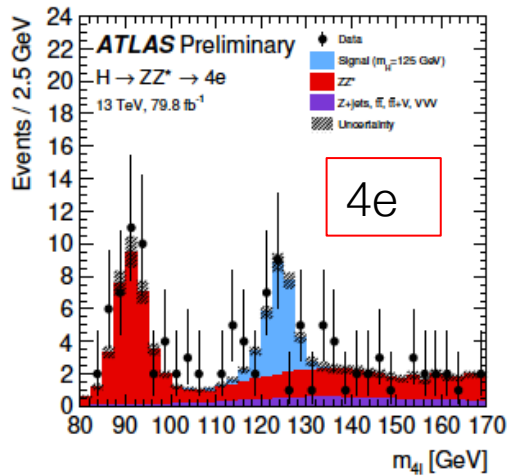
A correspondence is determined between reconstruction and simulation categories



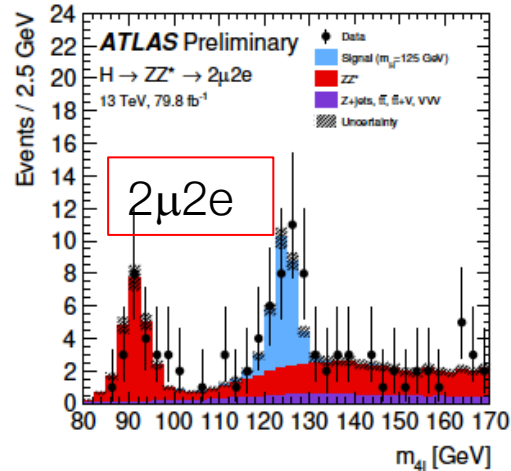
# Details of the Results



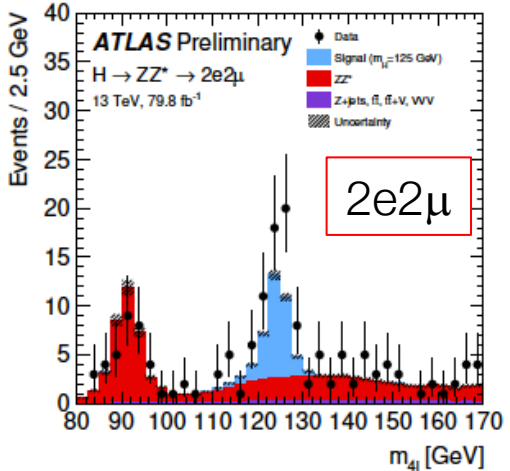
(a)



(b)



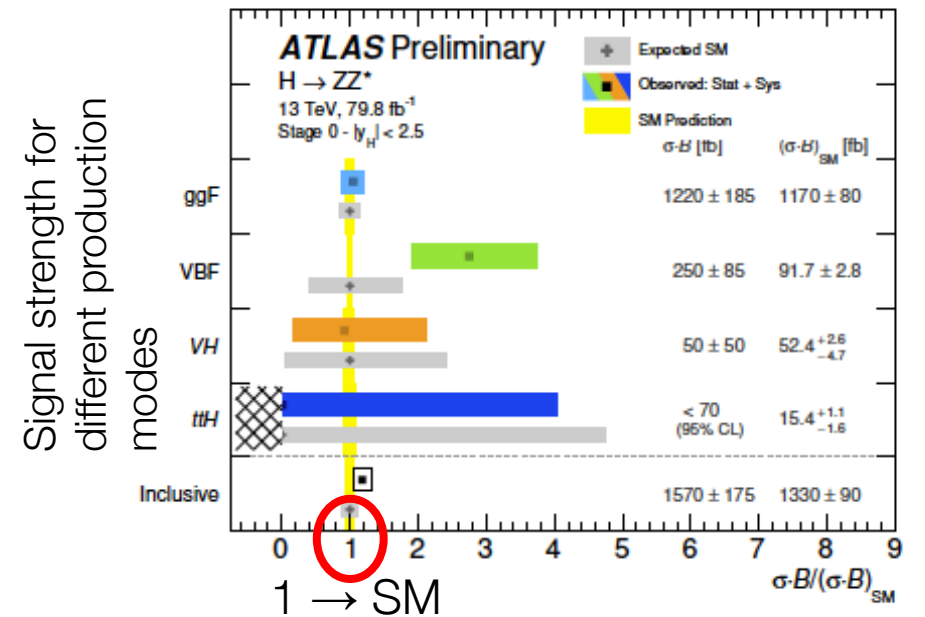
(c)



(d)

Cross section [fb]	Data ( $\pm$ (stat.) $\pm$ (syst.) )	Standard Model prediction	$p$ -value [%]
$\sigma_{4\mu}$	0.97 $\pm$ 0.17 $\pm$ 0.05	0.886 $\pm$ 0.039	62
$\sigma_{4e}$	0.61 $\pm$ 0.21 $\pm$ 0.07	0.886 $\pm$ 0.039	25
$\sigma_{2\mu 2e}$	0.88 $\pm$ 0.21 $\pm$ 0.08	0.786 $\pm$ 0.035	66
$\sigma_{2e 2\mu}$	1.37 $\pm$ 0.22 $\pm$ 0.07	0.786 $\pm$ 0.035	0.3

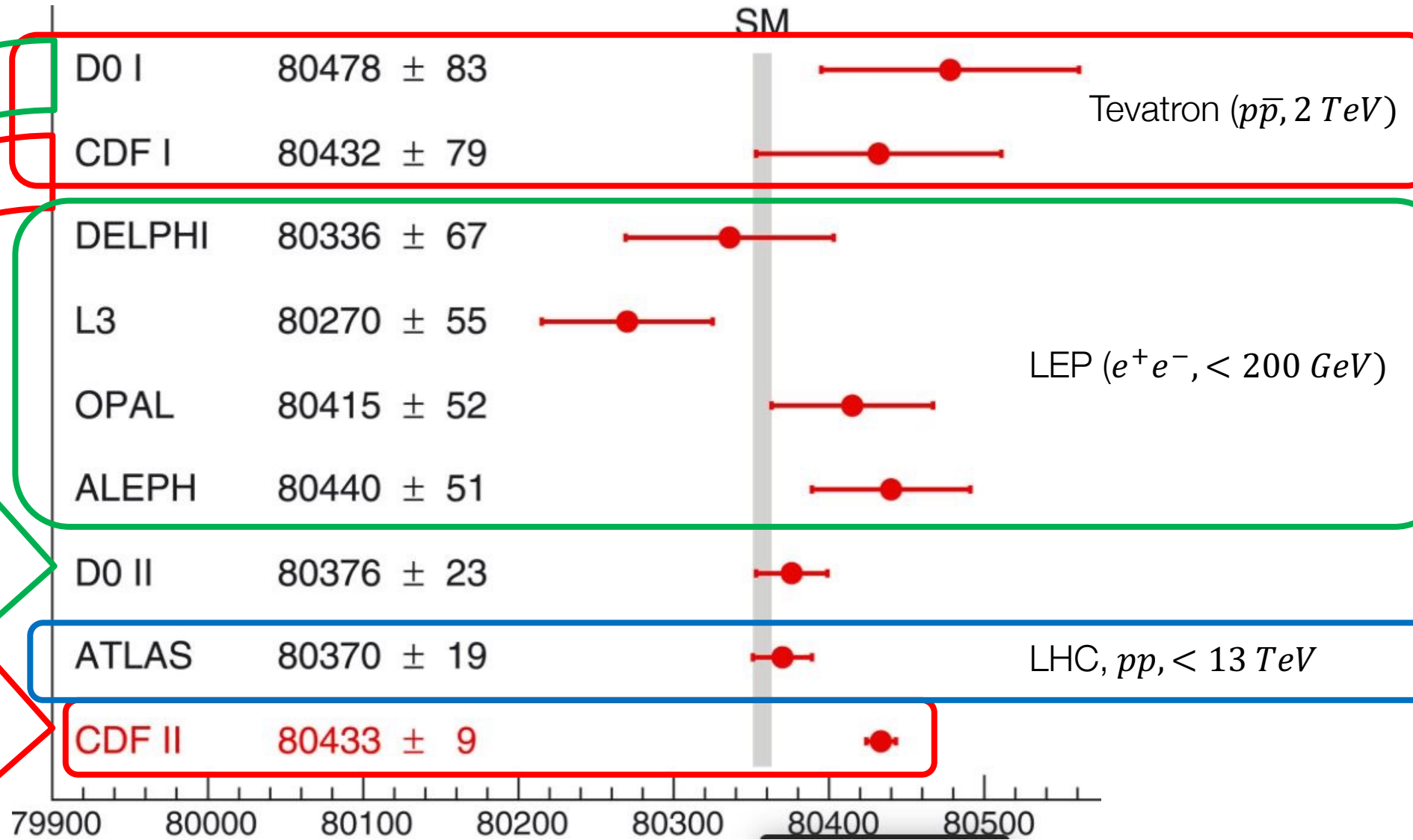
The signal strength  $\mu$  has also been calculated:  
 Integrated  $\mu = 1.19 \pm 0.12(\text{stat.}) \pm 0.06(\text{exp.})_{-0.07}^{+0.08}(\text{th.}) = 1.19_{-0.15}^{+0.16}$





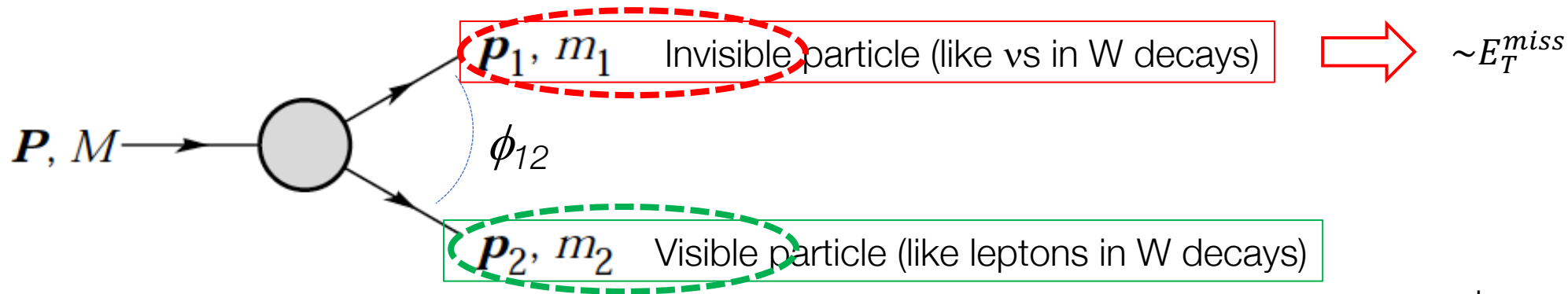


# *W mass measurement at Colliders*





# W Mass Measurements at Hadron Colliders



The mass of the parent particle can be constrained with the observable  $M_T$  defined by

$$M_T^2 \equiv [E_T(1) + E_T(2)]^2 - [p_T(1) + p_T(2)]^2$$

$$= m_1^2 + m_2^2 + 2[E_T(1)E_T(2) - p_T(1) \cdot p_T(2)]$$

where

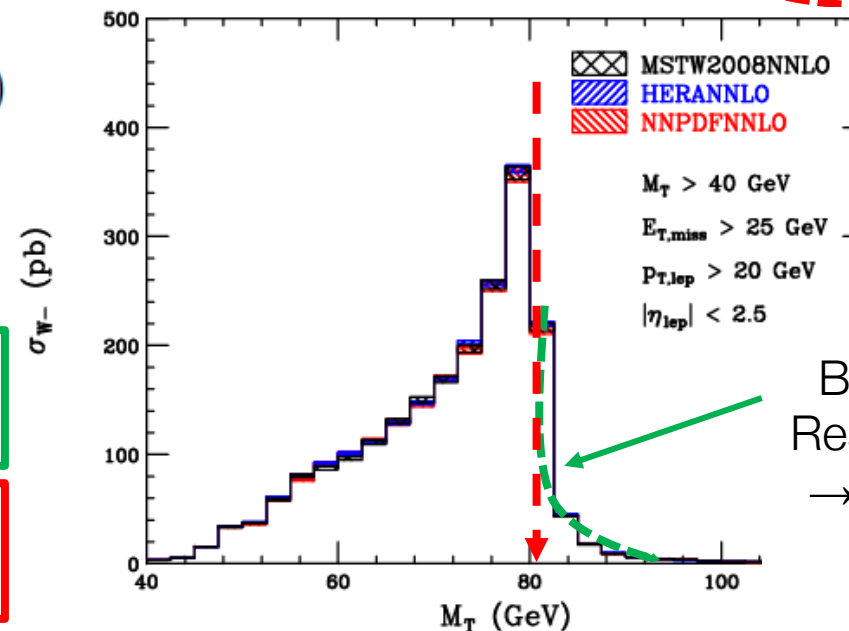
$$p_T(1) = E_T^{miss}$$

For  $m_1 \sim m_2 \sim 0 \rightarrow M_T^2 = 2|p_T(1)||p_T(2)|(1 - \cos \phi_{12})$

Important characteristic: the end point of this distribution is

$$M_T^{max} = M$$

Also the distribution of the  $p_T$  of the lepton has memory of  $m_W$ : the end-point is  $m_W/2$

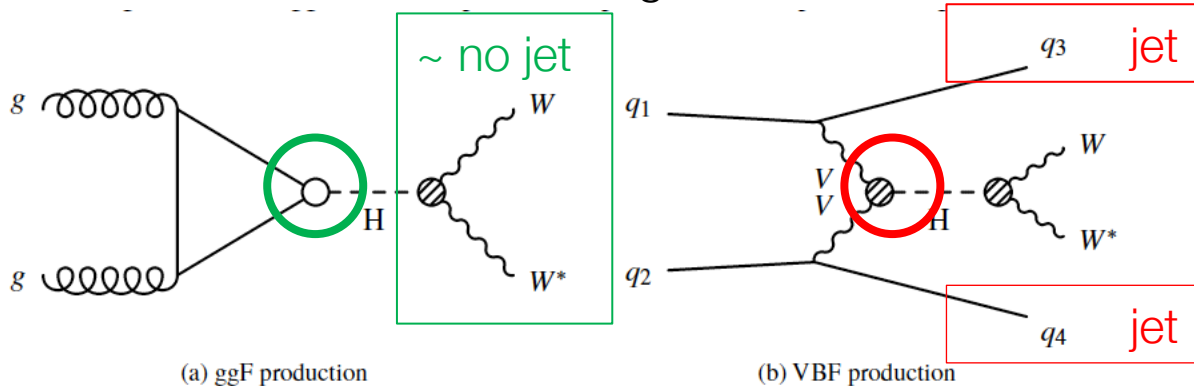


Breit-Wigner+ Resolution effect  
 $\rightarrow$  sharp fall  $\rightarrow$   
 $\rightarrow$  smooth fall



# Higgs $\rightarrow$ $WW \rightarrow e\nu \mu\nu$ ( $36 \text{ fb}^{-1}$ )

## Dominant diagrams



Coupling to fermions

Coupling to bosons

## Event selection

2 OFOS leptons ( $e\nu_e \mu\nu_\mu$ )  $\rightarrow$  no Z decay!

Category	$N_{\text{jet},(p_T>30 \text{ GeV})} = 0$ ggF	$N_{\text{jet},(p_T>30 \text{ GeV})} = 1$ ggF	$N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2$ VBF
Preselection	Two isolated, different-flavour leptons ( $\ell = e, \mu$ ) with opposite charge $p_T^{\text{lead}} > 22 \text{ GeV}, p_T^{\text{sublead}} > 15 \text{ GeV}$ $m_{\ell\ell} > 10 \text{ GeV}$ $p_T^{\text{miss}} > 20 \text{ GeV}$		
Background rejection	$\Delta\phi(\ell\ell, E_T^{\text{miss}}) > \pi/2$ $p_T^{\ell\ell} > 30 \text{ GeV}$	$\max(m_T^\ell) > 50 \text{ GeV}$ $m_{\tau\tau} < m_Z - 25 \text{ GeV}$	$N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$
$H \rightarrow WW^* \rightarrow e\nu\mu\nu$ topology	$m_{\ell\ell} < 55 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 1.8$	central jet veto outside lepton veto	
Discriminant variable	$m_T$		
BDT input variables	BDT $m_{jj}, \Delta y_{jj}, m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T, \sum_\ell C_\ell, \sum_{\ell,j} m_{\ell j}, p_T^{\text{tot}}$		

Background is computed using simulation. Control regions in data (orthogonal to the signal region) are used to normalise the MC predictions for most important backgrounds:

- Non resonant WW
- Top pairs production
- Di-bosons (WZ and ZZ) and Drell-Yan

CR	$N_{\text{jet},(p_T>30 \text{ GeV})} = 0$ ggF	$N_{\text{jet},(p_T>30 \text{ GeV})} = 1$ ggF	$N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2$ VBF
WW	$55 < m_{\ell\ell} < 110 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 2.6$	$m_{\ell\ell} > 80 \text{ GeV}$ $ m_{\tau\tau} - m_Z  > 25 \text{ GeV}$ $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ $\max(m_T^\ell) > 50 \text{ GeV}$	
$t\bar{t}/Wt$	$N_{b\text{-jet},(20 \text{ GeV}<p_T<30 \text{ GeV})} > 0$ $\Delta\phi(\ell\ell, E_T^{\text{miss}}) > \pi/2$ $p_T^{\ell\ell} > 30 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 2.8$	$N_{b\text{-jet},(p_T>30 \text{ GeV})} = 1$ $N_{b\text{-jet},(20 \text{ GeV}<p_T<30 \text{ GeV})} = 0$ $\max(m_T^\ell) > 50 \text{ GeV}$ $m_{\tau\tau} < m_Z - 25 \text{ GeV}$	$N_{b\text{-jet},(p_T>20 \text{ GeV})} = 1$ central jet veto outside lepton veto
$Z/\gamma^*$	$\Delta\phi_{\ell\ell} > 2.8$	no $p_T^{\text{miss}}$ requirement $\max(m_T^\ell) > 50 \text{ GeV}$ $m_{\tau\tau} > m_Z - 25 \text{ GeV}$	$N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ $m_{\ell\ell} < 80 \text{ GeV}$ central jet veto outside lepton veto $ m_{\tau\tau} - m_Z  \leq 25 \text{ GeV}$

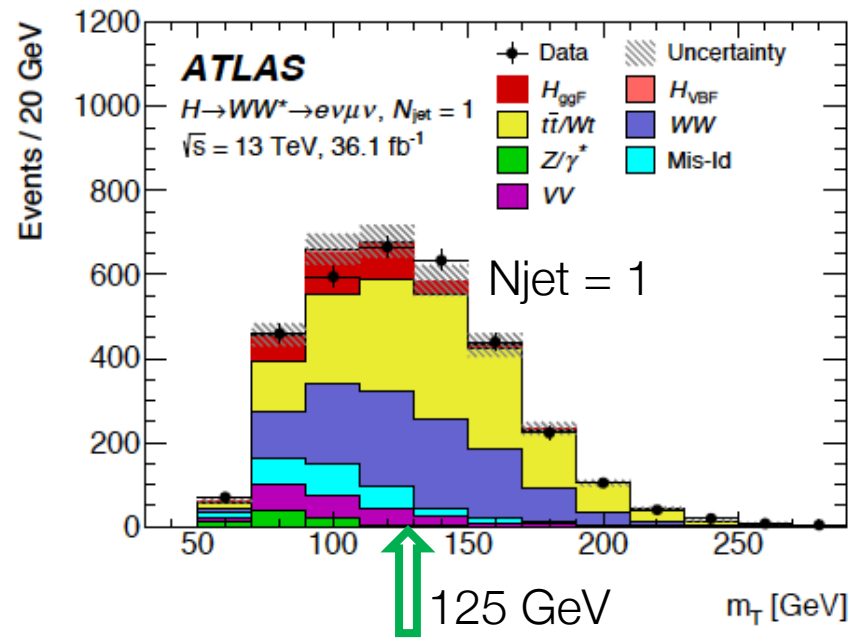
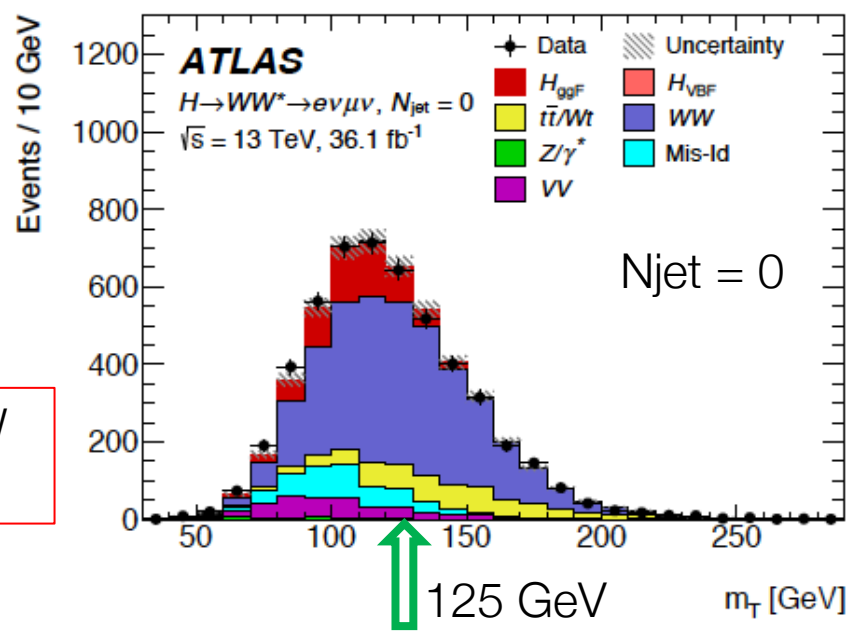


# Results: $m_T$ as proxy of $m_H$

Process	$N_{\text{jet}} = 0$ ggF	$N_{\text{jet}} = 1$ ggF	$N_{\text{jet}} \geq 2$ VBF	
			Inclusive	BDT: [0.86, 1.0]
$H_{\text{ggF}}$	$639 \pm 110$	$285 \pm 51$	$42 \pm 16$	$6 \pm 3$
$H_{\text{VBF}}$	$7 \pm 1$	$31 \pm 2$	$28 \pm 16$	$16 \pm 6$
$WW$	$3016 \pm 203$	$1053 \pm 206$	$400 \pm 60$	$11 \pm 2$
$VV$	$333 \pm 38$	$208 \pm 32$	$70 \pm 12$	$3 \pm 1$
$t\bar{t}/Wt$	$588 \pm 130$	$1397 \pm 179$	$1270 \pm 80$	$14 \pm 2$
Mis-Id	$447 \pm 77$	$234 \pm 49$	$90 \pm 30$	$6 \pm 2$
$Z/\gamma^*$	$27 \pm 11$	$76 \pm 24$	$280 \pm 40$	$4 \pm 1$
Total	$5067 \pm 80$	$3296 \pm 61$	$2170 \pm 50$	$60 \pm 10$
Observed	5089	3264	2164	60

ggF populates mostly the  $N_{\text{jet}}=0$  and  $N_{\text{jet}}=1$  region  
 VBF populates mostly the  $N_{\text{jet}}>1$  region

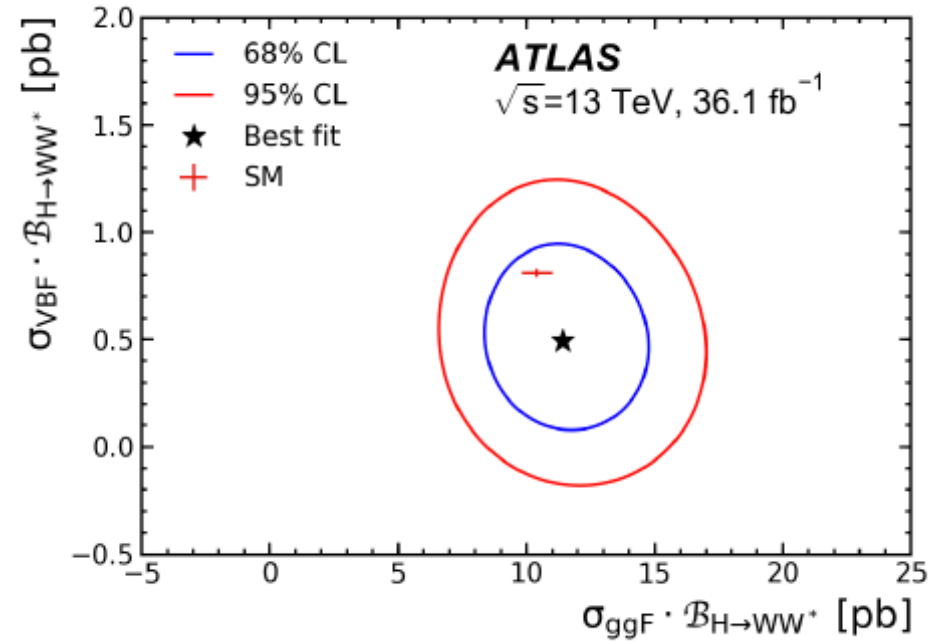
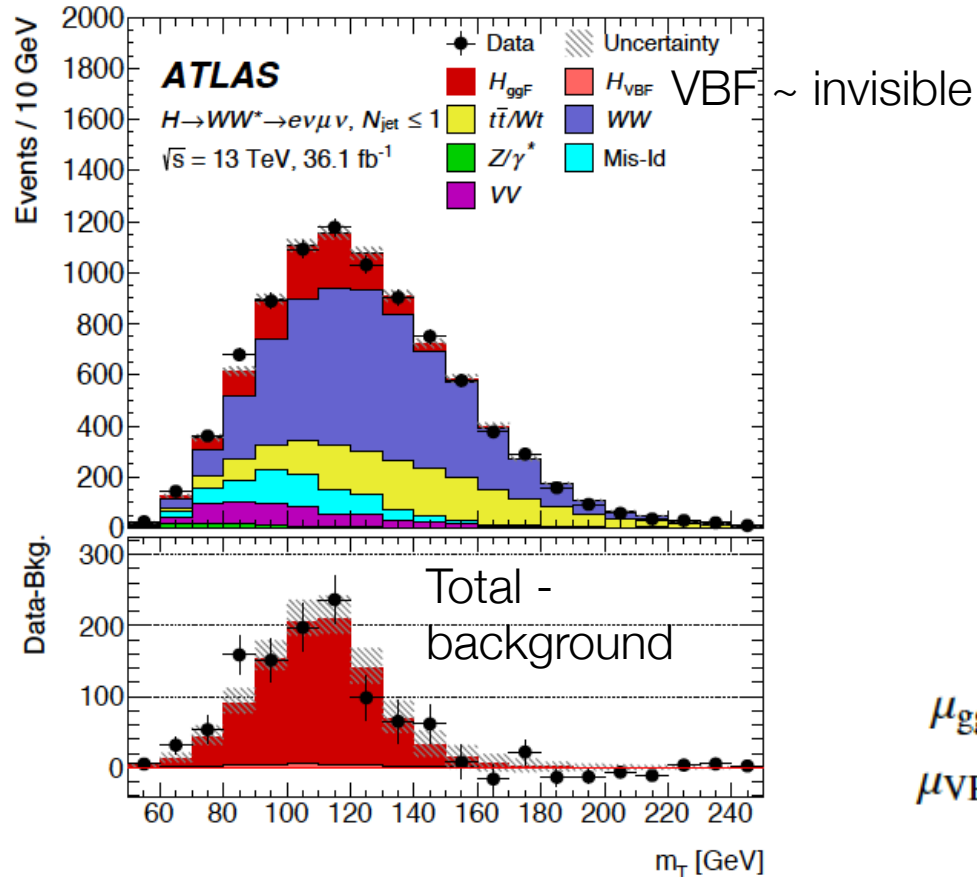
BDT (Boosted Decision Tree) is an analysis technique that combines many different variables in one unique indicator ranging between -1 and 1: the more negative (positive) values are the more background-like (signal-like) is the event





# Higgs to WW: Results

Combined  $m_T$  distribution for  $N_{\text{jet}} < 2$ .



$$\mu_{\text{ggF}} = 1.10^{+0.10}_{-0.09}(\text{stat.})^{+0.13}_{-0.11}(\text{theo syst.})^{+0.14}_{-0.13}(\text{exp syst.}) = 1.10^{+0.21}_{-0.20}$$

$$\mu_{\text{VBF}} = 0.62^{+0.29}_{-0.27}(\text{stat.})^{+0.12}_{-0.13}(\text{theo syst.}) \pm 0.15(\text{exp syst.}) = 0.62^{+0.36}_{-0.35}$$

Difference between the data and the estimated background for a SM Higgs boson with  $m_H = 125 \text{ GeV}$ .

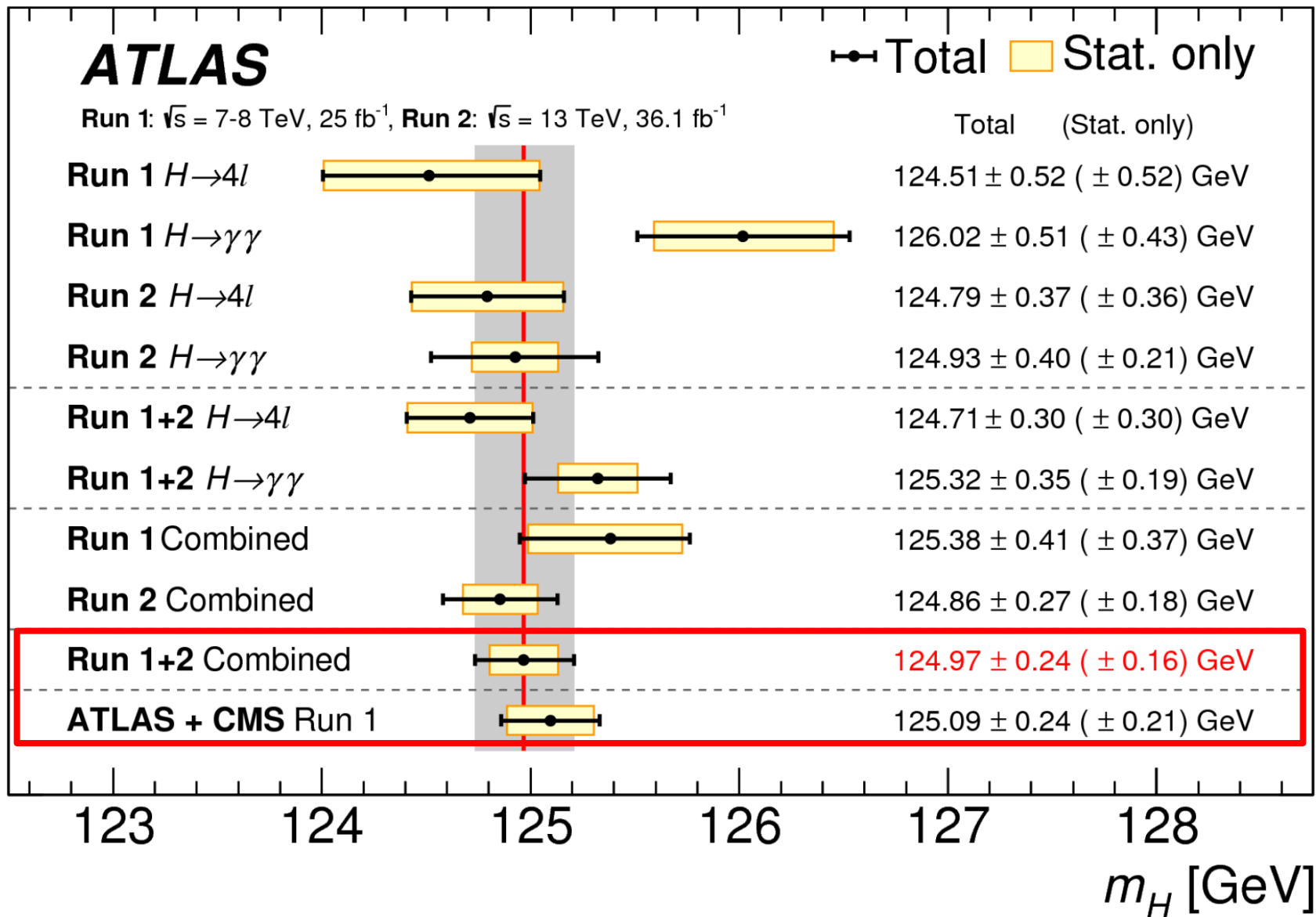
The signal is fitted to the data with a **floating signal strength**  $\rightarrow$  expect  $\mu = 1$  for SM Higgs



# Higgs mass: Results (@ Run 1 & Run 2)

Higgs decays to two photons and four leptons: well reconstructed  $m_H$  channels.

ATLAS & CMS





# ATLAS

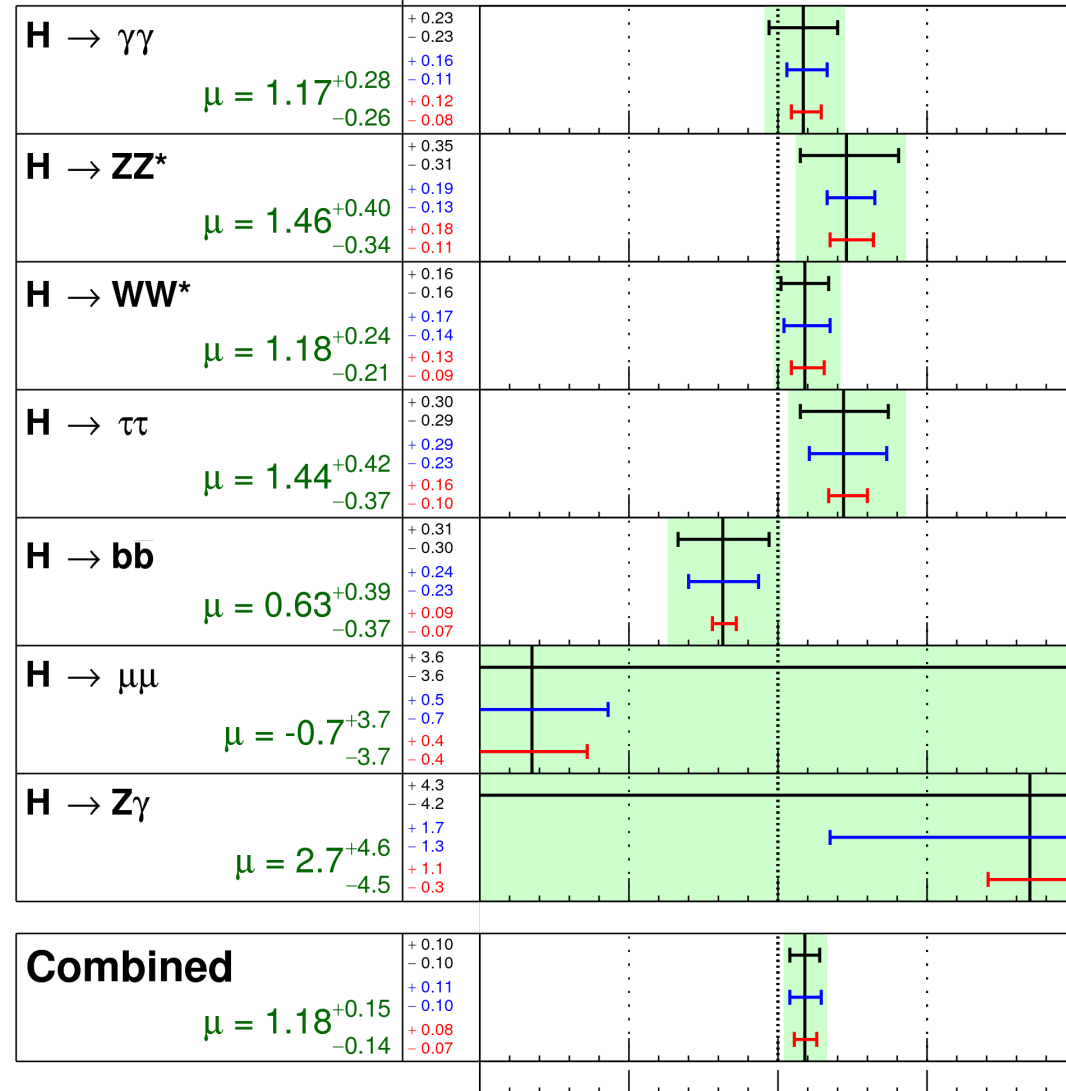
## ATLAS

$m_H = 125.36 \text{ GeV}$

—  $\sigma(\text{stat.})$   
 —  $\sigma(\text{sys inc.})$   
 —  $\sigma(\text{theory})$

Total uncertainty  
 $\pm 1\sigma$  on  $\mu$

# Higgs $\mu$



$\sqrt{s} = 7 \text{ TeV}, 4.5\text{-}4.7 \text{ fb}^{-1}$

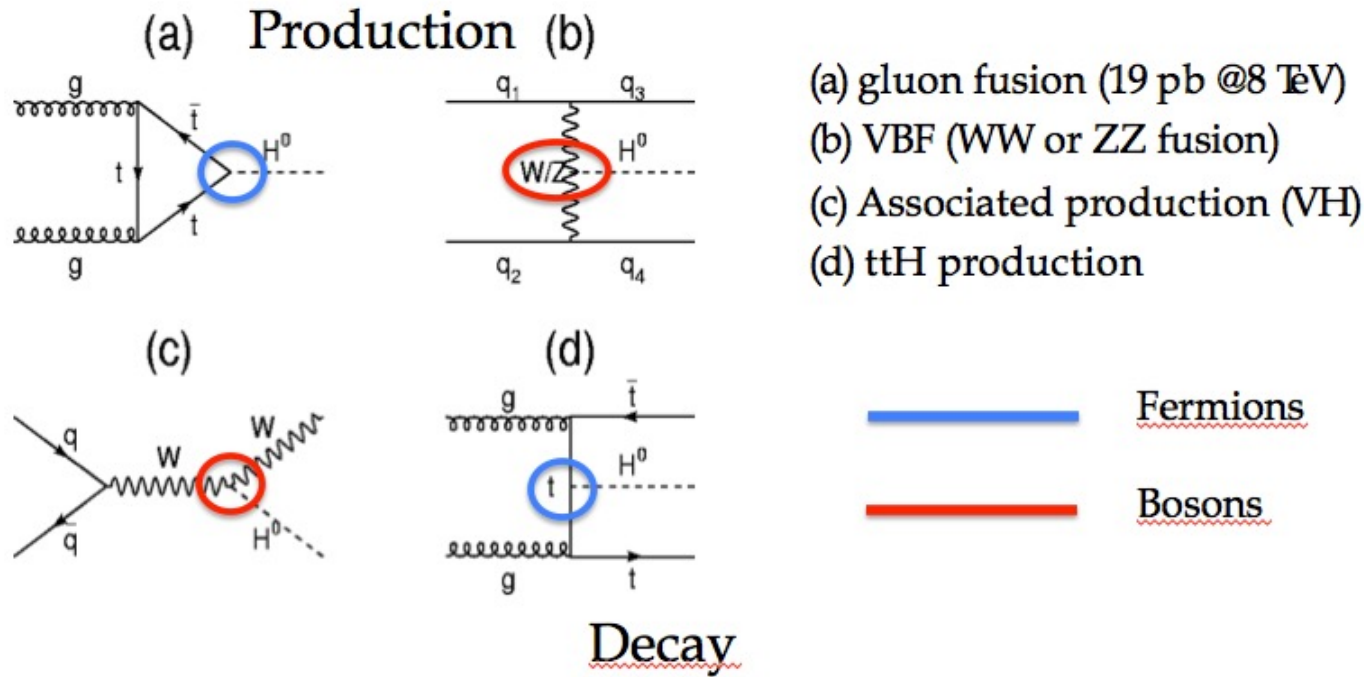
$\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$

Signal strength ( $\mu$ )



# Correlation Between Production and Decay of the Higgs

Toni Baroncelli: The Discovery of the Higgs

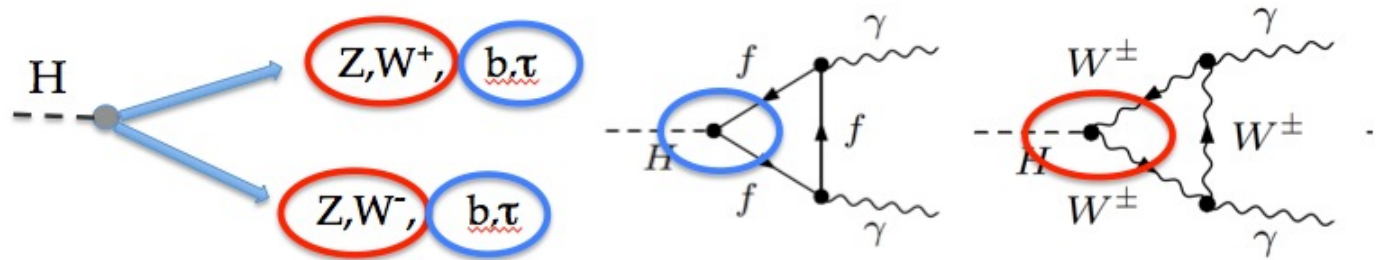


Same vertices in Higgs production and decay → in the SM they are dependent.

Use

$$\mu = \frac{\sigma^{measured} \times BR^{measured}}{\sigma^{SM} \times BR^{SM}}$$

→ if one vertex is modified by a correction factor **in production** the same correction factor has to modify also the corresponding vertex **in decay**



Use all possible measurements in the different categories and parametrise the correction factor with the signal strengths

**One unique  $\mu$**

The combination will give a better indication of how well SM describes data





# The Analysis Model

$$n_s^c = \left( \sum_{i,f} \underbrace{\mu_i}_{\text{production}} \cdot \sigma_i^{SM} \cdot A_{if}^c \cdot \varepsilon_{if}^c \cdot \underbrace{\mu_f}_{\text{decay}} \cdot BR_f^{SM} \right) \cdot \mathcal{L}^c$$

$n_s^c$  = number of events selected in one category  $c$   
Category is a sample of events selected by analysis cuts

- $\mu_i$  is the ratio between the observed cross section and the one predicted by the SM
- The production index  $i \in \{ggH, VBF, VH, ttH\}$  and the decay index  $f \in \{\gamma\gamma, WW, ZZ, bb, \tau\tau\}$
- $\sigma_i^{SM}$  and  $BR_f^{SM}$  production cross sections, decay branching fractions for a SM Higgs
- $A_{if}^c$  and  $\varepsilon_{if}^c$  are the signal acceptance and the reconstruction efficiency for given production and decay mode in the category  $c$ .
- $\mathcal{L}^c$  is the integrated luminosity used for that specific category.

Combination = fit  $\mu_i$  and  $\mu_i \rightarrow$  best agreement between data and (modified) SM prediction in different categories.

Includes Signal Regions and Control Regions

There are different ways of combination  $\rightarrow$  different assumptions ( $\sim$  simplifications)



# How Complex is the Fit?

$$n_S^C = \left( \sum_{i,f} \mu_i \cdot \sigma_i^{SM} \cdot A_{if}^C \cdot \varepsilon_{if}^C \cdot \mu_f \cdot BR_f^{SM} \right) \cdot \mathcal{L}^C$$

$i \in \{ggH, VBF, VH, ttH\}$

$f \in \{\gamma\gamma, WW, ZZ, bb, \tau\tau\}$

5 values of  $i$  ( $VH=WH+WZ$ )

5 values of  $f$

	$\gamma\gamma$	ZZ (4 $\ell$ )	WW ( $\ell\nu\ell\nu$ )	$\tau^+\tau^-$	$b\bar{b}$
ggF (high $p_T^H$ )	A	A	—	A	—
ggF (incl. or low $p_T^H$ )	A-C	A-C	A-C	A-C	—
ggF 1-jet	—	C	A-C	C	—
VBF	A-C	A-C	A-C	A-C	C
WH (1- $\ell$ )	A-C	A	A-C	C	A-C
WH (two jets)	A-C	A-C	A-C	—	—
ZH (0- $\ell$ )	A-C	A	—	—	A-C
ZH (2- $\ell$ )	A-C	A	A-C	C	A-C
ZH (two jets)	A-C	A-C	A-C	—	—
ttH (1- $\ell$ )	A-C	—	A-C	A-C	A-C
ttH (2- $\ell$ )	—	—	A-C	A-C	A-C
ttH (hadronic)	A-C	—	—	—	A

A=ATLAS, C=CMS

Possible ways of fitting:

1.  $i \times f$  makes 25 free parameters
2. One reference process (characterised by  $\sigma_{ref}$  and  $BR_{ref}$ ) +  $\sigma_f/\sigma_{ref} + BR_f/BR_{ref}$  ( $\rightarrow 1 + 4 + 4$  parameters)
3. Further assumptions (see later)
4. Effective Lagrangian, introduce vertex modifiers



# Signal Strength ( $\mu$ )

Production mechanism	$ggF$	$VBF$	$WH/ZH$	$t\bar{t}H$
Events produced at the LHC	500 K	40 K	20 K	3 K
Selected events	O(500)	O(500)	O(50)	
Events produced at the Tevatron	10 K		2 K	

For each decay channel “c” we define categories to maximise the sensitivity of the analysis to one particular production mode.

However a mixture of different mechanisms in one category is inevitable.  
 → This implies the cross section of one category is not the cross section of only one production mechanism.

$$n_s^c = \left( \sum_{i,f} \mu_i \cdot \sigma_i^{SM} \cdot A_{if}^c \cdot \epsilon_{if}^c \cdot \mu_f \cdot BR_f^{SM} \right) \cdot \mathcal{L}^c$$

Where  $m_c$  is ration between measured & expected events in that category and

$$i = gg, VBF, VH, t\bar{t} \quad \text{and} \quad f = gg, WW, ZZ, bb, tt$$

Measurement of  $\mu_c$  gives an indication of how well SM describes data

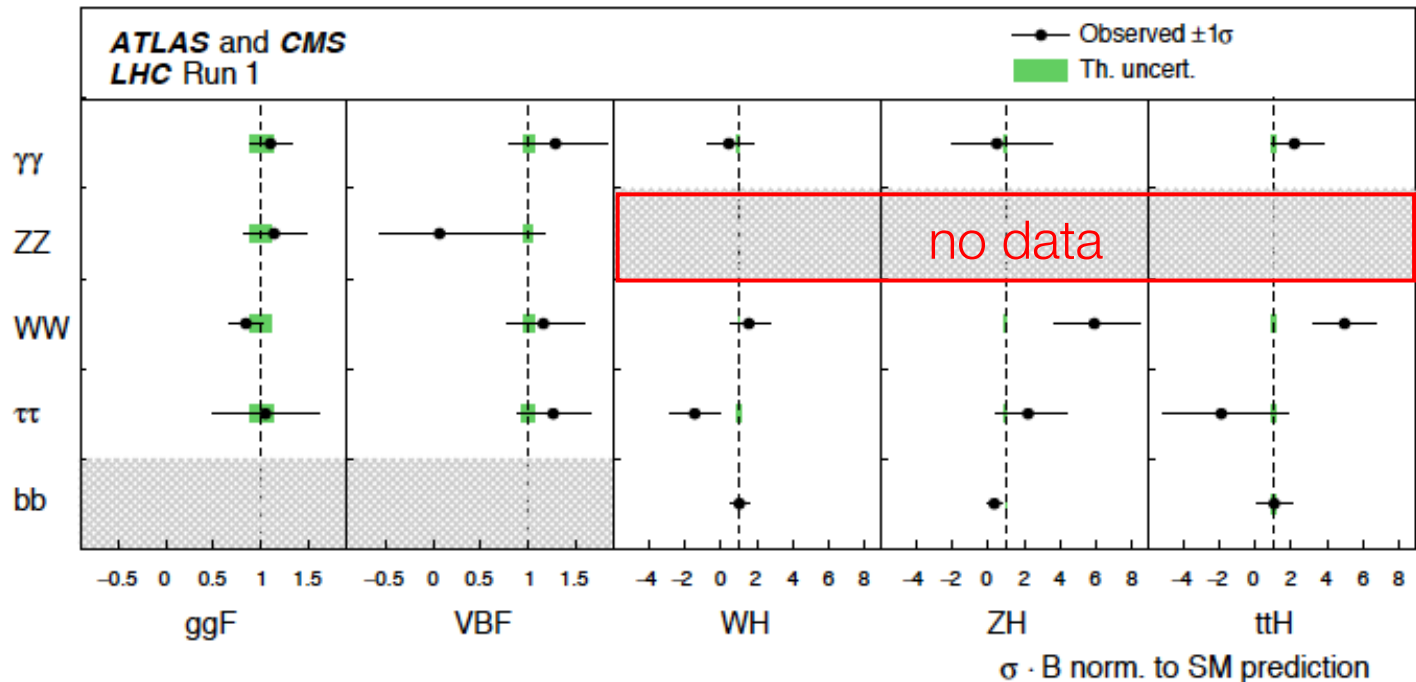


# Fit type 1 most general case

	$\gamma\gamma$	$ZZ (4\ell)$	$WW (l\nu l\nu)$	$\tau^+\tau^-$	$b\bar{b}$	Comb.
ggF	$1.10^{+0.22+0.07}_{-0.21-0.05}$	$1.13^{+0.33+0.09}_{-0.30-0.07}$	$0.84^{+0.12+0.12}_{-0.12-0.11}$	$1.00^{+0.4+0.4}_{-0.4-0.4}$	—	$1.03^{+0.16}_{-0.14}$
VBF	$1.3 \pm 0.5^{+0.2}_{-0.1}$	$0.1^{+1.1+0.2}_{-0.6-0.2}$	$1.2^{+0.4+0.2}_{-0.3-0.2}$	$1.3^{+0.3+0.2}_{-0.3-0.2}$	—	$1.18^{+0.25}_{-0.23}$
WH	$0.5^{+1.3+0.2}_{-1.2-0.1}$	—	$1.6^{+1.0+0.6}_{-0.9-0.5}$	$-1.4^{+1.2+0.7}_{-1.1-0.8}$	$1.0^{+0.4+0.3}_{-0.4-0.3}$	$0.89^{+0.40}_{-0.38}$
ZH	$0.5^{3.0}_{-2.5}^{+0.5}_{-0.2}$	—	$5.9^{+2.3+1.1}_{-2.1-0.8}$	$2.2^{+2.2+0.8}_{-1.7-0.6}$	$0.4^{+0.3+0.2}_{-0.3-0.2}$	$0.79^{+0.38}_{-0.36}$
ttH	$2.2^{1.6}_{-1.3}^{+0.2}_{-0.1}$	—	$5.0^{+1.5+1.0}_{-1.5-0.9}$	$-1.9^{+3.2+1.9}_{-2.7-1.8}$	$1.1^{+0.5+0.8}_{-0.5-0.8}$	$2.3^{+0.7}_{-0.6}$
Comb.	$1.14^{+0.19}_{-0.18}$	$1.29^{+0.26}_{-0.23}$	$1.09^{+0.18}_{-0.16}$	$1.11^{+0.24}_{-0.22}$	$0.70^{+0.29}_{-0.27}$	$1.09^{+0.11}_{-0.10}$

$$\text{Signal Strength } \mu = \frac{(\sigma \cdot B)_{\text{obs}}}{(\sigma \cdot B)_{\text{SM}}}$$

- No significant deviation from the SM
- 25 – 5 = 20 parameters determined





# Fit Type 2: Fewer Parameters

Take the Higgs produced via ggF and decaying to ZZ as reference

$i = \text{production}, f = \text{decay}$

$$\mu_i = \frac{\sigma_i}{\sigma_{ggF} \sigma_i^{SM}} \times \mu_{gg \rightarrow H \rightarrow ZZ}$$

$$\mu_f = \frac{BR_i}{BR_{ZZ} BR_i^{SM}}$$

$$\mu_{gg \rightarrow H \rightarrow ZZ} = \mu_{ggF} \times \mu_{ZZ} \times \sigma_i^{SM} \times BR_f^{SM}$$

Then, the master formula applies for all  $i$  and  $f$  indices except when

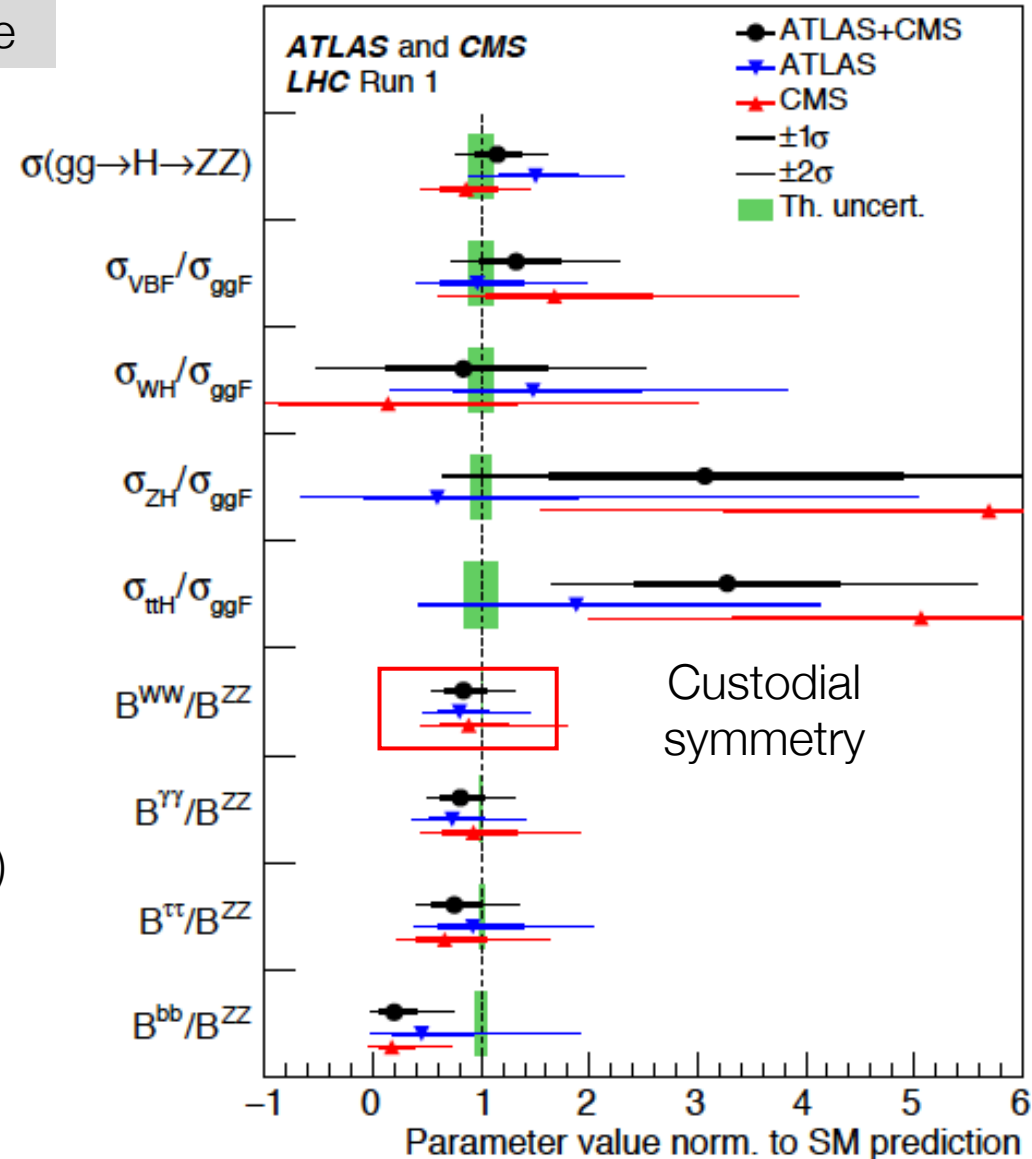
both  $i = \text{ggF}$  and  $f = \text{ZZ}$

$\rightarrow 8 + 1$  parameters

Improved precision in the fit, due to fewer parameters, shows no deviation from SM (assumption production and decay independent)

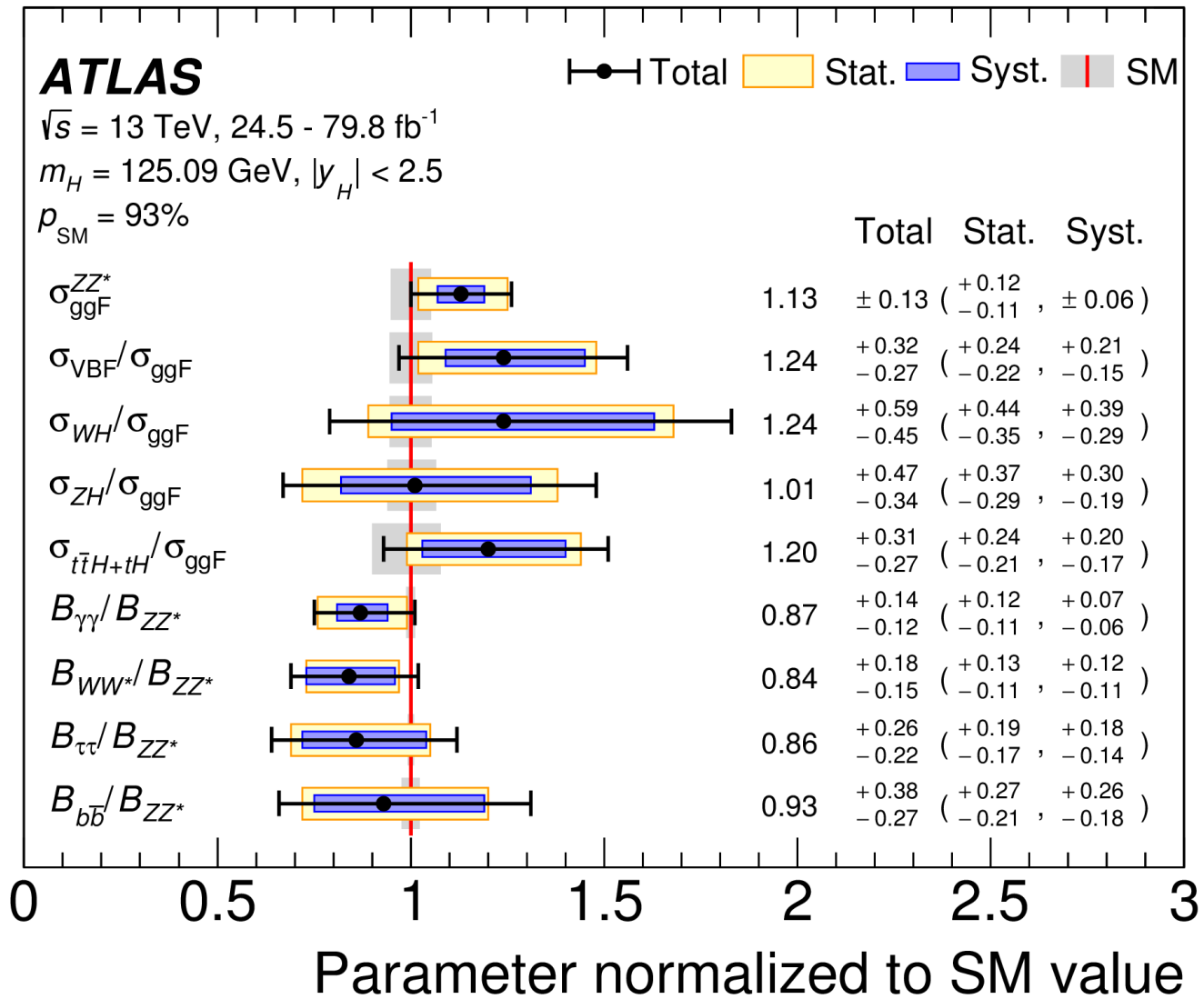
Finally if all deviations are included in a single parameter

$$\mu = 1.09 \pm 0.07 \pm 0.04(\text{expt}) \pm 0.03(\text{th. bkg}) \pm 0.07(\text{th. sig})$$



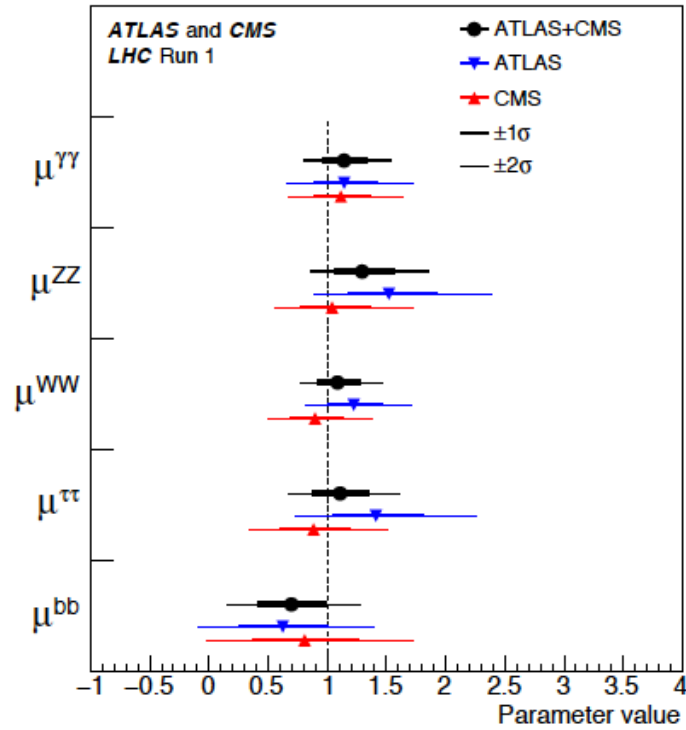


# ATLAS Recent Results: references

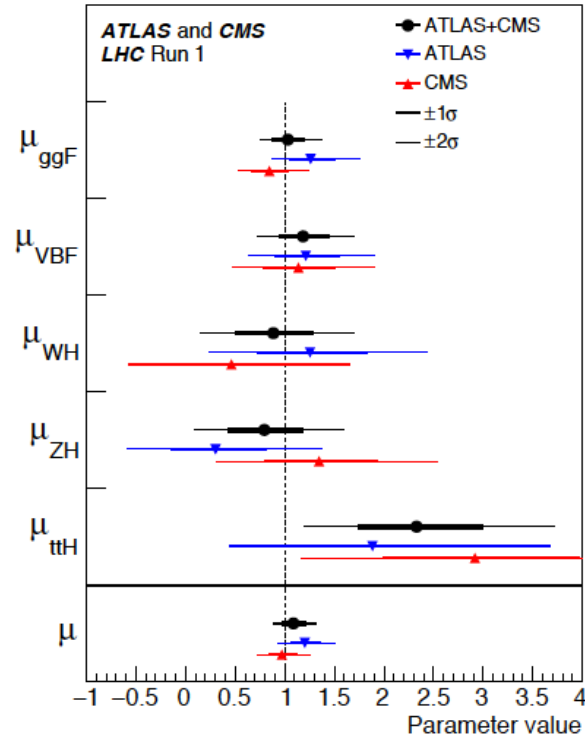




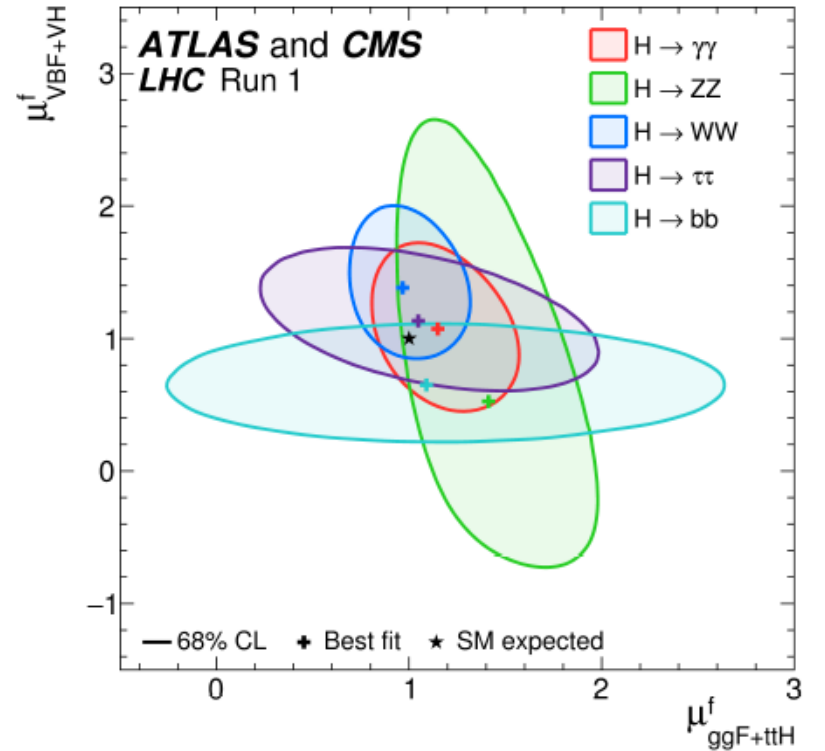
# More Simplifications.. Fewer Parameters



Assume all production cross sections are SM ones  $\rightarrow$  fit only modifiers of BR's



Assume all BRs are SM ones  $\rightarrow$  fit only modifiers of production cross-sections



One more assumption: vertices VBF & VH scale with one  $\mu$  and ggF and ttH with another  $\mu$

No (significant) deviation observed so far, need more statistics !  $\rightarrow$  LHC at High Luminosity



# A Different Approach

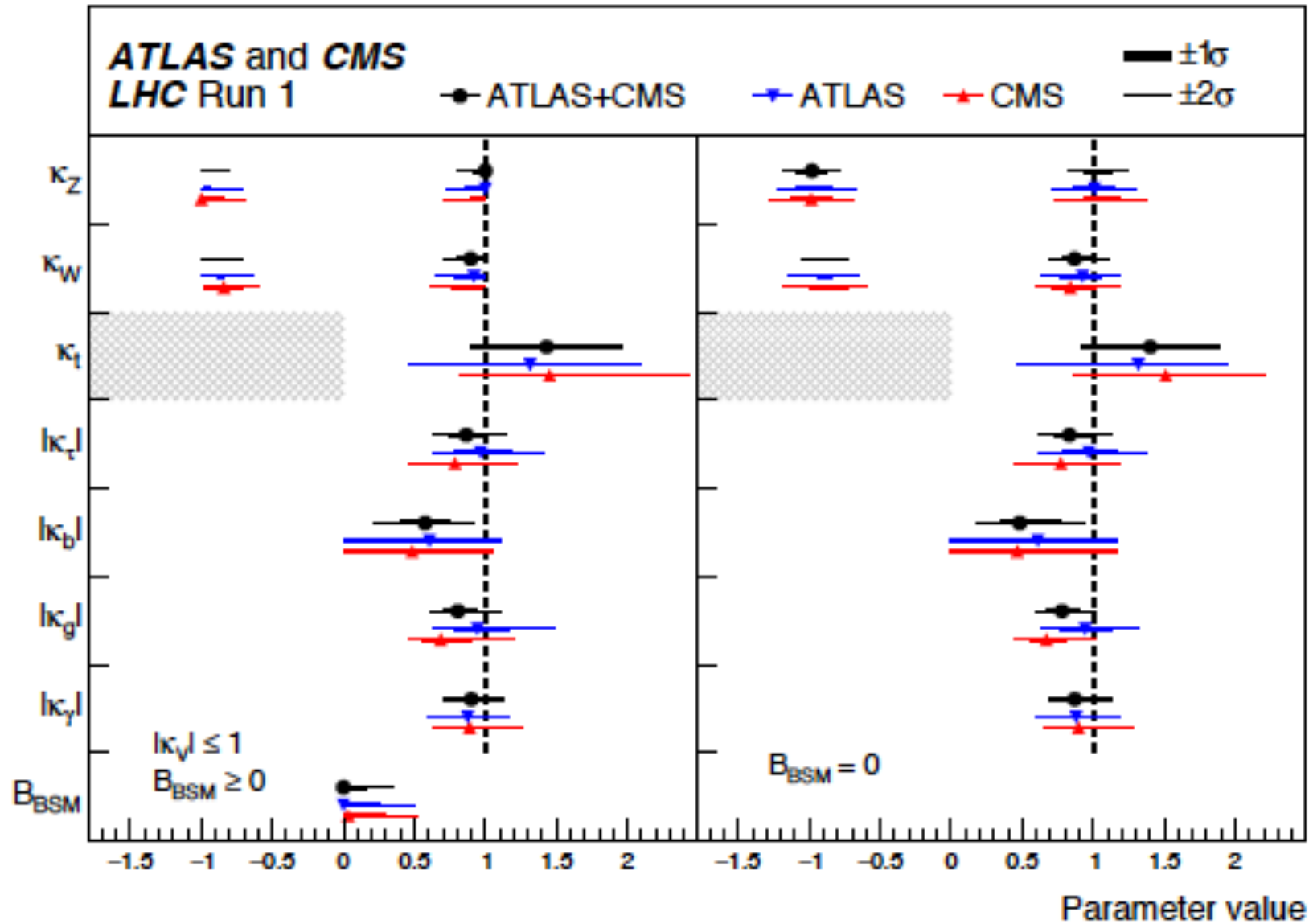
Elaborate even more: introduce modifiers  $k_x$  of vertices in the Lagrangian.  
In this way production and decay vertices are treated in the same way.

$$\begin{aligned} \mathcal{L} = & \boxed{\kappa_3} \frac{m_H^2}{2v} H^3 - \boxed{\kappa_Z} \frac{m_Z^2}{v} Z_\mu Z^\mu H + \boxed{\kappa_W} \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} H \\ & - \boxed{\kappa_g} \frac{\alpha_s}{2\pi v} G_{\mu\nu}^a G^{\mu\nu} H + \boxed{\kappa_\gamma} \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H + \boxed{\kappa_{Z\gamma}} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H \\ & + \boxed{\kappa_{VV}} \frac{\alpha}{2\pi v} \left( \cos^2 \theta_W Z_{\mu\nu} Z^{\mu\nu} + 2 W_{\mu\nu}^+ W^{-\mu\nu} \right) H \\ & - \left( \boxed{\kappa_t} \sum_{f=u,c,t} \frac{m_f}{v} f \bar{f} + \boxed{\kappa_b} \sum_{f=d,s,b} \frac{m_f}{v} f \bar{f} - \boxed{\kappa_\tau} \sum_{f=e,\mu,\tau} \frac{m_f}{v} f \bar{f} \right) H. \end{aligned}$$





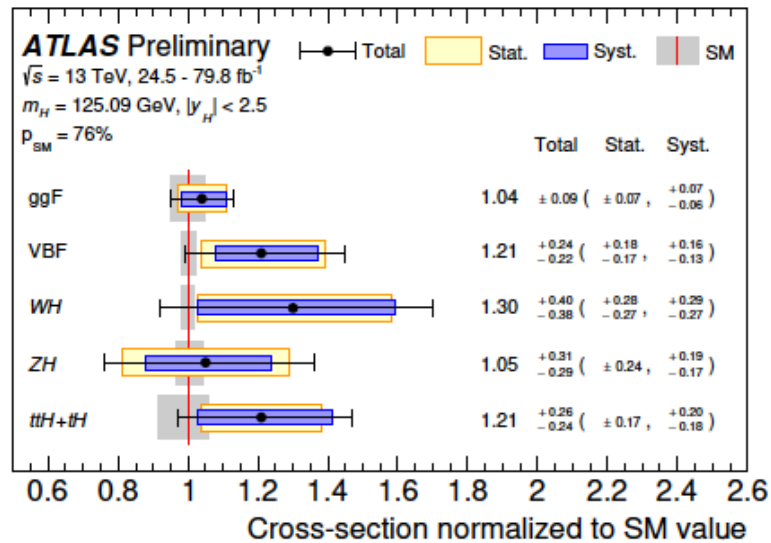
# Final Result





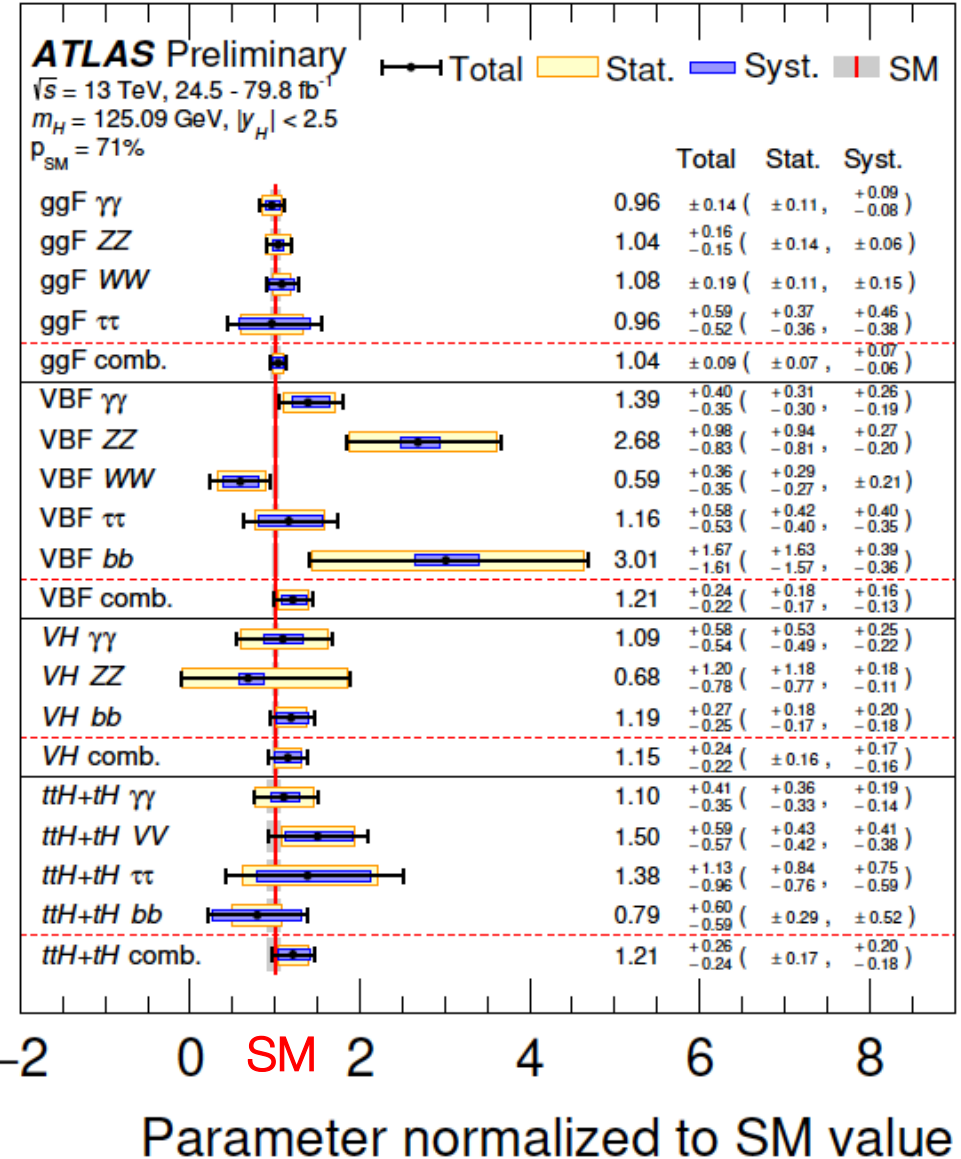
# Latest Results by ATLAS

Analysis	Integrated luminosity (fb <sup>-1</sup> )
$H \rightarrow \gamma\gamma$ (including $t\bar{t}H, H \rightarrow \gamma\gamma$ )	79.8
$H \rightarrow ZZ^* \rightarrow 4\ell$ (including $t\bar{t}H, H \rightarrow ZZ^* \rightarrow 4\ell$ )	79.8
$H \rightarrow WW^* \rightarrow e\nu\mu\nu$	36.1
$H \rightarrow \tau\tau$	36.1
$VH, H \rightarrow b\bar{b}$	79.8
VBF, $H \rightarrow b\bar{b}$	24.5 - 30.6
$H \rightarrow \mu\mu$	79.8
$t\bar{t}H, H \rightarrow b\bar{b}$ and $t\bar{t}H$ multilepton	36.1
$H \rightarrow$ invisible	36.1
Off-shell $H \rightarrow ZZ^* \rightarrow 4\ell$ and $H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	36.1



$$\mu = 1.11^{+0.09}_{-0.08} = 1.11 \pm 0.05 \text{ (stat.)}^{+0.05}_{-0.04} \text{ (exp.)}^{+0.05}_{-0.04} \text{ (sig. th.)} \pm 0.03 \text{ (bkg. th.)}$$

$$(\sigma \cdot B)_{obs} / (\sigma \cdot B)_{SM}$$

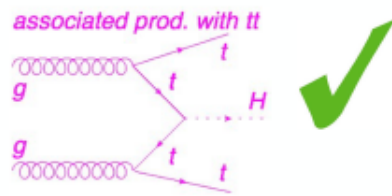
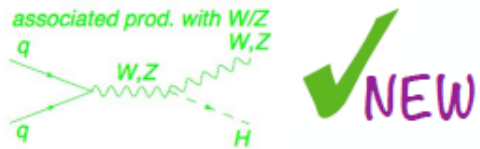
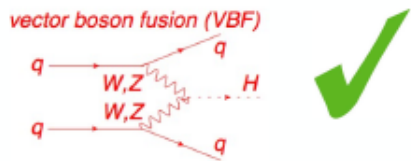
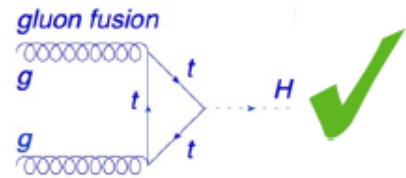




## Conclusions

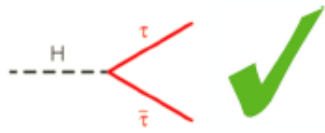
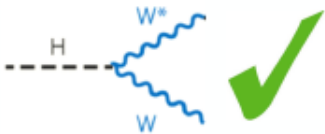
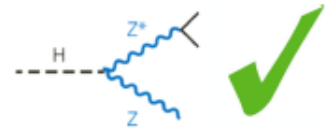
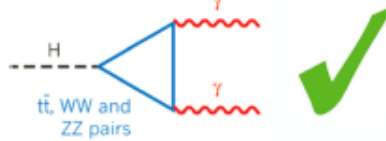
ATLAS and CMS

### Production



✓ = observed

### Decays



- Thanks to the first 36-80  $\text{fb}^{-1}$  of Run-2 data:
- The bosonic decay channels entered a precision era ( $\sim 3\times$  improvement w.r.t. Run-1)
- Direct observation achieved for all main production and decay modes!
- Direct confirmation of coupling to all 3rd generation fermions (top-quark, **bottom-quark**, taus)
- Sensitivity to double Higgs production approaching  $10 \times \text{SM}$  Higgs to 2 Higgses
- Higgs physics an important indirect probe for New Physics: so far no deviations from SM...
- But still at the beginning of a long journey! Only analyzed  $<3\%$  of the final LHC luminosity.



# *Higgs Discovery at LHC Part*

*Collider Physics*  
*Toni Baroncelli*  
*Haiping Peng*  
*USTC*

*End of Higgs Discovery at LHC Part*